# One-shot Multiparty Purity Distillation SuSTech-Nagoya workshop on Quantum Science

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## A very high level problem description:

#### $\rightarrow$ Single party:

Given a quantum mixed state  $\rho^A$  how much purity can Alice distill by application of only <u>local unitaries</u> ? (Devetak, 2004)

#### $\rightarrow\,$ Two parties:

Given a mixed quantum state  $\rho^{AB}$  where state on system A is with Alice and that on B is with Bob, what is the amount of **TOTAL** purity that can Alice and Bob can distill using only <u>local unitaries</u>?

#### $\rightarrow\,$ Multi-party generalization

Given a multipartite density matrix  $\rho^{A_1,A_2,...,A_n}$ , where *n* parties  $A_1,...,A_n$  are given the marginal states each, how much **TOTAL** purity can they distill by application of only <u>local unitaries</u>?

#### A detailed description of the problem:

Given a density matrix (say to Alice)  $\rho^{A}$  on an |A| dimensional Hilbert space :

- Distill purity  $\Leftrightarrow$  Extract <u>pure states</u>, specifically  $|0\rangle^{A_p}$  where  $|A_p| \leq |A|$ . Trivial checks:  $|A_p| = |A| \Rightarrow \rho^A = |0\rangle^A \langle 0|$ ;  $|A_p| = 0 \Rightarrow \rho^A = \frac{I^A}{|A|}$ .
- Allowed operations to Alice: Only a unitary operator on *A*; If Alice wish to use ancilla, then she can do so *only catalytically*, i.e., with a promise to return it at the end of the protocol.
- Pure state ⇔ |0⟩; since, any other pure state can be created by applying another local unitary.
   Aliter: Not charged for local unitary operations.
- Multi-party case: Allow <u>unidirectional</u> classical communication. (Reason to be described shortly!)

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## Single Party Case: (aka Local Protocol)

Aim: Find maximum 
$$\log |A_p|$$
 s.t. for a small  $\epsilon > 0$ :  
 $\|U^A \rho^A (U^{\dagger})^A - |0\rangle \langle 0|^{A_p} \otimes \rho'^{A/A_p}\|_1 \le \epsilon$  (1)

Definition: A rate  $R := \log |A_p|$  is said to be  $\epsilon$ -achievable for one-shot purity distillation protocol  $\iff \exists$  a unitary  $U^A$  that satisfies Eq (1).

#### Theorem

Given a state  $\rho^{A}$  and  $\epsilon > 0$ , there exists a unitary on the system  $A \rightarrow A_{p} \otimes A_{q}$  such that:  $\|U\rho U^{\dagger} - (\Pi\rho\Pi)^{A_{q}} \otimes |0\rangle \langle 0|^{A_{p}}\|_{1} \leq \epsilon$  and

$$\boxed{ \begin{array}{c} R = \log |A_{\rho}| = \log |A| - \widetilde{H}_{\max}^{\epsilon}(A)_{\rho} \ , \ where} \\ \hline \widetilde{H}_{\max}^{\epsilon}(A)_{\rho} := \log \operatorname{supp}(\rho'^{A}) \end{array} }$$

 $\rho'^A$  is the matrix obtained by zeroing out the smallest eigen vectors of  $\rho$  that sum up to  $\leq \epsilon.$ 

#### Proof of local protocol

$$A \cong A_p \otimes A_q : A_q = eigen space(\rho')$$
 (thm. above).

• Thus, 
$$|A_q| = 2^{\widetilde{H}_{\max}^{\epsilon}(A)_{\rho}}$$
.

•  $\Pi$  projector onto  $A_q 
ightarrow (1-\epsilon)$  probability eigen space(
ho) :

$$\operatorname{Tr}(\mathbf{\Pi}) = |\mathbf{A}_{\mathbf{q}}|; \operatorname{Tr}[\mathbf{\Pi}\rho] \geq \mathbf{1} - \epsilon.$$

- $:: [\Pi, \rho] = 0 \Rightarrow A_q \cong \text{eig. space}(\Pi \rho \Pi).$ Hence, eigen vectors  $\{|i\rangle^{A_q}\}$  of  $\Pi \rho \Pi \in Basis(A_q).$
- Let { |i'⟩<sup>A</sup>}<sub>i'∈A</sub> = eigen vectors(ρ<sup>A</sup>).
   ∴ [Π, ρ] = 0 ⇒ ∃ embedding from span { |i⟩<sup>A<sub>q</sub></sup>} into span { |i'⟩<sup>A</sup>} as:

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## Proof of local protocol

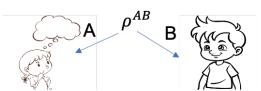
$$\begin{aligned} \frac{\left| U: A \to A_{p} \otimes A_{q}: \left| i' \right\rangle^{A} \mapsto \left| i \right\rangle^{A_{q}} \left| 0 \right\rangle^{A_{p}}; U(\Pi \rho \Pi)^{A} U^{\dagger} &= (\Pi \rho \Pi)^{A_{q}} \otimes \left| 0 \right\rangle \left\langle 0 \right|^{A_{p}}. \\ \| \rho^{A} - (\Pi \rho \Pi)^{A} \|_{1} &\leq \epsilon \Rightarrow \| (U \rho U^{\dagger})^{A} - (\Pi \rho \Pi)^{A_{q}} \otimes \left| 0 \right\rangle \left\langle 0 \right|^{A_{p}} \|_{1} \overset{uni.}{\leq} 2\sqrt{\epsilon} \\ R &= \log |A_{p}| = \log \frac{|A|}{|A_{q}|} = \log |A| - \widetilde{H}_{\max}^{\epsilon}(A)_{\rho} \quad [QED] \end{aligned}$$

• Devetak (2004) showed that in asymptotic iid setting:

 $\max \mathbf{R} = \log |\mathbf{A}| - \mathbf{H}(\mathbf{A})_{\rho}$ , with a matching converse.

• In either one-shot and asymptotic iid case, R=0, if  $\rho = \frac{I}{|A|}$  (intuitively).

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Alice and Bob want to distill max # pure states from their respective halves of  $\rho^{AB}$ .

• Naively, if they apply the local (single party) protocol, they can distill a total purity of  $\log |A| - \widetilde{H}_{max}^{\epsilon}(A)_{\rho} + \log |B| - \widetilde{H}_{max}^{\epsilon}(B)_{\rho}$ .

• Can classical correlations help to increase the rate?

- e.g. let  $|A| = |B| = 2, \epsilon < \frac{1}{2}$ ;  $\rho^{AB} = \frac{|0\rangle\langle 0|^A \otimes |0\rangle\langle 0|^B}{2} + \frac{|1\rangle\langle 1|^A \otimes |1\rangle\langle 1|^B}{2}$ , then  $\widetilde{H}^{\epsilon}_{max}(A)_{\rho} = 1 = \widetilde{H}^{\epsilon}_{max}(B)_{\rho}$ .  $\Rightarrow R_{AB} = 0$  by local protocol!
- If classical communication is allowed from Alice → Bob, then Alice can simply send her A system to Bob and the Bob can apply the controlled unitary CNOT on A(rx. from Alice), B. Then, Alice distills 0 purity but Bob distills 1 bit of purity and R<sub>AB</sub> = 1!!

- Can do better by utilizing classical correlations in  $\rho^{AB}$ .
- A way to create such a classical correlated state is via <u>measurement</u>!
- Thus, further allowed: catalytic use of ancilla for Alice; Alice <u>1-way classical</u> <u>communication</u> Bob; Alice can perform measurements, leading to:

#### Theorem

Given  $\rho^{AB} \xrightarrow{purification} |\rho\rangle^{RAB}$ . A one-shot achievable total rate for distilling purity under the above conditions is:

 $R_{AB} = \log |A||B| - I_{\mathsf{max}}^{\epsilon}(X:RB) - \widetilde{H}_{\mathsf{max}}^{\epsilon}(B) + I_{H}^{\epsilon_{0}/2}(X:B) + 0(\epsilon,\epsilon_{0}) - O(1)$ 

for a given  $\epsilon > 0$  and  $\epsilon_0 = O(\epsilon)$  and the above entropic quantities are evaluated wrt the  $\rho^{XRB} = (I^{RB} \otimes \Lambda^A) |\rho\rangle^{RAB}$ , where  $\Lambda^A$  are rank one measurement operators on A. The maximum achievable rate can be obtained by optimizing over the  $\Lambda^A$ .

• Asymptotic iid rate (with a matching converse, Devetak, 2004) =  $\log |A| - H(A) + \log |B| - H(B) + \lim_{n \to \infty} \frac{\max_{n \to \infty} I(X^n; B^n)_{(I^B \otimes \Lambda^A) \otimes n(\rho^{AB}) \otimes n}}{n}$ Protocol for one-shot achievability:

Step 1: Alice applies normalized rank 1-POVM  $\Lambda^{A} = \{|\psi_{x}\rangle \langle \psi_{x}|\}_{\{x \in \mathcal{X}\}}$ on her system *A* "coherently" by borrowing log  $|\mathcal{X}|$  ancilla  $\rightarrow$  applies unitary:  $U_{1}^{AX} : |0\rangle^{X} |\rho\rangle^{ABR} (\stackrel{Schmidt}{=} \sum_{i} \lambda_{i} |i\rangle^{A} |\varphi_{i}\rangle^{BR}) \mapsto \sum_{x,i} |x\rangle^{X} (|\psi_{x}\rangle \langle \psi_{x}| |a_{i}\rangle)^{A} |\tilde{\varphi}_{i}\rangle^{RB} = \sum_{x} |x\rangle^{X} |\psi_{x}\rangle^{A} |\tilde{\varphi}_{x}\rangle^{RB}$ . Step 2: Alice now applies controlled unitary  $\sum_{x} |x\rangle \langle x|^{X} \otimes U_{x}^{A} : |x\rangle^{X} |\psi_{x}\rangle^{A} \mapsto |x\rangle^{X} |0\rangle^{A}$ , thus distilling log  $|A| - \log |\mathcal{X}|$  purity. (Note: log  $|\mathcal{X}|$  catalyst is be returned.) Remark: Now the state is  $\sum_{x} |x\rangle^{A} |\tilde{\varphi}_{x}\rangle^{RB} \Rightarrow I(X; RB) = H(A)_{\rho}$ .

Step 3: Alice 
$$\xrightarrow{\{|x\rangle\}}{\text{measure}} \tau^{XB} = \sum_{x} p_X(x) |x\rangle \langle x|^X \otimes \rho_x^B$$
.  
Note:  $\tau^B = \sum_{x} p_X(x) \rho_x^B = \rho^B$  (unaltered).

Alice local protocol on X
Alice's log |X| - identified in X
Alice's overall purity is |A| - log |X| + log |X| - H̃<sub>max</sub><sup>ϵ</sup>(X) = log |A| - H̃<sub>max</sub><sup>ϵ</sup>(X).
Resultant state (S = (1 - ϵ) prob. subset of x's under p<sub>X</sub>):
σ<sup>X<sub>1</sub>B</sup> = Σ<sub>x∈S</sub> P<sub>X</sub>(x)|x⟩⟨x|<sup>X<sub>1</sub>⊗ρ<sup>B</sup></sup><sub>x</sub>; Π<sub>S</sub> : proj.(S). Alice now aims at transferring these correlations to Bob for him to be able to distill some purity from these classical correlations. Naively: Alice sends X<sub>1</sub> to Bob and Bob tries to distill purity from X<sub>1</sub>!

Naively: Alice sends  $X_1$  to Bob and Bob tries to distill purity from  $X_1$ ! **Won't work naively:** S is a high prob. set, with almost no redundancies. What else can Bob do to distill purity from classical correlations? New idea: One shot classical data compression with quantum side information.

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Tool: One-shot classical data-compression with quantum side information

Lemma (Chakraborty et al., 2022)

Given 
$$\rho^{XB} = \sum_{x} P_X(x) |x\rangle \langle x|^X \otimes \rho_x^B \exists$$
 bijection  $\sigma : \mathcal{X} \to [M] \times [N]$ , s.t.  
a) for  $\epsilon > 0$ ,  $\log |N| \le I_H^{\epsilon}(X : B) + 2\log \epsilon$ ;  
b)  $\rho^{XB} \xrightarrow{\sigma} \sum_{m,n} P_{MN}(m,n) |m,n\rangle \langle m,n|^{MN} \otimes \rho_{mn}^B$ ; then  $\forall m \in [M] \exists a$   
POVM  $\Theta_n(m)$  s.t. for  $\epsilon_0 := O(\epsilon^{1/4})$   
 $\sum_{m,n} P_{MN}(m,n) \|\rho_{mn}^B - \sqrt{\Theta_n(m)}\rho_{mn}^B \sqrt{\Theta_n(m)}\|_1 \le \epsilon_0$ .

Proof: Binning at the encoder; one-shot cq packing lemma at the decoder. Corollary

For 
$$\rho^{MNB} \exists$$
 a unitary  $W^{MNB}$ :  $\|\operatorname{Tr}_{BM} (W^{MNB} \cdot \sigma^{MNB}) - |0\rangle \langle 0|^{N} \|_{1} \leq 2\epsilon_{0}$ 

Step 4: 
$$\tau^{X_1B} \xrightarrow{\text{Alice}}_{\text{bij},\sigma} \sigma^{MNB} = \sum_{m,n} P_{MN}(m,n) |m,n\rangle \langle m,n|^{MN} \otimes \rho_{mn}^B$$
  
s.t.  $\log |N| \leq I_H^{\sqrt{\epsilon_0}}(X_1:B) + \log \epsilon_0$ ;  $\xrightarrow{MN}$  Bob using  $\tilde{H}_{\max}^{\epsilon}(X)$  bits.  
Step 5: Bob applies unitary  $W^{MNB}$  from Cor. [4] and extracts  
 $\log |N| \approx I_H^{\sqrt{\epsilon_0}}(X_1:B) + \log \epsilon_0$  purity. Further, by lemma [3] and  
cor. [4] and fact that  $S$  being a high prob. set implies: (non-trivial)  
 $\|\sum_x p_X(x)\rho_X^B - \frac{1}{\text{Tr}[\Pi_S \tau]}\sum_{x \in S} p_X(x)\rho_X^B\|_1 \leq O(\epsilon_0^{1/8}).$   
 $\Rightarrow$  Bob's state  $\approx \rho^B$ ; Distills  $I_H^{\sqrt{\epsilon_0}}(X_1;B)$  purity.  
Lemma:  $I_H^{\sqrt{\epsilon_0}}(X_1:B)_{\sigma^{X_1B}} \geq I_H^{\epsilon_0/2}(X:B)_{\rho^{XB}} - O(1) + O(\log(1-\epsilon_0))$   
Step 6: Bob now applies the local protocol on approximate version of  
his local state and distill  $\log |B| - \tilde{H}_{\max}^{\epsilon}(B)$  purity.  
Thus,  $R_{AB} \stackrel{(\epsilon,\epsilon_0)}{\approx} \log |A| - \tilde{H}^{\epsilon}(X) + I_H^{\epsilon_0/2}(X:B) + \log |B| - \tilde{H}_{\max}^{\epsilon}(B)$ .

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#### Bipartite purity distillation and further

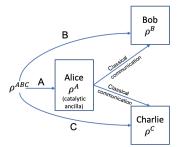
- Can recover the aysmptotic iid rate :: entropic quantities are smooth.
- The amount of classical communication  $\log |supp(X_1)| = \tilde{H}_{\max}^{\epsilon}(X)_{p_X}$ . Can we save on this?
  - YES! By using a modification to a recent **one-shot measurement compression** (Chakraborty et al., 2022).
  - Idea: Come up with a smaller outcome rank 1-measurement,  $\{\widetilde{\Lambda_y}^A\}_{\{y \in \mathcal{Y}\}}, \ |\mathcal{Y}| \leq |\mathcal{X}| \text{ from } \Lambda^A$ :

$$\left|I_{H}^{\sqrt{\epsilon_{0}}}(Y:B)_{\sigma} \geq I_{H}^{\epsilon_{0}/2}(X:B)_{
ho} - O(1) + O(\log(1-\epsilon_{0}))
ight|$$
 and find a

high prob. set (under  $p_Y$  obtained from  $\widetilde{\Lambda}$ )  $S_{\mathcal{Y}} : |S_{\mathcal{Y}}| \leq 2^{I_{\max}^{\epsilon}(X:RB)_{\rho}XRB}$ .

•  $\Rightarrow$  Catalytic ancilla  $\downarrow$ -classical communication  $\downarrow$  and Alice's purity from correlations  $\uparrow$  to log  $|\mathcal{Y}| - I^{\epsilon}_{\max}(X : RB)_{\rho^{XRB}}$ .

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Alice, Bob and Charlie want to distill max # pure states from their respective parts of  $\rho^{ABC}$ .

Need: A multiparty version of classical data compression with quantum side information at Bob and Charlie, respectively. Can modify a similar theorem from Chakraborty et al. (2022).
 Proposed rate is:

$$R_{ABC} \stackrel{\epsilon,\epsilon_{0},\epsilon_{1}}{\approx} \log |A| - I_{\max}^{\epsilon}(X,Y:B,C) + \log |B| - \tilde{H}_{\max}^{\epsilon}(B) + I_{H}^{\epsilon_{0}}(X:B) + \log |C| - \tilde{H}_{\max}^{\epsilon}(C) + I_{H}^{\epsilon_{1}}(Y:C)$$

## Conclusion and Future Prospects

- We have obtained the rate for one-shot purity distillation of pure states of form  $|0\rangle\,\langle 0|.$ 
  - We have used a modified version of the recent one-shot measurement compression theorem.
  - We have used the strategy for classical data compression with quantum side information.
  - Using the insights from above two lemmas, we have shown inequalities between the entropic quantities that show up in the rate expression.
- **Ongoing Work:** To prove the suggested rate for the tripartite case and in general, for the multipartite case.
- For future work: Investigate if this protocol or some minor variant of it can serve as an operational meaning to the so-called Grünwald information gain. This is definfined as:  $H(R)_{\sigma} H(B|X)_{\sigma}$  for the state  $\sigma^{RXB} := \sum_{x} \Lambda_{x}^{A \to B} (|\rho\rangle^{RA}) \otimes |x\rangle \langle x|^{X}$ , where  $\Lambda^{A \to BX}$  is a quantum instrument.

# Thank you and Questions ?? (To be out on arXiv shortly!!)

Aditya Nema

One-shot Multiparty Purity Distillation