

Langevin equation

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Exercise 5.1.7. Consider the Itô process satisfying

$$dX_t = -\beta X_t dt + \alpha dB_t, \quad X_0 = x_0 \quad (1)$$

for $\alpha \in \mathbb{R}$ and $\beta > 0$. This equation is called the Langevin equation. Note that this equation can be written equivalently

$$X_t = x_0 + \alpha B_t - \beta \int_0^t X_s ds.$$

Show that the solution of this equation reads

$$X_t = x_0 e^{-\beta t} + \alpha \int_0^t e^{-\beta(t-u)} dB_u.$$

Consider a deterministic differential equation as follows:

$$x'(t) + \beta x(t) = f(t),$$

for some constant β and function f . We solve it by multiplying both sides by $e^{\beta t}$ and get

$$\begin{aligned} e^{\beta t} x'(t) + \beta e^{\beta t} x(t) &= f(t) e^{\beta t} \\ \Leftrightarrow \frac{d}{dt} [e^{\beta t} x] (t) &= f(t) e^{\beta t} \\ \Leftrightarrow e^{\beta t} x(t) &= \int f(t) e^{\beta t} dt + C \\ \Leftrightarrow x(t) &= e^{-\beta t} \left[\int f(t) e^{\beta t} dt + C \right]. \end{aligned}$$

Taking inspiration from this, we try an ansatz $X_t = e^{-\beta t} Z_t$ for some Itô process $(Z_t)_{t \in [0, T]}$ such that $Z_0 = x_0$. Since we have

$$\begin{aligned} d[e^{-\beta t}] &= 0 dB_t - \beta e^{-\beta t} dt, \\ dZ_t &= V_t dB_t + D_t dt, \end{aligned}$$

Lemma 5.1.9¹ dictates that

$$\begin{aligned} dX_t &= d[e^{-\beta t} Z_t] = Z_t d[e^{-\beta t}] + e^{-\beta t} dZ_t + 0 \cdot V_t dt \\ &= -\beta e^{-\beta t} Z_t dt + e^{-\beta t} dZ_t \\ &= -\beta X_t dt + e^{-\beta t} dZ_t. \end{aligned}$$

Comparing this with (1), we obtain Z_t :

$$e^{-\beta t} dZ_t = \alpha dB_t \Leftrightarrow dZ_t = \alpha e^{\beta t} dB_t \Leftrightarrow Z_t = Z_0 + \alpha \int_0^t e^{\beta u} dB_u.$$

As a result, we obtain a solution of (1) by using the initial condition $Z_0 = x_0$:

$$X_t = e^{-\beta t} Z_t = x_0 e^{-\beta t} + \alpha \int_0^t e^{-\beta(t-u)} dB_u.$$

The process $(X_t)_{t \in [0, T]}$ satisfying (1) is also called the *Ornstein-Uhlenbeck process*.

Now, we want to find the expectation value $\mathbb{E}(X_t)$ and the autocovariance $\text{Cov}(X_t, X_s)$ of this process. Observe that

$$\int_0^T e^{2\beta u} du = \frac{e^{2\beta u}}{2\beta} \Big|_{u=0}^{u=T} = \frac{e^{2\beta T} - 1}{2\beta} < \infty.$$

Since the process $(e^{\beta u})_{u \in [0, T]}$ belongs to $M^2([0, T])$, we obtain the expectation value of X_t by using **Proposition 4.2.10**,

$$\mathbb{E}(X_t) = \mathbb{E}(x_0 e^{-\beta t}) + \alpha e^{-\beta t} \mathbb{E}\left(\int_0^t e^{\beta u} dB_u\right) = x_0 e^{-\beta t} + 0 = x_0 e^{-\beta t}.$$

Now, before finding the autocovariance $\text{Cov}(X_t, X_s)$, we prove that increments of a martingale $(M_t)_{t \in [0, T]}$ are uncorrelated [2, Exercise 5.4]. Consider any times t_1, t_2, t_3, t_4 such that $t_1 \leq t_2 \leq t_3 \leq t_4$. Then, we have by using the martingale property,

$$\mathbb{E}(M_{t_4} - M_{t_3} \mid \mathcal{F}_{t_2}) = \mathbb{E}(M_{t_4} \mid \mathcal{F}_{t_2}) - \mathbb{E}(M_{t_3} \mid \mathcal{F}_{t_2}) = M_{t_2} - M_{t_2} = \mathbf{0}.$$

Since $\mathcal{F}_{t_1} \subset \mathcal{F}_{t_2}$ and M_{t_1} is \mathcal{F}_{t_1} -measurable, M_{t_1} is also \mathcal{F}_{t_2} -measurable. As a result $M_{t_2} - M_{t_1}$ is \mathcal{F}_{t_2} -measurable and we have

$$\begin{aligned} \mathbb{E}[(M_{t_2} - M_{t_1})(M_{t_4} - M_{t_3})] &= \mathbb{E}\{\mathbb{E}[(M_{t_2} - M_{t_1})(M_{t_4} - M_{t_3}) \mid \mathcal{F}_{t_2}]\} \\ &= \mathbb{E}[(M_{t_2} - M_{t_1})\mathbb{E}(M_{t_4} - M_{t_3} \mid \mathcal{F}_{t_2})] \quad (\text{Proposition 3.1.3, 4.}) \\ &= \mathbb{E}[(M_{t_2} - M_{t_1}) \cdot \mathbf{0}] = 0. \end{aligned}$$

Note that $(M_t)_{t \in [0, T]}$ defined by $M_t = \int_0^t Y_u dB_u$ is a martingale for any adapted stochastic process $(Y_t)_{t \in [0, T]}$ belonging to $M^2([0, T])$ (**Theorem 4.3.1**). Then, using **Proposition 4.2.10** and the

¹All the statement in bold are from the lecture notes [1].

above property of martingales, we have for any $s \leq t$,

$$\begin{aligned} \mathbb{E} \left[\left(\int_0^t Y_u \, dB_u \right) \left(\int_0^s Y_u \, dB_u \right) \right] &= \mathbb{E} \left[\left(\int_0^s Y_u \, dB_u \right)^2 \right] + \mathbb{E} \left[\left(\int_s^t Y_u \, dB_u \right) \left(\int_0^s Y_u \, dB_u \right) \right] \\ &= \int_0^s \mathbb{E}(Y_u^2) \, du + \mathbb{E}[(M_t - M_s)(M_s - M_0)] \\ &= \int_0^s \mathbb{E}(Y_u^2) \, du. \end{aligned}$$

Now, coming back to the Ornstein-Uhlenbeck process $(X_t)_{t \in [0, T]}$, we have for any $s \leq t$,

$$\begin{aligned} \text{Cov}(X_t, X_s) &= \mathbb{E}[(X_t - \mathbb{E}(X_t))(X_s - \mathbb{E}(X_s))] \\ &= \mathbb{E} \left[\left(\alpha \int_0^t e^{-\beta(t-u)} \, dB_u \right) \left(\alpha \int_0^s e^{-\beta(s-u)} \, dB_u \right) \right] \\ &= \alpha^2 e^{-\beta(t+s)} \mathbb{E} \left[\left(\int_0^t e^{\beta u} \, dB_u \right) \left(\int_0^s e^{\beta u} \, dB_u \right) \right] \\ &= \alpha^2 e^{-\beta(t+s)} \int_0^s e^{2\beta u} \, du \\ &= \alpha^2 e^{-\beta(t+s)} \frac{e^{2\beta s} - 1}{2\beta} \\ &= \frac{\alpha^2}{2\beta} e^{-\beta(t-s)} (1 - e^{-2\beta s}). \end{aligned}$$

Furthermore, $\left(\int_0^t e^{\beta u} \, dB_u \right)_{t \in [0, T]}$ is a Gaussian process by **Theorem 4.3.4**. We know that if X is a Gaussian random variable, then $aX + b$ is also a Gaussian random variable for any a, b with $a \neq 0$. So $(X_t)_{t \in [0, T]}$ is also a Gaussian process since $X_t = x_0 e^{-\beta t} + \alpha e^{-\beta t} \int_0^t e^{\beta u} \, dB_u$ for any $t \in [0, T]$.

If $\alpha = \beta = 1$ and $x_0 = 0$, then we get $\mathbb{E}(X_t) = 0$ and $\text{Cov}(X_t, X_s) = \frac{1}{2} e^{-(t-s)} (1 - e^{-2s})$ for $s \leq t$, which is the Gaussian process given in **Example 2.2.6**.

Digression: Langevin equation in physics

Originally, the Langevin equation was proposed as the equation of motion for a small particle in Brownian motion (see [3, Chapter 7]). Consider a particle of mass m undergoing Brownian motion in a fluid. The particle is assumed to be larger and heavier than the constituents of the fluid. In addition, the particle is subjected to a damping force proportional to the particle's velocity with the coefficient γ , and a random force $\xi(t)$ due to the thermal agitation of molecules composing the fluid. The equation of motion according to Newton's 2nd law is

$$m \frac{dv(t)}{dt} = -\gamma v(t) + \xi(t), \tag{2}$$

which we write in one dimension for simplicity. The random force $\xi(t)$ is also called *white noise*, which is assumed to be a Gaussian process with stationary (invariant under time shifts) and Markov

properties. $\xi(t)$ also has the following averaged properties

$$\langle \xi(t) \rangle = 0, \quad \langle \xi(t)\xi(t') \rangle = g\delta(t-t'), \quad (3)$$

where δ is the delta Dirac function and g is the measure of the noise's strength. In the notation of this course, the equations (2) and (3) can be rewritten as

$$m dV_t = -\gamma V_t dt + \Xi_t dt, \quad (4)$$

$$\mathbb{E}(\Xi_t) = 0, \quad \mathbb{E}(\Xi_t \Xi_{t'}) = g\delta(t-t'), \quad (5)$$

where we replace v and ξ with their capital counterparts. Comparing (4) with (1), we get $\beta = \gamma/m$ for the first term and formally for the second term,

$$\Xi_t \text{ " = " } m\alpha \frac{dB_t}{dt}.$$

Since we know that the Brownian motion (B_t) is nowhere differentiable, this expression of Ξ_t is not well-defined. However, the version (4) of the Langevin equation is widely used in physics. Therefore, a heuristic explanation (see [4]) can be made to justify the usage of Ξ_t (or $\xi(t)$) and its properties (5). From the definition of the Brownian motion, we have for any $t, \delta t > 0$,

$$(B_{t+\delta t} - B_t) \sim N(0, \delta t) \quad \Rightarrow \quad \Xi_t^{\delta t} := m\alpha \frac{B_{t+\delta t} - B_t}{\delta t} \sim N\left(0, \frac{m^2\alpha^2}{\delta t}\right).$$

It is clear that $\mathbb{E}(\Xi_t^{\delta t}) = 0$ for any $t, \delta t > 0$. Since the Brownian motion B_t is also a martingale, we can use the property that the increments are uncorrelated to compute the autocovariance $\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t})$. Fix the value of t . For any t' such that $t' + \delta t < t$ or $t + \delta t < t'$, we have

$$\mathbb{E}[(B_{t+\delta t} - B_t)(B_{t'+\delta t} - B_{t'})] = 0 \quad \Rightarrow \quad \mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t}) = 0.$$

Now, consider t' such that $t' \leq t \leq t' + \delta t$. Then, we have $t' \leq t \leq t' + \delta t \leq t + \delta t$ and

$$\begin{aligned} \mathbb{E}[(B_{t+\delta t} - B_t)(B_{t'+\delta t} - B_{t'})] &= \mathbb{E}[(B_{t+\delta t} - B_{t'+\delta t} + B_{t'+\delta t} - B_t)(B_{t'+\delta t} - B_{t'} + B_t - B_{t'})] \\ &= \mathbb{E}[(B_{t+\delta t} - B_{t'+\delta t})(B_{t'+\delta t} - B_t)] + \mathbb{E}[(B_{t'+\delta t} - B_t)^2] \\ &\quad + \mathbb{E}[(B_{t+\delta t} - B_{t'+\delta t})(B_t - B_{t'})] + \mathbb{E}[(B_{t'+\delta t} - B_t)(B_t - B_{t'})] \\ &= 0 + (t' + \delta t - t) + 0 + 0 \\ &= \delta t - (t - t'), \end{aligned}$$

where we use the fact that $(B_{t'+\delta t} - B_t) \sim N(0, t' + \delta t - t)$. For the case $t \leq t' \leq t + \delta t$, since we also have $t \leq t' \leq t + \delta t \leq t' + \delta t$, we can use the above result and interchange t and t' :

$$\mathbb{E}[(B_{t+\delta t} - B_t)(B_{t'+\delta t} - B_{t'})] = \delta t + (t - t').$$

As a result, we get the autocovariance depending only on $(t - t')$,

$$\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t}) = \begin{cases} m^2\alpha^2 \frac{\delta t - |t-t'|}{\delta t^2}, & |t-t'| \leq \delta t, \\ 0, & |t-t'| > \delta t. \end{cases}$$

In addition, we have

$$\begin{aligned} \int_{t-\delta t}^{t+\delta t} (\delta t - |t - t'|) dt' &= \int_{t-\delta t}^t (\delta t - t + t') dt' + \int_t^{t+\delta t} (\delta t + t - t') dt' \\ &= \left(\delta t^2 - t\delta t + \frac{2t\delta t - \delta t^2}{2} \right) + \left(\delta t^2 + t\delta t - \frac{2t\delta t + \delta t^2}{2} \right) \\ &= \delta t^2. \end{aligned}$$

Therefore, we get

$$\int_{-\infty}^{\infty} \mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t}) dt' = m^2 \alpha^2 \frac{\delta t^2}{\delta t^2} = m^2 \alpha^2.$$

As δt approaches 0, the region, where $\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t})$ has nonzero values, shrinks while $\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t}) = m^2 \alpha^2 / \delta t \rightarrow \infty$ for $t = t'$. Meanwhile, the integral of $\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t})$ with respect to t or t' on \mathbb{R} remains a constant. Hence, $\mathbb{E}(\Xi_t^{\delta t} \Xi_{t'}^{\delta t})$ formally converges to $m^2 \alpha^2 \delta(t - t')$ as δt goes to 0. Comparing with (5), we get the equality $\alpha^2 = g/m^2$. Consequently, we can define formally

$$\Xi_t := \lim_{\delta t \rightarrow 0} \Xi_t^{\delta t} = \sqrt{g} \lim_{\delta t \rightarrow 0} \frac{B_{t+\delta t} - B_t}{\delta t} \quad “ = ” \quad \sqrt{g} \frac{dB_t}{dt},$$

and get all the properties required. In hindsight, we can rewrite (4) in a more rigorous form:

$$m dV_t = -\gamma V_t dt + \sqrt{g} dB_t. \tag{6}$$

Even though (4) is mostly used in physics, it is a mnemonic for (6), which is more mathematically precise.

References

- [1] Serge Richard, Lecture Notes: Introduction to Stochastic Calculus (Fall 2023)
- [2] J.-L. Arguin, A first course in stochastic calculus
- [3] L. E. Reichl, A modern course in statistical physics
- [4] Gillespie, Daniel T. “The Mathematics of Brownian Motion and Johnson Noise.” *American Journal of Physics* 64, no. 3 (March 1, 1996): 225–40. <https://doi.org/10.1119/1.18210>.