

Lyapunov and asymptotical stability

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By definition, a fixed point z of $f(x)$ is asymptotically stable if it is Lyapunov stable and there exists $B_\epsilon(z)$ such that $\forall x_0 \in B_\epsilon(z)$,

$$\lim_{t \rightarrow \infty} \|\phi(t, x_0) - z\| = 0. \quad (1)$$

However, the condition given by (1) does not imply Lyapunov stability, and this can be seen by considering the following system in \mathbb{R}^2 (from [T], p.200, Problem 6.16)

$$\begin{cases} \dot{x} = x - y - x(x^2 + y^2) + \frac{xy}{\sqrt{x^2 + y^2}}, \\ \dot{y} = x + y - y(x^2 + y^2) - \frac{x^2}{\sqrt{x^2 + y^2}}. \end{cases}$$

By changing to polar coordinates, i.e. for $r \in (0, \infty)$ and $\theta \in [0, 2\pi)$, we set

$$x = r \cos \theta, \quad y = r \sin \theta,$$

the differential equations can be rewritten as

$$\begin{cases} (r \dot{\cos} \theta) = r \cos \theta - r \sin \theta - r^3 \cos \theta + r \cos \theta \sin \theta, \\ (r \dot{\sin} \theta) = r \cos \theta + r \sin \theta - r^3 \sin \theta - r \cos^2 \theta. \end{cases}$$

Expand the left hand side of the equations

$$\begin{cases} \dot{r} \cos \theta - r \sin \theta \dot{\theta} = r \cos \theta - r \sin \theta - r^3 \cos \theta + r \cos \theta \sin \theta, \\ \dot{r} \sin \theta + r \cos \theta \dot{\theta} = r \cos \theta + r \sin \theta - r^3 \sin \theta - r \cos^2 \theta. \end{cases}$$

By multiplying the first equation by $\cos \theta$ and the second one by $\sin \theta$ then adding them together, we have

$$\begin{aligned} \dot{r}(\cos^2 \theta + \sin^2 \theta) &= r(\cos^2 \theta + \sin^2 \theta) - r^3(\cos^2 \theta + \sin^2 \theta) \\ \text{or } \dot{r} &= r - r^3 = r(1 - r)(1 + r). \end{aligned} \quad (2)$$

On the other hand, multiplying the first equation by $-\sin \theta$ and the second one by $\cos \theta$ then adding them together gives

$$r(\sin^2 \theta + \cos^2 \theta) \dot{\theta} = r(\sin^2 \theta + \cos^2 \theta) - r \cos \theta (\sin^2 \theta + \cos^2 \theta)$$

$$\text{or } \dot{\theta} = 1 - \cos \theta = 2 \sin^2 \frac{\theta}{2}. \quad (3)$$

The change to polar coordinates has, therefore, decoupled the original ODEs into two separate ODEs that can be solved independently. Notice that $(r, \theta) = (1, 0)$ - which corresponds to $(x, y) = (1, 0)$ - is a fixed point of the system, and let this point be denoted by z . Also, whenever the right hand sides of equations (2) and (3) are non-zero, i.e. $r \neq 1$ and $\theta \neq 0$, the equations are separable and can be solved implicitly by evaluating

$$\int_{r_0}^r \frac{dr}{r(1-r)(1+r)} = \int_0^t dt = t,$$

$$\int_{\theta_0}^{\theta} \frac{d\theta}{2 \sin^2(\theta/2)} = \int_0^t dt = t,$$

where $(r_0, \theta_0)_{r, \theta}$ is the initial condition given at $t = 0$. We can calculate each integral in the left hand side as follows

$$\begin{aligned} \int_{r_0}^r \frac{dr}{r(1-r)(1+r)} &= \int_{r_0}^r \left[\frac{1}{r} + \frac{1}{2(1-r)} - \frac{1}{2(1+r)} \right] dr \\ &= \int_{r_0}^r \frac{dr}{r} + \frac{1}{2} \int_{r_0}^r \frac{dr}{1-r} - \frac{1}{2} \int_{r_0}^r \frac{dr}{1+r} \\ &= \ln \left| \frac{r}{r_0} \right| - \frac{1}{2} \ln \left| \frac{1-r}{1-r_0} \right| - \frac{1}{2} \ln \left| \frac{1+r}{1+r_0} \right| \\ &= -\frac{1}{2} \ln \left| \left(\frac{r_0}{r} \right)^2 \cdot \frac{1-r}{1-r_0} \cdot \frac{1+r}{1+r_0} \right| \\ &= -\frac{1}{2} \ln \left| \frac{1/r^2 - 1}{1/r_0^2 - 1} \right|, \end{aligned}$$

$$\begin{aligned} \int_{\theta_0}^{\theta} \frac{d\theta}{2 \sin^2(\theta/2)} &= \int_{\theta_0}^{\theta} \frac{d(\theta/2)}{\sin^2(\theta/2)} \\ &= - \left(\cot \frac{\theta}{2} - \cot \frac{\theta_0}{2} \right). \end{aligned}$$

These then give

$$-\frac{1}{2} \ln \left| \frac{1/r^2 - 1}{1/r_0^2 - 1} \right| = t, \quad - \left(\cot \frac{\theta}{2} - \cot \frac{\theta_0}{2} \right) = t.$$

By rearranging the terms, it is possible to show (with appropriate conditions) that

$$r(t) = \left[1 + \left(\frac{1}{r_0^2} - 1 \right) e^{-2t} \right]^{-\frac{1}{2}},$$

$$\theta(t) = 2 \cot^{-1} \left(\cot \frac{\theta_0}{2} - t \right).$$

As $t \rightarrow \infty$, we have

$$\lim_{t \rightarrow \infty} r(t) = 1,$$

$$\lim_{t \rightarrow \infty} \theta(t) = 2\pi,$$

which are independent of the initial condition and also correspond to the fixed point z . Therefore, if we define $\phi(t, x_0) = (x(t), y(t))$ with initial condition $(x(0), y(0)) = x_0 \equiv (r_0, \theta_0)_{r, \theta}$ then

$$\lim_{t \rightarrow \infty} \phi(t, x_0) = z,$$

so condition (1) is satisfied. Now we consider the initial condition where $r_0 = 1$ and $\theta_0 \neq 0$, meaning that the system starts at somewhere along the unit circle except at the fixed point z , and since $\dot{r} = 0$, $\phi(t, x_0)$ will stay on the unit circle. By the cosine law, the distance from $\phi(t, x_0)$ to z is given by

$$\|\phi(t, x_0) - z\| = [2 - 2 \cos \theta(t)]^{\frac{1}{2}}.$$

Since $\dot{\theta} = 2 \sin^2(\theta/2) > 0$, $\theta(t)$ is strictly increasing. Therefore, in the time interval where $0 < \theta(t) < \pi$, $\|\phi(t, x_0) - z\|$ is strictly increasing, i.e. $\phi(t, x_0)$ is always moving further away from z . We have the distance from x_0 to z being

$$\|x_0 - z\| = [2 - 2 \cos \theta_0]^{\frac{1}{2}},$$

and because θ_0 can be chosen arbitrarily close to 0, x_0 can also be chosen arbitrarily close to z . Hence, there exists $\delta > 0$ such that $\forall \epsilon > 0$ and $\epsilon < \delta$, we can always find $t > 0$ and $x_0 \in B_\epsilon(z)$ so that

$$\phi(t, x_0) \notin B_\delta(z),$$

so z is not Lyapunov stable even when condition (1) is satisfied.

□