

VI General relativity

Def. a PSEUDO-RIEMANNIAN MANIFOLD is a pair (M, ϕ) with

M a smooth manifold and $\phi \in J^2(M)$, symmetric and non-degenerate.

⚠ No «positive definite» required!

$$\phi(X, Y) = \phi(Y, X) \quad ; \quad \phi(X, Y) = 0 \quad \forall Y \in X(M)$$

a LORENTZIAN MANIFOLD is a pseudo-Riemannian manifold $X=0$

with $(g_{ij}) = \begin{pmatrix} -1 & & & 0 \\ & 1 & & \\ & & \ddots & \\ 0 & & & 1 \end{pmatrix}$ in suitable coordinates (locally).
 \rightarrow signature (trace) = $n-2$

Facts for pseudo-Riemannian manifolds

1) unicity of Levi Civita connection when the 2 conditions are imposed.

2) Koszul formula still holds.

3) Hopf-Rinow thm + geodesically complete are **no more valid**.

\Rightarrow We don't have a metric space anymore.

4) Cartan structure thm are still valid.

Recall that the length of a vector is not changed under parallel transport along a curve.

Geodesics c satisfy that \dot{c} is parallel transported along c .

$$\Rightarrow \phi(\dot{c}, \dot{c}) = \text{cst}$$

Def. A geodesic c on a pseudo-Riemannian manifold (M, ϕ) is

TIMELIKE, NULL, or SPACELIKE if

$$\phi(\dot{c}, \dot{c}) < 0, \quad \phi(\dot{c}, \dot{c}) = 0 \quad \text{or} \quad \phi(\dot{c}, \dot{c}) > 0$$

< 0 and $= 0$ are allowed by PSEUDO METRIC ϕ

Remark: these expressions come from special relativity with $M = \mathbb{R}^4$ and

$$(g_{\nu\mu}) = \begin{pmatrix} -1 & & & 0 \\ & 1 & & \\ & & \ddots & \\ 0 & & & 1 \end{pmatrix} \quad \text{a special case of a Lorentzian manifold.}$$

$\mu, \nu = 0, 1, 2, 3$

For a Lorentzian manifold (M, ϕ) of dim 4, with the Levi Civita connection, the Einstein field equation reads

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu} \quad \text{for } \mu, \nu = 0, 1, 2, 3 \quad (*)$$

$R_{\mu\nu}$ Ricci curvature tensor
 $\frac{1}{2} R g_{\mu\nu}$ scalar curvature
 $\Lambda g_{\mu\nu}$ cosmological constant > 0
 $T_{\mu\nu}$ stress-energy or energy-momentum tensor
 G = gravitation constant
 c = speed of light

(about geometry) $G_{\mu\nu}$ Einstein tensor

contains the physics (energy + matter)

⚠ Not so much freedom for writing a meaningful equations.

This is a system of 10 equations because of symmetry between μ and ν .

In addition the thms of structure reduces the number of indep eq.

Remark: These equations define the pseudo metric tensor $g_{\mu\nu}$.

Indeed, $R_{\mu\nu\gamma}{}^\delta$ and $R_{\mu\nu}$ can be expressed in terms of $\Gamma_{\mu\nu}{}^\delta$ and its derivatives.

And $\Gamma_{\mu\nu}{}^\delta$ can be expressed in terms of $g_{\mu\nu}$ and its derivatives.

$\Rightarrow \textcircled{*}$ is a system of non linear partial differential equations for $g_{\mu\nu}$.

Schwarzschild solution

Assumptions: $\circ T_{\mu\nu} = 0$

$\circ g_{\mu\nu}$ is time independent (static solution)

\circ spherically symmetric in space (\equiv in the indices 1, 2, 3)

$\circ M = \mathbb{R} \times \mathbb{R}_+ \times \mathbb{S}^2$ \mathbb{R} for t and $\mathbb{R}_+ \times \mathbb{S}^2$ is \mathbb{R}^3 in spherical coordinates

Suppose that

$$g = -A^2(r) dt \otimes dt + B^2(r) dr \otimes dr + r^2 d\theta \otimes d\theta + r^2 \sin^2(\theta) d\varphi \otimes d\varphi \in J^2(M)$$

with $A, B: \mathbb{R}_+ \mapsto \mathbb{R}$ unknown,

$(dt, dr, d\theta, d\varphi) \in J^1(M)$ generate a basis of $T^*(M)$.

$\left\{ \left(\frac{\partial}{\partial x^i} \right)_p \right\}_{i=1}^n$ is a basis of $T_p(M)$, and $\{dx_p^i\}_{i=1}^n$ is a basis of $T_p^*(M)$.

$\Rightarrow \{dx^i \otimes dx^j\}_{i,j}$ is a basis of $J^2(M)$.

Question: can we find A, B such that $\textcircled{*}$ is satisfied (with $T_{\mu\nu} = 0$)?

Two approaches:

1) Express $\Gamma_{\mu\nu}{}^\delta \rightsquigarrow R_{\mu\nu\sigma}{}^\rho$ and then $R_{\mu\nu}$ and R in terms of $g_{\mu\nu}$, and solve $\textcircled{*}$

2) Set $\theta^0 := A(r) dt$ $\theta^2 := r d\theta$
 $\theta^1 := B(r) dr$ $\theta^3 := r \sin(\theta) d\varphi$ $\} \in J^1(M)$ and observe that

$$g = \sum_{\mu,\nu} \eta_{\mu,\nu} \theta^\mu \otimes \theta^\nu \text{ and } \eta_{\mu\nu} = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \text{ and that}$$

$\{\theta^0, \theta^1, \theta^2, \theta^3\}$ is an orthonormal coframe a basis of $T^*(M)$

• Define $\theta_\mu{}^\nu$ and $\Omega_\mu{}^\nu$ (connection and curvature tensors) and write the structure relations of Cartan.

• One obtains some differential equations for A and B , which can be solved.

• $A(r) = \left(1 - \frac{2m}{r}\right)^{\frac{1}{2}}$ and $B(r) = \left(1 - \frac{2m}{r}\right)^{-\frac{1}{2}}$ with $m \in \mathbb{R}$ an integration const

Conclusion

Textbooks on general relativity are now accessible

(but still the theory is complicated).