# Exact WKB analysis and cluster algebras 

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The pdf file of this slide (or updated one) will be available at my web site.

## exact WKB analysis and cluster algebras

WKB approximation
Wentzel, Kramers, Brilloin (1926)
semiclassical approximation method for Schrödinger equation
exact WKB analysis ( $80 \sim$ )
$\fallingdotseq$ study of WKB solution of 1d (complex) Schrödinger equation by Borel resummation
Voros (83)
Aoki-Kawai-Takei (91)
Dellabaere-Dillinger-Pham (DDP) (93)
cluster algebras
Fomin-Zelevinsky (00 ~)
combinatorial structure in representation theory in several contexts
appearing in several areas in mathematics

## Overview \& Keywords

Hitchin
system


3-CY
categories
[Kontsevich08]
differentials (Stokes graphs)

| Schrödinger | $[$ Kawai05] |
| :--- | :--- |
| equation | [Gekhtman05,Fock03,Fomin08] |

cluster algebras
"pentagon relation" [DDP93]

$$
\mathfrak{S}_{\gamma_{1}} \mathfrak{S}_{\gamma_{2}}=\mathfrak{S}_{\gamma_{2}} \mathfrak{S}_{\gamma_{1}+\gamma_{2}} \mathfrak{S}_{\gamma_{1}}, \quad \mathfrak{S}_{\gamma}: \text { Stokes automorphism for cycle } \gamma
$$

"There is a striking similarity between our [their] wall-crossing formula and identities for the Stokes automorphisms in the theory of WKB asymptotics..." [Kontsevich-Soibelman08]
$\left\{\begin{array}{l}\text { Voros symbols (2) } \\ \text { Stokes phenomenon }\end{array}+\left\{\begin{array}{l}\text { Stokes graph (1) } \\ \text { flip }\end{array} \quad \longrightarrow \quad\left\{\begin{array}{l}\text { seed }\left\{\begin{array}{l}\text { quiver (1) } \\ \text { cluster variables (2) } \\ \text { mutation }\end{array}\right.\end{array}\right.\right.\right.$

## WKB solutions (1)

T. Kawai and Y. Takei, Algebraic analysis of singular perturbation theory, AMS, 2005. Schrödinger equation

$$
\left(\frac{d^{2}}{d z^{2}}-\eta^{2} Q(z, \eta)\right) \psi(z, \eta)=0
$$

$z$ : complex (local) coordinate, $\quad \eta=\hbar^{-1}$ : large parameter

$$
\begin{gathered}
\psi(z, \eta)=\exp \left(\int^{z} S(z, \eta) d z\right) \\
\frac{d S}{d z}+S^{2}=\eta^{2} Q \quad \text { (Riccati equation) } \\
\left\{\begin{array}{l}
Q(z, \eta)=Q_{0}(z)+\eta^{-1} Q_{1}(z)+\cdots \\
S(z, \eta)=\eta S_{-1}(z)+S_{0}(z)+\cdots \\
S_{-1}^{2}=Q_{0}, \quad \frac{d S_{-1}}{d z}+2 S_{-1} S_{0}=Q_{1} \\
S_{ \pm}(z, \eta) \\
= \pm \eta \sqrt{Q_{0}(z)}+\cdots \\
=S_{\text {even }}(z, \eta) \pm S_{\text {odd }}(z, \eta) \\
S_{\text {odd }}(z, \eta)=\eta \sqrt{Q_{0}(z)}+\cdots \\
S_{\text {even }}(z, \eta)=-\frac{1}{2} \frac{d}{d z} \log S_{\text {odd }}(z, \eta)
\end{array}\right.
\end{gathered}
$$

## WKB solutions (2)

(in previous page)

$$
\begin{aligned}
\psi(z, \eta) & =\exp \left(\int^{z} S(z, \eta) d z\right) \\
S_{ \pm}(z, \eta) & = \pm \eta \sqrt{Q_{0}(z)}+\cdots \\
& =S_{\text {even }}(z, \eta) \pm S_{\text {odd }}(z, \eta) \\
S_{\text {odd }}(z, \eta) & =\eta \sqrt{Q_{0}(z)}+\cdots \\
S_{\text {even }}(z, \eta) & =-\frac{1}{2} \frac{d}{d z} \log S_{\text {odd }}(z, \eta)
\end{aligned}
$$

Hence

$$
\begin{aligned}
& \psi_{ \pm}(z, \eta)= \frac{1}{\sqrt{S_{\text {odd }}(z, \eta)}} \exp \left( \pm \int^{z} S_{\text {odd }}(z, \eta) d z\right) \quad \text { WKB solutions } \\
&= \frac{1}{\sqrt{\eta \sqrt{Q_{0}(z)}}} \exp \left( \pm \int^{z} \eta \sqrt{Q_{0}(z)} d z\right)\left(1+O\left(\eta^{-1}\right)\right) \\
& \quad \text { WKB approximation } \quad \text { divergent series! }
\end{aligned}
$$

The exact WKB analysis manages this divergent series by Borel resummation.

## Borel resummation

$$
\begin{aligned}
f(\eta)= & \sum_{n=0}^{\infty} f_{n} \eta^{-n} \quad \text { (possibly divergent) formal series } \\
f_{B}(y)= & \sum_{n=1}^{\infty} \frac{f_{n}}{(n-1)!} y^{n-1} \quad \text { Borel transform of } f \\
\mathcal{S}[f](\eta)= & f_{0}+\int_{0}^{\infty} e^{-\eta y} f_{B}(y) d y \quad \text { Borel sum of } f \\
& \text { (not necessarily convergent) }
\end{aligned}
$$

## Example.

(1) $f(\eta)=\eta^{-n} \quad \Longrightarrow \quad \mathcal{S}[f](\eta)=\eta^{-n}$
(2) $f(\eta)$ : a convergent series of $\eta^{-1} \Longrightarrow \mathcal{S}[f](\eta)=f(\eta)$ near $\eta=\infty$.

## Borel summability

Definition. A formal series $f(\eta)=\sum_{n=0}^{\infty} f_{n} \eta^{-n}$ is Borel summable if

- $f_{B}(y)$ is a convergent series of $y$.
- $f_{B}(y)$ is analytically continued in the domain $\Omega$.
- $\left|f_{B}(y)\right| \leq c_{1} e^{c_{2}|y|}$ for some $c_{1}, c_{2}>0$.


Example. The following $f(\eta)$ is divergent, but Borel summable.

$$
\begin{aligned}
f(\eta) & =\sum_{n=0}^{\infty}(-1)^{n-1}(n-1)!\eta^{-n} \\
f_{B}(y) & =\sum_{n=1}^{\infty}(-1)^{n-1} y^{n-1} \quad\left(=\frac{1}{1+y} \quad \text { for }|y|<1\right)
\end{aligned}
$$

Theorem. (e.g. [Costin08]) If $f(\eta)$ is Borel summable, then the Borel sum $\mathcal{S}[f](\eta)$ converges near $\eta=\infty$; moreover, it is asymptotically expanded to $f(\eta)$,

$$
\mathcal{S}[f](\eta) \sim f(\eta) \quad(\eta \rightarrow \infty)
$$

## Stokes phenomenon

Formulation of Stokes phenomenon by Borel resummation

$\times$ : isolated singularity of $f_{B}(y)$ $(f(\eta)$ is not Borel summable)

| analytic function | $\mathcal{S}_{+\delta}[f]$ | $=$ | $\mathcal{S}_{-\delta}[f]$ | + |
| :---: | :---: | :---: | :---: | :---: |
| asymptotic expansion | $f$ |  | $f$ | + |


$g$ jump
(Stokes phenomenon)

This method will be applied to the WKB solutions of the Schrödinger equation

$$
\begin{gathered}
\left(\frac{d^{2}}{d z^{2}}-\eta^{2} Q(z, \eta)\right) \psi(z, \eta)=0 \\
\psi_{ \pm}(z, \eta)=\frac{1}{\sqrt{S_{\text {odd }}(z, \eta)}} \exp \left( \pm \int^{z} S_{\text {odd }}(z, \eta) d z\right)
\end{gathered}
$$

$S^{1}$-action on $y \Longrightarrow S^{1}$-action on $\eta \Longrightarrow S^{1}$-action on $z$, or $Q(z, \eta)$
(various viewpoints for Stokes phenomenon)

## Trajectories and Stokes curves

Schrödinger equation on a compact Riemann surface $\Sigma$

$$
\left(\frac{d^{2}}{d z^{2}}-\eta^{2} Q(z, \eta)\right) \psi(z, \eta)=0
$$

In this section, concentrate on the classic situation

$$
\Sigma=\text { Riemann sphere }, \quad Q(z, \eta)=Q_{0}(z) \text { polynomial in } z .
$$

Assume that $Q_{0}(z)$ has only simple zeros.
$Q_{0}(z)$ determines a foliation in $\Sigma$ :

$$
\text { leaf (trajectory) } \operatorname{Im} \int_{a}^{z} \sqrt{Q_{0}(z)} d z=\text { const } \quad\left(a \text { : zero of } Q_{0}(z)\right)
$$

$$
\text { Stokes curve } \quad \operatorname{Im} \int_{a}^{z} \sqrt{Q_{0}(z)} d z=0 \quad\left(a \text { : zero of } Q_{0}(z)\right)
$$

Around a simple zero $\times$

## Stokes graph

Stokes graph: graph on $\Sigma$ with vertices = zeros \& pole $(\infty)$, edges $=$ Stokes curves Stokes region: domain on $\Sigma$ surrounded by Stokes curves
Example: $Q_{0}(z)=i(z-1)(z+1)(z-2 i+2)(z-i)(z+i)$


Assume that the Stokes graph is saddle-free (i.e., without saddle trajectory
 Then, Stokes regions fall into the following two classes [Strebel84].

(regular) horizontal strip

half plane

## From Stokes graph to triangulation

To each saddle-free Stokes graph, one can associate a triangulation of a polygon.
Example: $Q_{0}(z)=i(z-1)(z+1)(z-2 i+2)(z-i)(z+i)$


| Stokes graph | triangulation |
| :---: | :---: |
| horizontal strip | (internal) arc |
| half plane | (boundary) edge |
| simple zero | triangle |

## Voros' connection formula

Theorem. [Voros83, Koike-Schäfke (to appear)]
Assume that the Stokes graph is saddle-free.
(1) The WKB solutions $\psi_{ \pm}(z, \eta)$ are Borel summable in each Stokes region $D$, so that the Borel sums $\mathcal{S}\left[\psi_{ \pm}\right](z, \eta)$ for $|\eta| \gg 1$ define analytic functions of $z$ in each $D$.
(2) The following connection formula holds:

$$
\begin{aligned}
& \mathcal{S}\left[\psi_{+}^{D_{1}}\right]=\mathcal{S}\left[\psi_{+}^{D_{2}}\right]+i \mathcal{S}\left[\psi_{-}^{D_{2}}\right] \\
& \mathcal{S}\left[\psi_{-}^{D_{1}}\right]=\mathcal{S}\left[\psi_{-}^{D_{2}}\right]
\end{aligned}
$$

where

and we normalize the WKB solutions at $a$, i.e.,

$$
\psi_{ \pm}(z, \eta)=\frac{1}{\sqrt{S_{\mathrm{odd}}(z, \eta)}} \exp \left( \pm \int_{a}^{z} S_{\mathrm{odd}}(z, \eta) d z\right)
$$

## DDP's jump formula

Assume that the Stokes graph has the following unique saddle trajectory:

$\gamma_{0}$ : saddle class

For each path $\beta \in H_{1}\left(\hat{\Sigma} \backslash \hat{P}_{0}, \hat{P}_{\infty}\right)$, we define

$$
\begin{aligned}
& W_{\beta}(\eta):=\int_{\beta}\left(S_{\text {odd }}(z, \eta)-\eta \sqrt{Q_{0}(z)}\right) d z \quad \text { Voros coefficient } \\
& e^{W_{\beta}} \quad \text { Voros symbol }
\end{aligned}
$$

Let $\left\langle\gamma_{0}, \beta\right\rangle$ be the intersection number of $\gamma_{0}$ and $\beta$.
Theorem. [Dellabaere-Dillinger-Pham93]
(1) If $\left\langle\gamma_{0}, \beta\right\rangle=0$, then $e^{W_{\beta}}$ is Borel summable.
(2) If $\left\langle\gamma_{0}, \beta\right\rangle \neq 0$, then $e^{W_{\beta}}$ is not Borel summable; the following jump formula holds:

$$
\mathcal{S}_{-\delta}\left[e^{W_{\beta}}\right]=\mathcal{S}_{+\delta}\left[e^{W_{\beta}}\right]\left(1+\mathcal{S}_{+\delta}\left[e^{V_{\gamma_{0}}}\right]\right)^{-\left\langle\gamma_{0}, \beta\right\rangle}, \quad V_{\gamma_{0}}(\eta):=\int_{\gamma_{0}} S_{\mathrm{odd}}(z, \eta) d z
$$

## Quadratic differential

Schrödinger equation on a surface $\Sigma$ :

$$
\left(\frac{d^{2}}{d z^{2}}-\eta^{2} Q(z, \eta)\right) \psi(z, \eta)=0
$$

In this section, we consider a more general situation than the classic one:
$\Sigma$ : compact Riemann surface,

$$
Q(z, \eta)=Q_{0}(z)+Q_{1}(z) \eta^{-1}+\cdots+Q_{k}(z) \eta^{-k}, \quad Q_{n}(z): \text { meromorphic }
$$

Under the coordinate transformation, the leading term $Q_{0}(z)$ transforms as a quadratic differential (e.g. [Kawai-Takei05]),

$$
Q_{0}(z) d z^{\otimes 2}
$$

$\Longrightarrow \sqrt{Q_{0}(z)} d z:$ 1-form single-valued on the double covering $\hat{\Sigma}$ of $\Sigma$

Assume

- Every zero of $Q_{0}(z)$ is simple.
- Every pole of $Q_{0}(z)$ is not simple.

We also assume some technical condition on the poles of $Q_{n}(z)(n \geq 2)$.

## Stokes graphs

Patterns of foliations around a pole [Strebel84]

(a)

(b)

(c)

$m-2$ tangent lines
pole of order $m \geq 3$ ( $m=6$ )

The Stokes graph and the Stokes regions are defined as before.
Assume that the Stokes graph is saddle-free (i.e., without saddle trajectory). Then, Stokes regions fall into three classes [Strebel84].

regular horizontal strip

degenerate horizontal strip

half plane

## From Stokes graph to triangulation

(in the previous page)

regular horizontal strip

degenerate horizontal strip

half plane

Assume that the Stokes graph is saddle-free.
(a part of) Stokes graph

(a part of) triangulation


These triangulations fit the surface realization of cluster algebras.

## Mutation of Stokes graphs

Under a continuous deformation of the potential $Q(z, \eta)$, the Stokes graph may change its topology. (= mutation of Stokes graphs)

## Example.


$e^{+\pi i / 10}\left(1-z^{2}\right)$

$1-z^{2}$

$e^{-\pi i / 10}\left(1-z^{2}\right)$


## Mutation of Stokes graphs

Example. $\Sigma=\mathbb{P}^{1}, Q(z)=e^{\delta \pi i}(z-1)(z+1)(z-2 i+2)(z-i)(z+i)$

$\delta=0.5$

$\delta=0.25$

## Mutation of Stokes graphs

## Mutation of Stokes graphs

- Under a continuous deformation of the potential $Q(z, \eta)$, the Stokes graph may change its topology. (= mutation of Stokes graphs)
- When the mutation occurs, one or more saddle connections appear.
- If only one saddle connection simultaneously appears during the mutation, it locally reduces to two types of elementary mutations called flip and pop [Gaiotto-Neitzke-Moore09], [Bridgeland-Smith13].
- They are refined to signed flip and signed pop.

Singed flip $\mu_{k}^{(\varepsilon)}$


It is induced from the $S^{1}$-action on $Q(z, \eta)$ for $G_{0}, Q^{(\theta)}(z, \eta):=e^{2 i \theta} Q\left(z, e^{i \theta} \eta\right)$.

## Simple paths and simple cycles

For simplicity, assume that the Stokes graph has no degenerate horizontal strip.
( $\leftrightarrow$ no self-folding triangle in the triangulation)
To each Stokes region $D_{i}$, we assign simple path $\beta_{i} \in H_{1}\left(\hat{\Sigma} \backslash \hat{P}_{0}, \hat{P}_{\infty}\right)$
simple cycle $\gamma_{i} \in H_{1}\left(\hat{\Sigma} \backslash \hat{P}_{0} \sqcup \hat{P}_{\infty}\right)$


$$
\text { duality }\left\langle\gamma_{i}, \beta_{j}\right\rangle=\delta_{i j}
$$

Under the signed flip $\mu_{k}^{(\varepsilon)}$, they mutate as

$$
\begin{aligned}
& \beta_{i}^{\prime}= \begin{cases}-\beta_{k}+\sum_{j=1}^{n}\left[-\varepsilon b_{j k}\right]+\beta_{j} & i=k \\
\beta_{i} & i \neq k\end{cases} \\
& \gamma_{i}^{\prime}= \begin{cases}-\gamma_{k} & g \text {-vector-like) } \\
\gamma_{i}+\left[\varepsilon b_{k i}\right]_{+} \gamma_{k} & i \neq k\end{cases} \\
& \text { (c-vector-like) }
\end{aligned}
$$

## Main Theorem: Mutation formula of Voros symbols

We set

$$
\begin{array}{lll}
x_{i}=e^{W_{i}}, & W_{i}=\int_{\beta_{i}}\left(S_{\text {odd }}(z, \eta)-\eta \sqrt{Q_{0}(z)}\right) d z & \text { Voros symbol for } \beta_{i}, \\
\hat{y}_{i}=e^{V_{i}}, & V_{i}=\int_{\gamma_{i}} S_{\text {odd }}(z, \eta) d z & \text { Voros symbol for } \gamma_{i}, \\
y_{i}=e^{v_{i}}, & v_{i}=\int_{\gamma_{i}} \eta \sqrt{Q_{0}(z)} d z, & \hat{y}_{i}=y_{i} \prod_{j=1}^{n} x_{j}^{b_{j i}}
\end{array}
$$

where we follow Fomin-Zelevinsky's notation for cluster algebras with coefficients.
Theorem. [IN14] Under the $S^{1}$-action $Q^{(\theta)}(z, \eta)=e^{2 i \theta} Q\left(z, e^{i \theta} \eta\right)$ which induces the signed mutation $\mu_{k}^{(\varepsilon)}$ of Stokes graphs, the Voros symbols "mutate" as

$$
\begin{aligned}
& y_{i}^{\prime} \rightsquigarrow \begin{cases}y_{k}^{-1} & i=k \\
y_{i} y_{k}\left[\varepsilon b_{k i}\right]_{+} & i \neq k\end{cases} \\
& x_{i}^{\prime} \rightsquigarrow \begin{cases}x_{k}^{-1}\left(\prod_{j=1}^{n} x_{j}^{\left[-\varepsilon b_{j k}\right]+}\right)\left(1+\hat{y}_{k}^{\varepsilon}\right) & i=k \\
x_{i} & i \neq k\end{cases} \\
& \hat{y}_{i}^{\prime} \rightsquigarrow \begin{cases}\hat{y}_{k}^{-1} & i=k \\
\hat{y}_{i} \hat{y}_{k}\end{cases} \\
& {\left[\varepsilon b_{k i}\right]+\left(1+\hat{y}_{k}^{\varepsilon}\right)^{-b_{k i}}} \\
& i \neq k .
\end{aligned}
$$

Remark. The jump terms come form DDP's jump formula.

