Quantum phase transition and Resurgence: Lessons from 3d N=4 SQED

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[arXiv: 2103.13654]

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Hiyoshi Phys. Keio U.[A], YITP[B], NCBJ[C], RECNS Keio[D], Akita->Kindai U.[E]

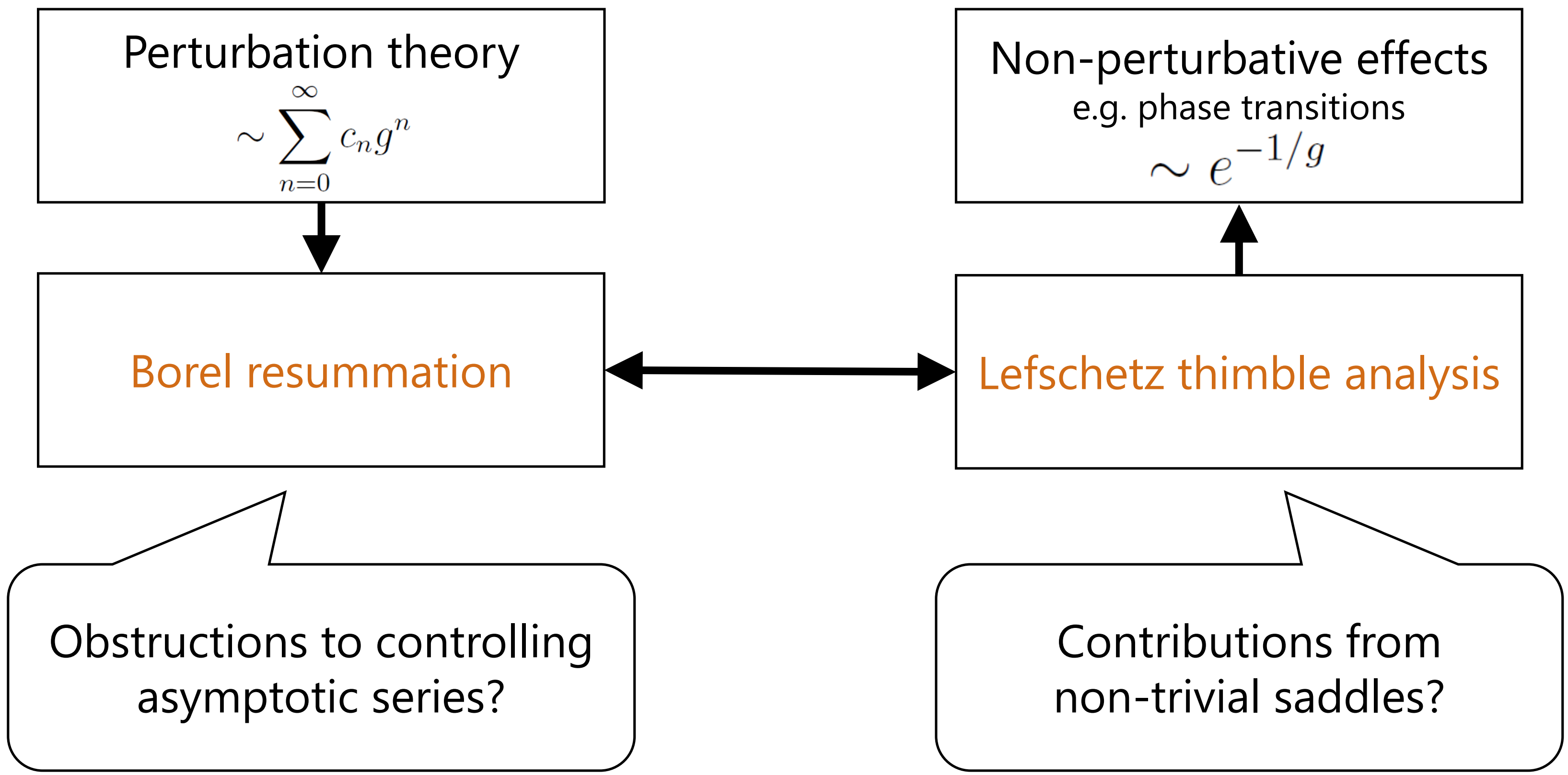
May./20/2021 Seminar @Nagoya University

Resurgence theory

[J. Ecalle, 81]

Lectures and reviews, e.g.
[M. Marino, 12]
[D. Dorigoni, 19]
[I. Aniceto, G. Basar, R. Schiappa, 19]

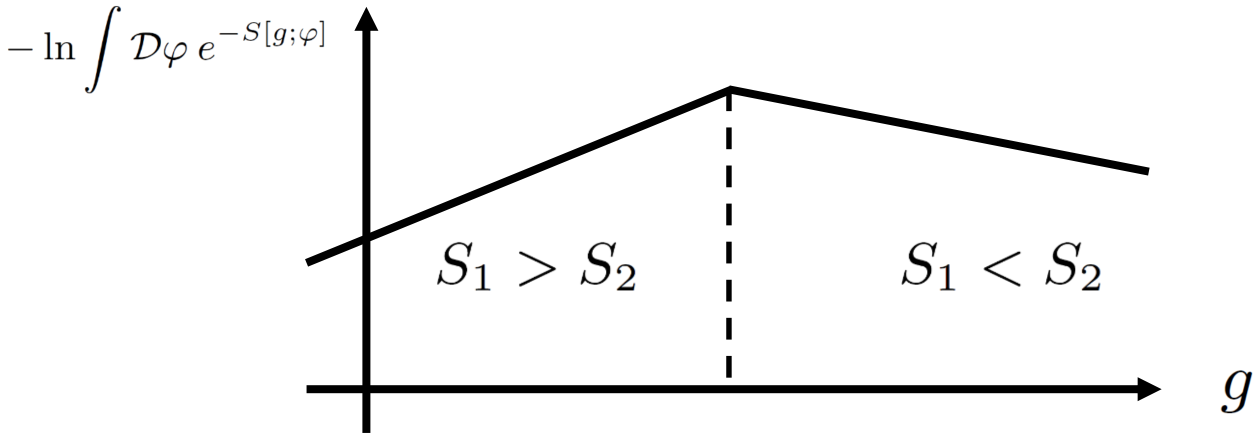
- One of the approaches to non-perturbative physics
- Decodes *non-perturbative information* from *perturbation theory*



Phase transitions and resurgence

Common story:

1st order phase transition = Change of dominant saddles



This is also within the scope of resurgence theory

Recent works:

- 0-dim Gross-Neveu, Nambu-Jona-Lasinio like model [T. Kanazawa, Y. Tanizaki, 15]
- 2dim Yang-Mills on lattice (reduced to Gross-Witten-Wadia) [G. Dunne et al., 16, 17, 18]

etc.

→ Is resurgence theory applicable to 2nd order phase transitions or more realistic QFTs?

Brief summary

Model

[Russo, Tierz, 17]

- $3\text{dim}\ \mathcal{N}=4\ U(1)$ SUSY gauge theory + $2N_f$ hypermultiplets with charge 1
- Fayet-Illiopoulos parameter η, flavor mass m

$\longrightarrow$ 2nd order quantum phase transition at the large-flavor limit

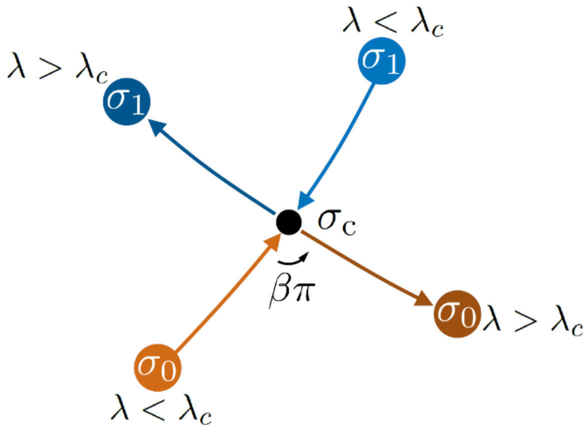
Result: resurgence is applicable!

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

2nd order phase transition = Simultaneous Stokes and Anti-Stokes phenomena

Change of the form of asymptotic expansion
and change of dominant saddles

- The order of the phase transition is determined by the collision angle of saddles
- Such information is encoded in a perturbative series



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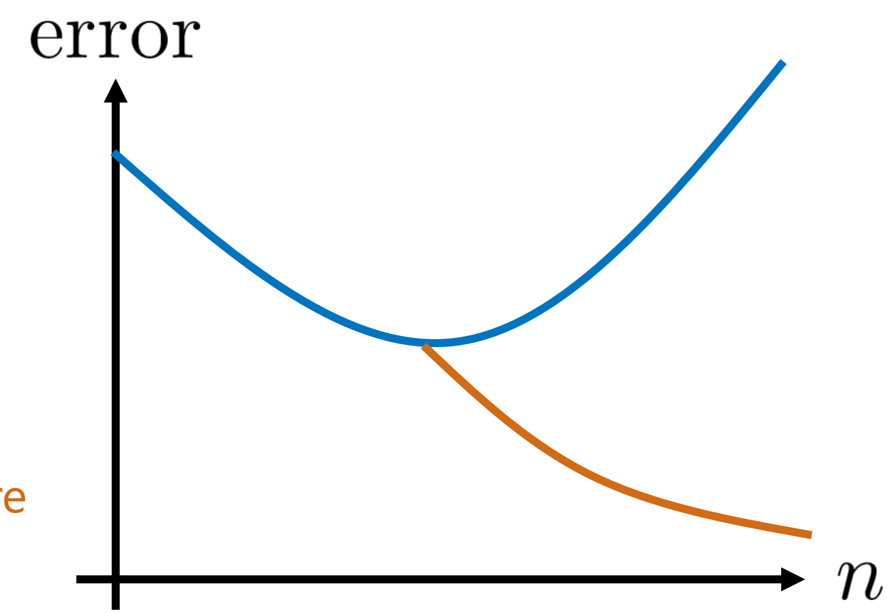
Before these, let me provide
lightning introduction to resurgence

Asymptotic series and resurgence structure

$$Z(g)=\sum_{n=0}^{\infty} c_n g^{n+r}, \quad c_n \sim n!$$



Resurgence structure

$$Z(g)=\boxed{\sum_{n=0}^{\infty} c_n^{(0)} g^n}+\boxed{\sum_{n=0}^{\infty} c_n^{(1)} g^n e^{-S_1/g}+\cdots}$$

[J. Ecalle, 81]
[M. Marino, 12]
[Cherman, Dorigoni, Unsal, 14]
[Cherman, Koroteev, Unsal, 14]
[D. Dorigoni, 19]
[I. Aniceto, G. Basar, R. Schiappa, 19]

Resurgence theory:

the perturbative part knows the non-perturbative parts

Example: 0dim Sine-Gordon model

Partition function:

$$Z(g) := \sqrt{\frac{\pi}{2g}}\, e^{-1/4g} I_0(1/4g) = \frac{1}{\sqrt{2\pi g}} \int_{-\pi/2}^{\pi/2} \mathrm{d}\varphi\, e^{-\frac{1}{2g}\sin^2\varphi} \qquad \left(-\frac{\pi}{2} < \pm\arg(1/4g) < \frac{3\pi}{2}\right)$$

$$= \underbrace{\sum_{n=0}^{\infty} \frac{(+2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1}}_{\text{Perturbative part}} \pm \underbrace{i e^{-1/2g} \sum_{n=0}^{\infty} \frac{(-2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1}}_{\text{Non-perturbative part}}$$

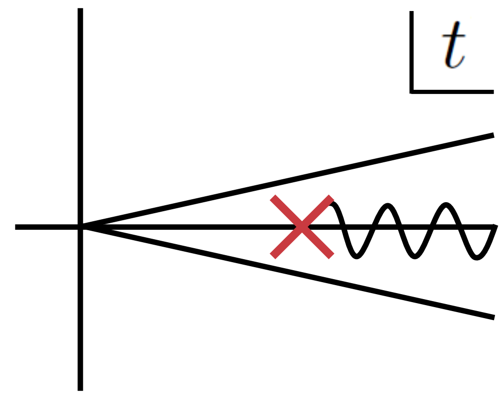
Observation 1

$$\sum_{n=0}^{\infty} c_n^{(0)} g^{n+1} \to \int_0^{\infty} \mathrm{d}t\, e^{-t/g} \sum_{n=0}^{\infty} \frac{c_n^{(0)}}{\Gamma(n+1)} t^n$$

Borel resummation

$$= \int_0^{\infty} \mathrm{d}t\, e^{-t/g}\, {}_2F_1\left(\frac{1}{2}, \frac{1}{2}, 1; 2t\right)$$

"Ambiguity" $\sim \boxed{ie^{-1/2g}}$



Recall

$$\sum_{n=0}^{\infty} \frac{(+2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1} \pm \boxed{ie^{-1/2g}} \sum_{n=0}^{\infty} \frac{(-2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1}$$

Observation 2

$$\underline{c_n^{(0)}} = \frac{(+2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)}$$

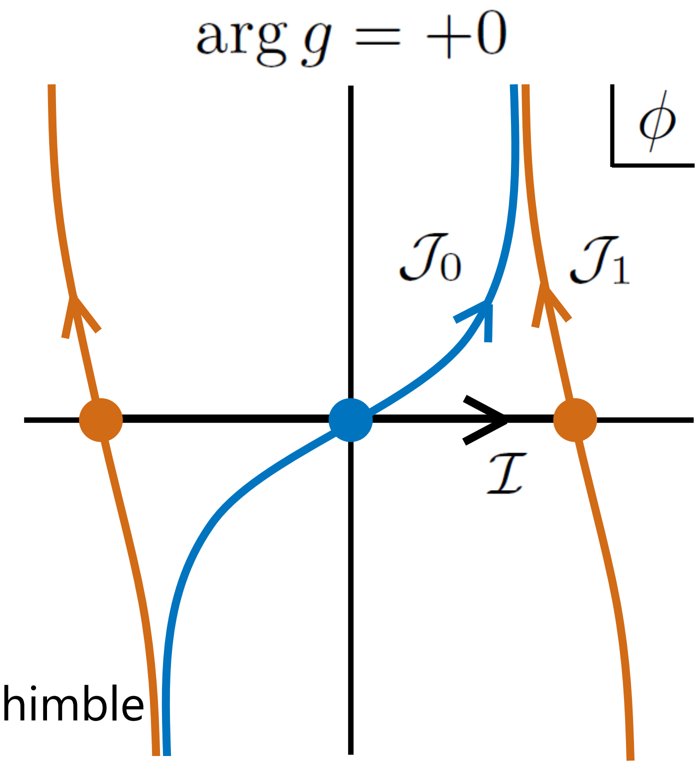
$$\sim \frac{2^n \Gamma(n)}{\Gamma(1/2)^2} \left[\boxed{1} + \frac{\boxed{-1/4}}{n-1} + \frac{\boxed{9/32}}{(n-1)(n-2)} + \frac{\boxed{-75/128}}{(n-1)(n-2)(n-3)} + \cdots \right]$$

$$\underline{c_n^{(1)}} = \frac{(-2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)}$$

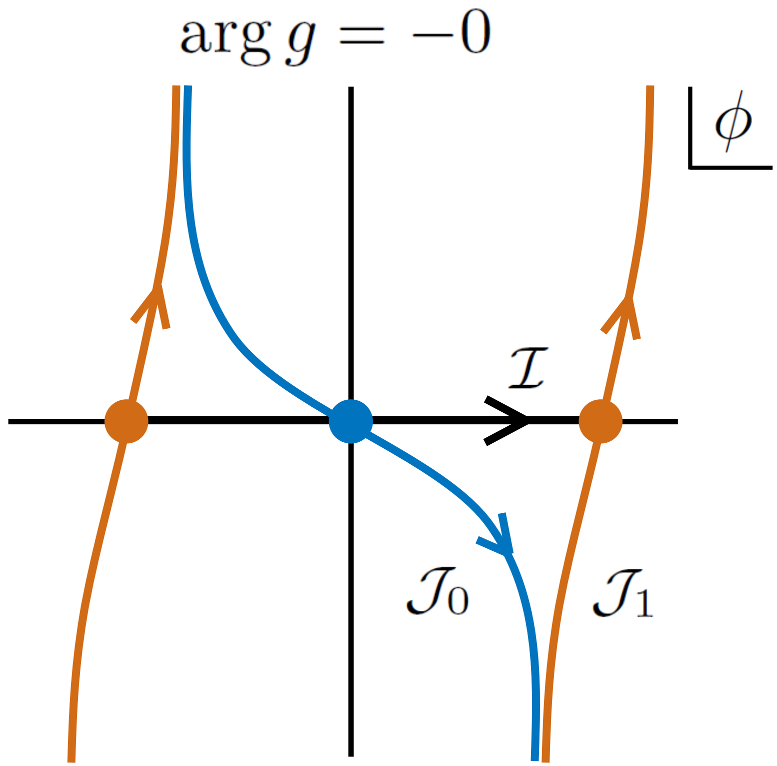
$$c_0^{(1)} = \boxed{1},\ c_1^{(1)} = 2^1 \left(\boxed{\frac{-1}{4}} \right),\ c_2^{(1)} = 2^2 \left(\boxed{\frac{9}{32}} \right),\ c_3^{(1)} = 2^3 \left(\boxed{\frac{-75}{128}} \right)$$

Observation 3

$$Z(g) = \frac{1}{\sqrt{2\pi g}} \int_{-\pi/2}^{\pi/2} \mathrm{d}\varphi \, e^{-\frac{1}{2g}\sin^2\varphi} = \underline{\sum_{n=0}^{\infty} c_n^{(0)} g^{n+1}} \boxed{\pm} \underline{ie^{-1/2g} \sum_{n=0}^{\infty} c_n^{(1)} g^{n+1}}$$



$$\mathcal{I} = \mathcal{J}_0 \boxed{-} \mathcal{J}_1$$



$$\mathcal{I} = \mathcal{J}_0 \boxed{+} \mathcal{J}_1$$

Resurgence theory

[J. Ecalle, 81]

Lectures and reviews, e.g.
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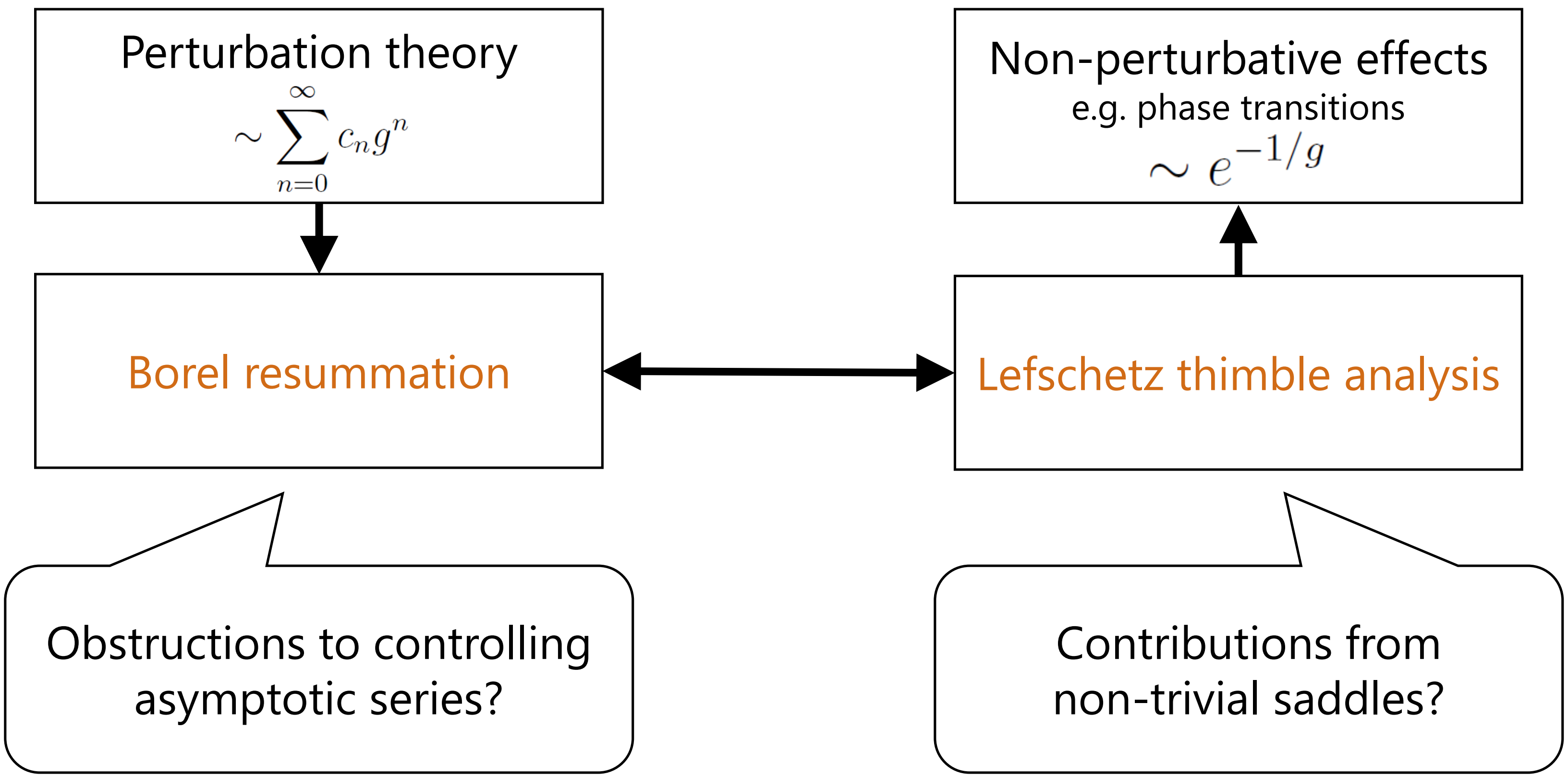
[D. Dorigoni, 19]
[I. Aniceto, G. Basar, R. Schiappa, 19]

- One of the approaches to non-perturbative physics
- Decodes *non-perturbative information* from *perturbation theory*



Contents

SQED3

Model:

$3\text{dim}\ \mathcal{N}=4\ U(1)$ SUSY gauge theory with

- $2N_f$ hypermultiplets (charge 1)
- Fayet-Illiopoulos parameter η
- flavor masses $\pm m$

Partition function:

Exactly computed on S^3 by SUSY localization technique

[Pestun, 12]
[A. Kapustin, B. Willett, I. Yaakov, 10]
[N. Hama, K. Hosomichi, S. Lee, 11]
[D. L. Jafferis, 12]

$$Z=\int_{-\infty}^{\infty} d\sigma \frac{e^{i\eta\sigma}}{\left[2\cosh\frac{\sigma+m}{2}\cdot 2\cosh\frac{\sigma-m}{2}\right]^{N_f}}$$

σ : Coulomb branch parameter
i.e. constant configuration of the scalar belonging to the vector multiplet in $\mathcal{N}=2$ Language

Saddles

't Hooft like limit:

[Russo, Tierz, 17]

$$N_f \to \infty, \quad \lambda \equiv \frac{\eta}{N_f} = \text{fixed}.$$

"Action":

$$S(\sigma) = N_f \Big[-i\lambda\sigma + \ln(\cosh\sigma + \cosh m) \Big]$$

Saddles:

$$\sigma_n^{\pm} = \log\left(\frac{-\lambda\cosh m \pm i\Delta(\lambda, m)}{i+\lambda} \right) + 2\pi i n \quad (n \in \mathbb{Z}),$$

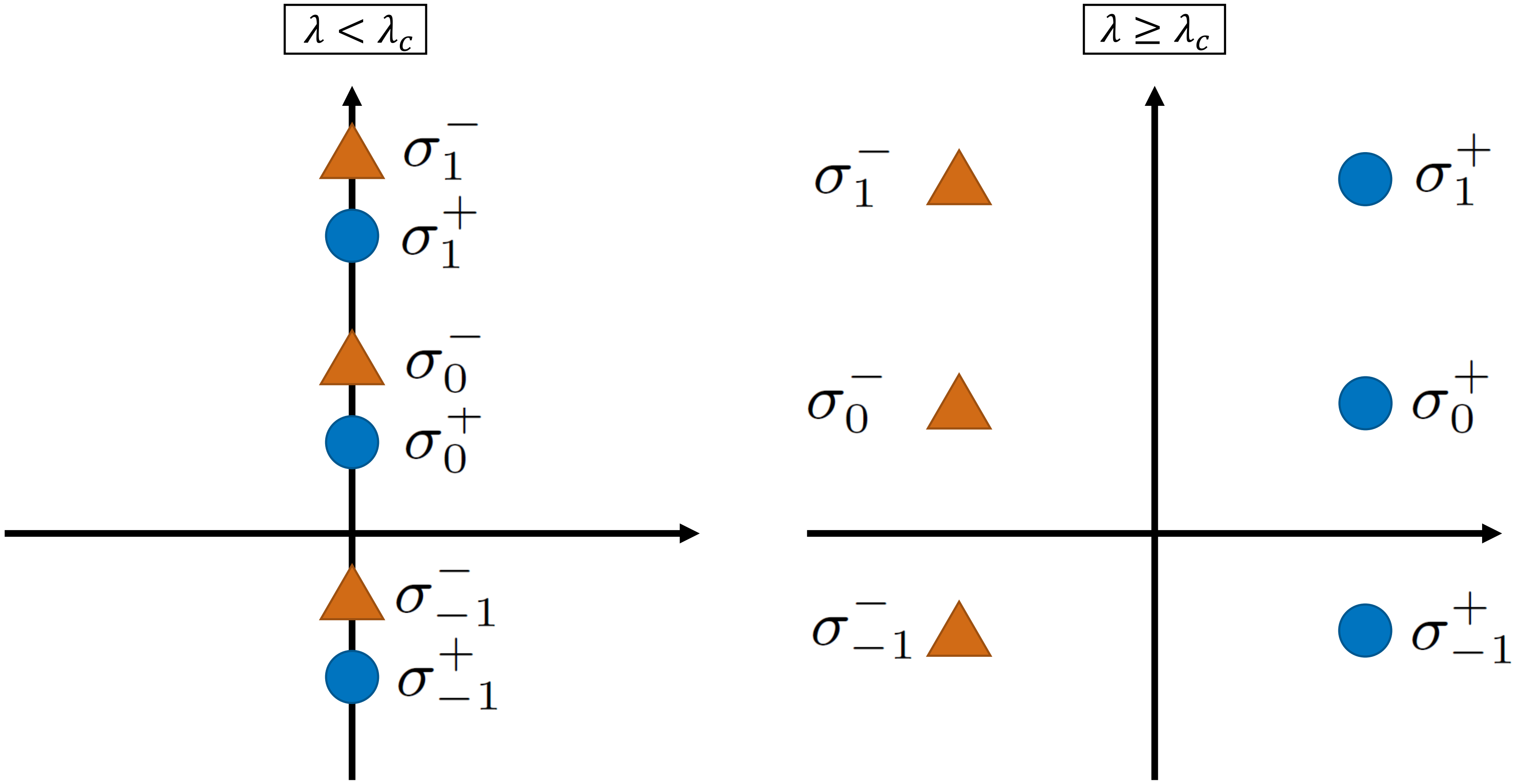
$$\Delta(\lambda, m) = \sqrt{1 - \lambda^2 \sinh^2 m}.$$

Something must happen at

$$\lambda_c \equiv \frac{1}{\sinh m}$$

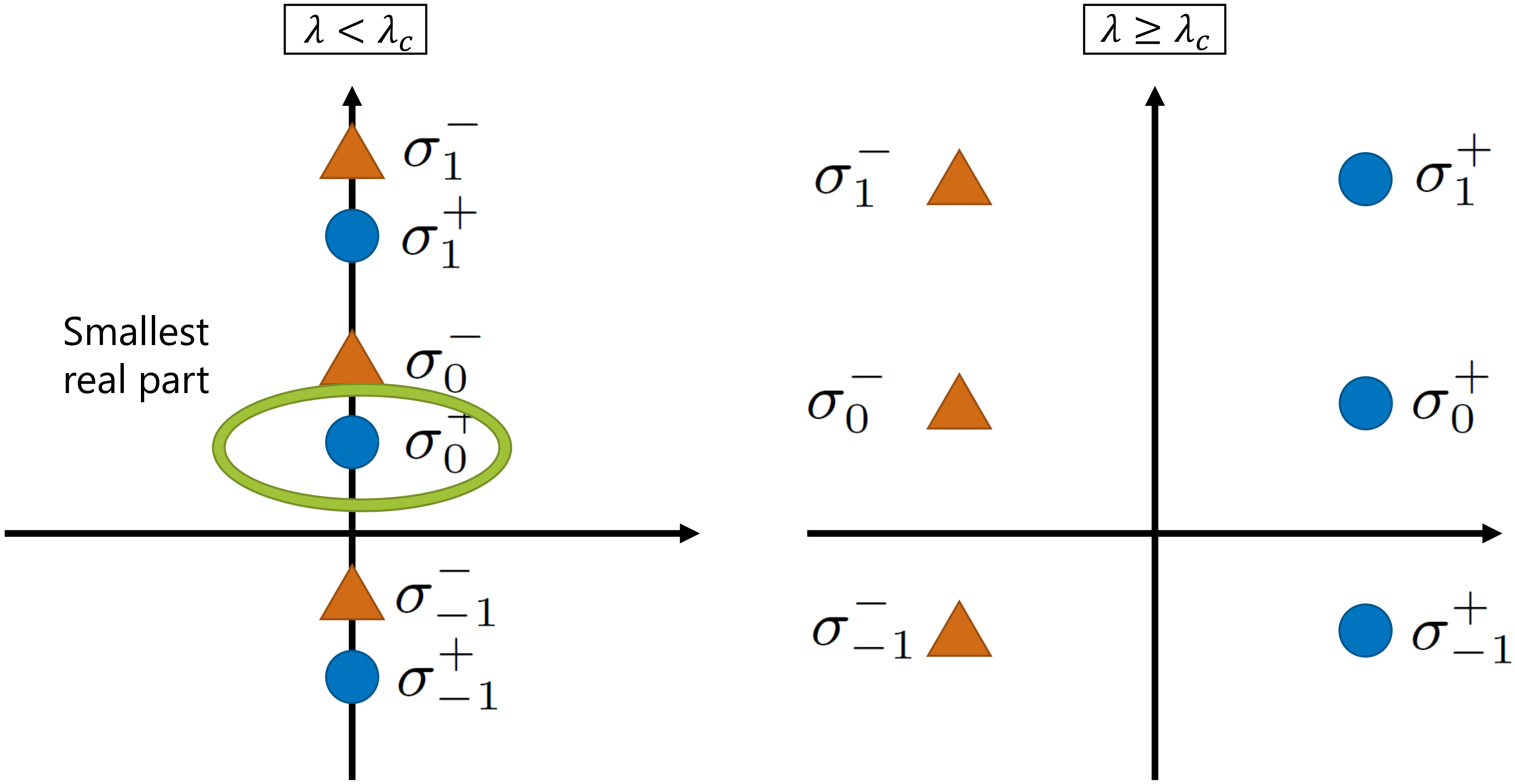
Saddles

[Russo, Tierz, 17]



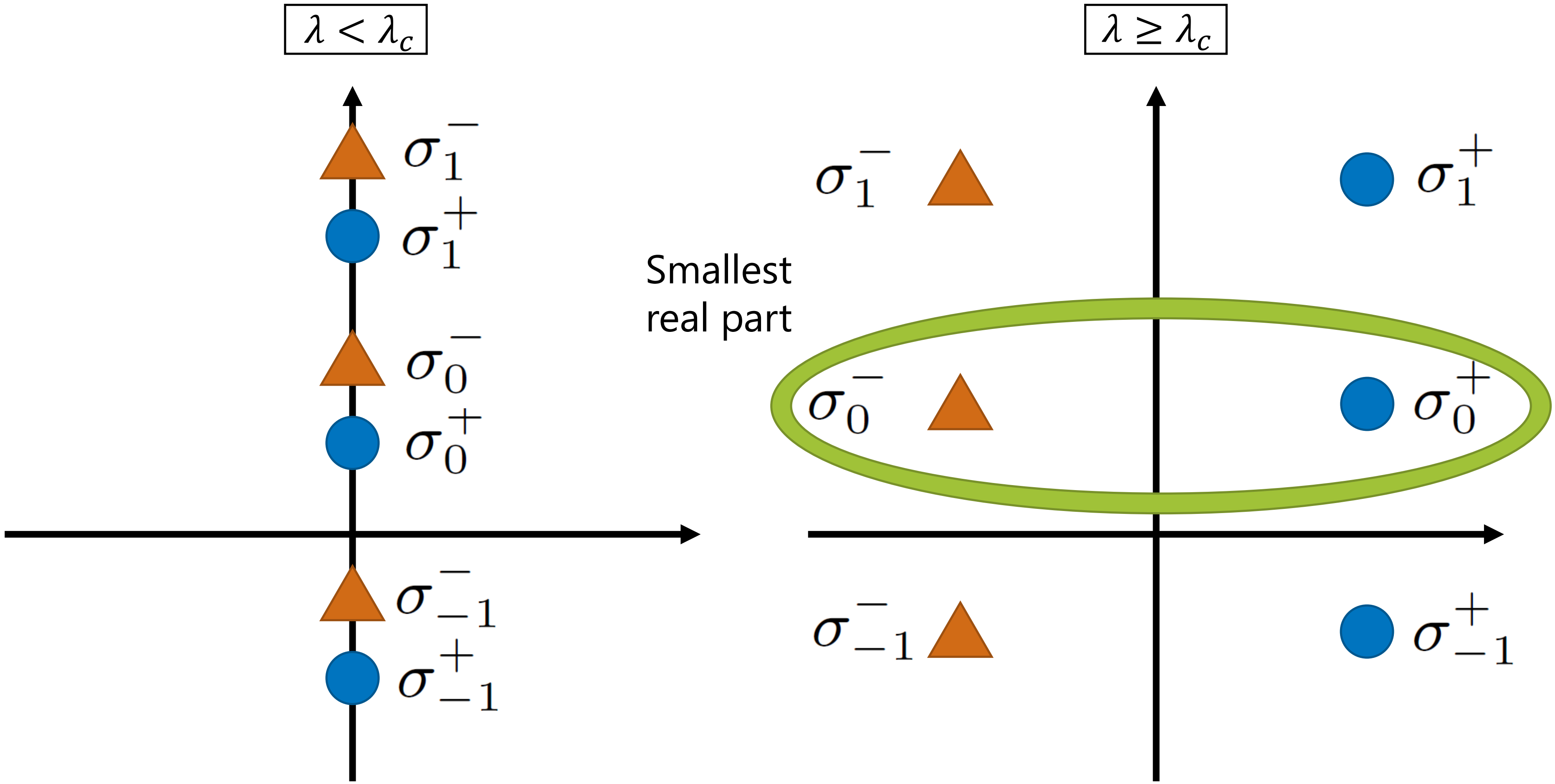
Saddles

[Russo, Tierz, 17]



Saddles

[Russo, Tierz, 17]



2^{nd} order phase transition

If the saddles of smallest real part contribute, [Russo, Tierz, 17]

$$\frac{d^2 F}{d\lambda^2} = \begin{cases} \frac{N_f}{1+\lambda^2}\left(1 + \frac{\cosh m}{\sqrt{1-\lambda^2 \sinh^2 m}}\right) & \lambda < \lambda_c \\ \frac{N_f}{1+\lambda^2} & \lambda \geq \lambda_c \end{cases}$$

$\longrightarrow$ 2^{nd} order phase transition

Questions:

- Smallest real part does not necessarily mean that such saddles contribute to the path integral. Can we ***justify it in more precise way***?
- Can we interpret the 2^{nd} order phase transition ***from the viewpoint of resurgence***, and ***draw lessons*** for generic QFTs?

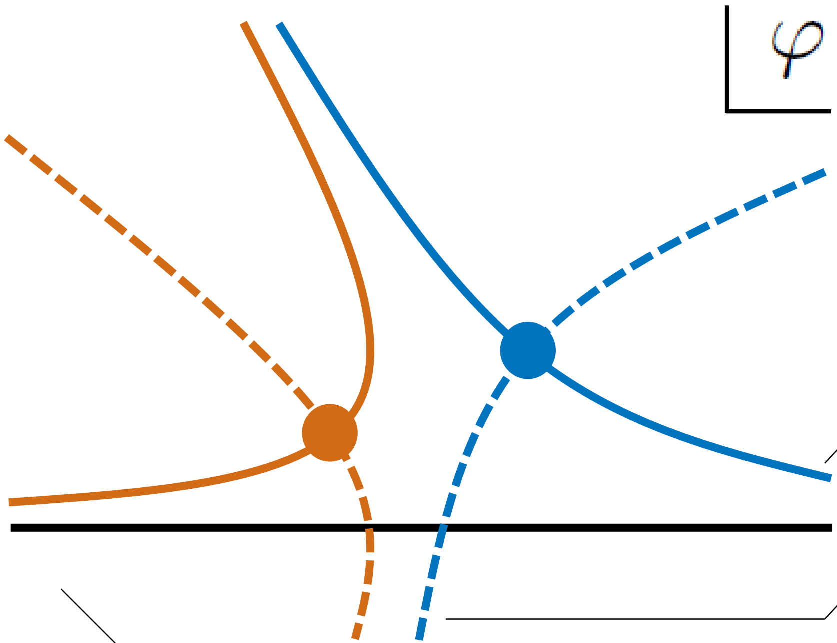
We will answer to these questions, Yes!

Contents

Lefschetz thimbles and dual-thimbles

Lefschetz thimbles = "Steepest descents" in configuration space
Dual-thimble = "Steepest ascents" in configuration space

[E. Witten, 11]
[M. Cristoforetti et al., 12,13,14]
[H. Fujii et al., 13]
[G. Aarts, 13]
[A. Alexandru et al, 16]



$$\mathcal{J}_i : \frac{\mathrm{d}\varphi(t)}{\mathrm{d}t} = +\overline{\frac{\delta S[\varphi]}{\delta\varphi}}, \quad \varphi(-\infty) = \varphi_i$$

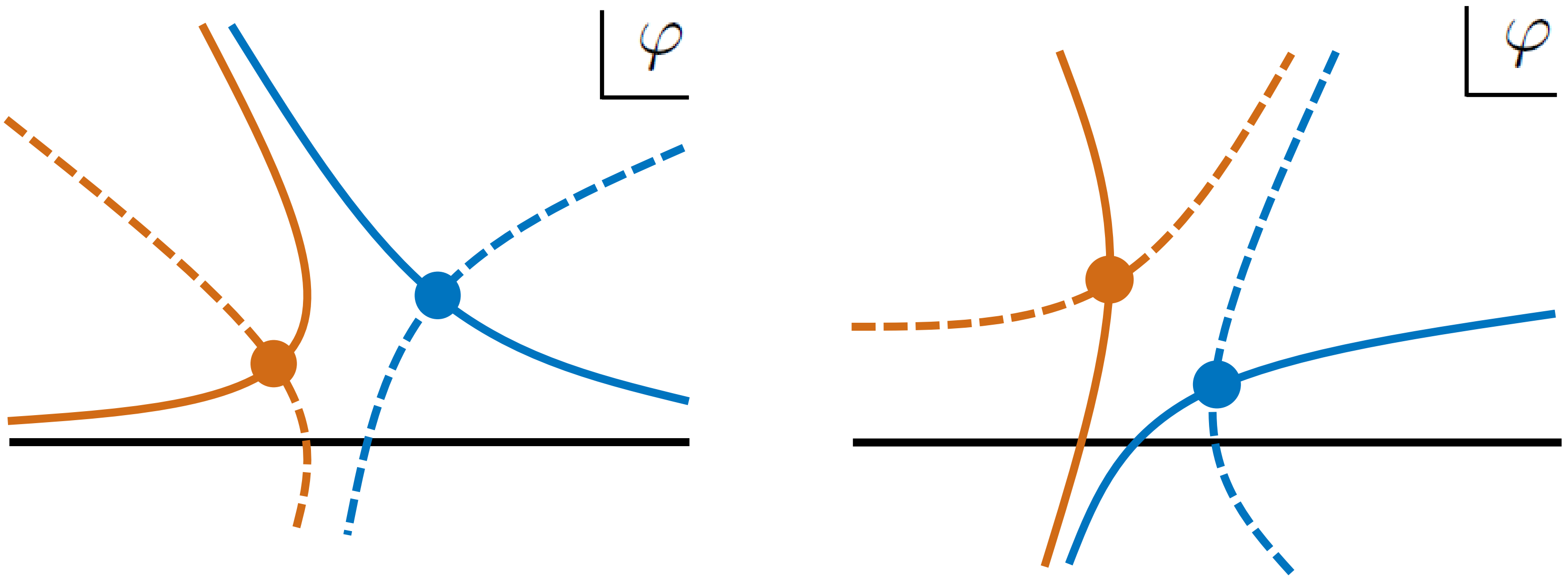
$$\mathcal{K}_i : \frac{\mathrm{d}\varphi(t)}{\mathrm{d}t} = -\overline{\frac{\delta S[\varphi]}{\delta\varphi}}, \quad \varphi(-\infty) = \varphi_i$$

$$\left.\frac{\delta S[\varphi]}{\delta\varphi}\right|_{\varphi=\varphi_i} = 0$$

$$\mathcal{I} = \sum_i n_i \mathcal{J}_i$$ (thimble decomposition)

intersection number

Stokes and anti-Stokes phenomena



Stokes phenomenon : Change of an intersection number, which occurs at

$$\mathrm{Im}[S[\varphi_i]] = \mathrm{Im}[S[\varphi_j]]$$

Anti-Stokes phenomenon: Change of dominant saddles, which occurs at

$$\mathrm{Re}[S[\varphi_i]] = \mathrm{Re}[S[\varphi_j]]$$

$\longleftrightarrow$ 1st order phase transition

What we do

Recall our goal:

- Smallest real part does not necessarily mean that such saddles contribute to the path integral. Can we ***justify it in more precise way***?
- Can we interpret the 2nd order phase transition ***from the viewpoint of resurgence***, and ***draw lessons*** for generic QFTs?

Sub-questions in this part:

- Is the 2nd order phase transition justified by ***thimble decomposition***?
- Is the 2nd order phase transition interpreted as ***(anti-)Stokes phenomena***?

Application to the "path integral"

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

Partition function:

$$Z = \int_{-\infty}^{\infty} \mathrm{d}\sigma \; e^{-S(\sigma)} \qquad S(\sigma) = N_f \left[-i\lambda\sigma + \ln(\cosh\sigma + \cosh m) \right]$$

Thimble/dual-thimble equations:

$$\mathcal{J}_n^{\pm} : \; \frac{\mathrm{d}\sigma(t)}{\mathrm{d}t} = +\overline{\frac{\mathrm{d}S(\sigma)}{\mathrm{d}\sigma}}, \quad \sigma(-\infty) = \sigma_n^{\pm}$$

$$\mathcal{K}_n^{\pm} : \; \frac{\mathrm{d}\sigma(t)}{\mathrm{d}t} = -\overline{\frac{\mathrm{d}S(\sigma)}{\mathrm{d}\sigma}}, \quad \sigma(-\infty) = \sigma_n^{\pm}$$

Lefschetz thimble structure

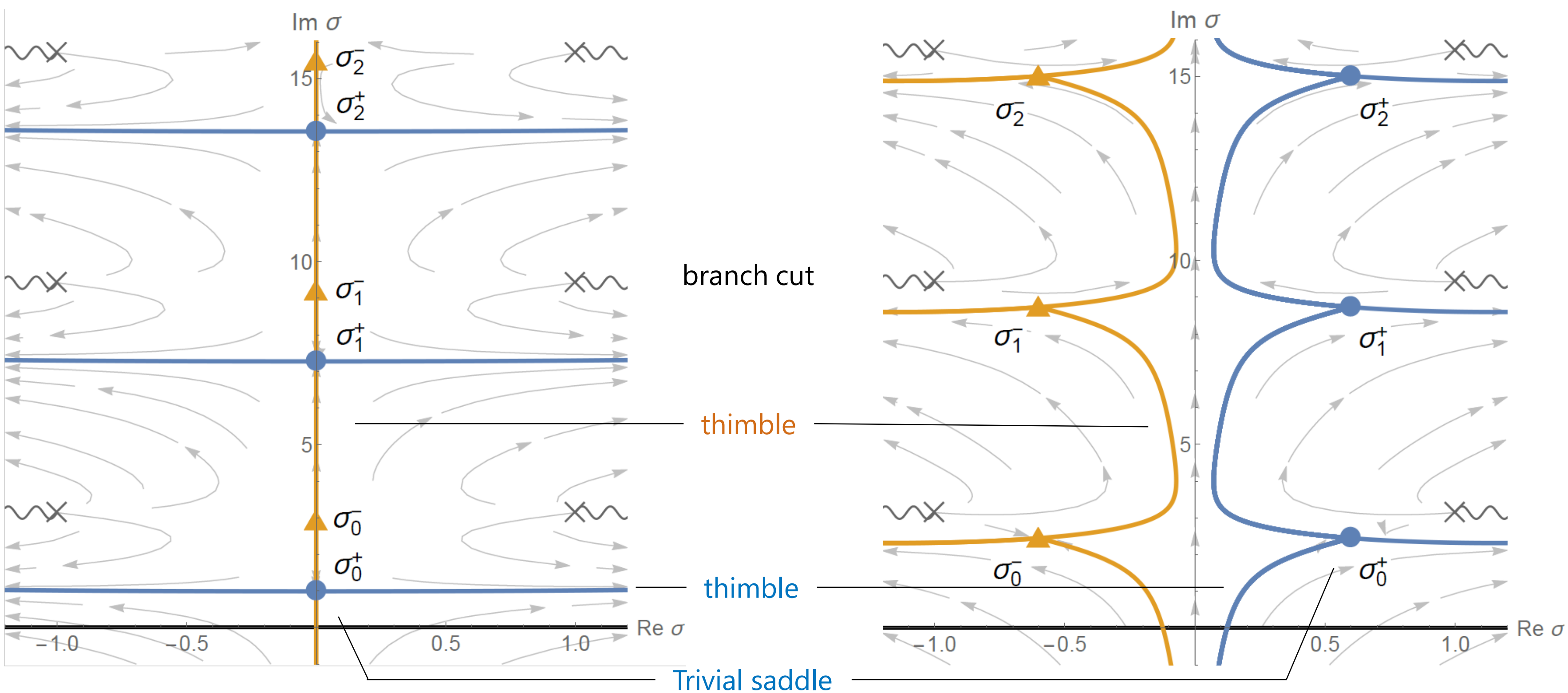
[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

$\boxed{\lambda < \lambda_c}$

$\arg N = 0, \lambda = 0.4, m = 1$

$\boxed{\lambda \geq \lambda_c}$

$\arg N = 0, \lambda = 1.2, m = 1$

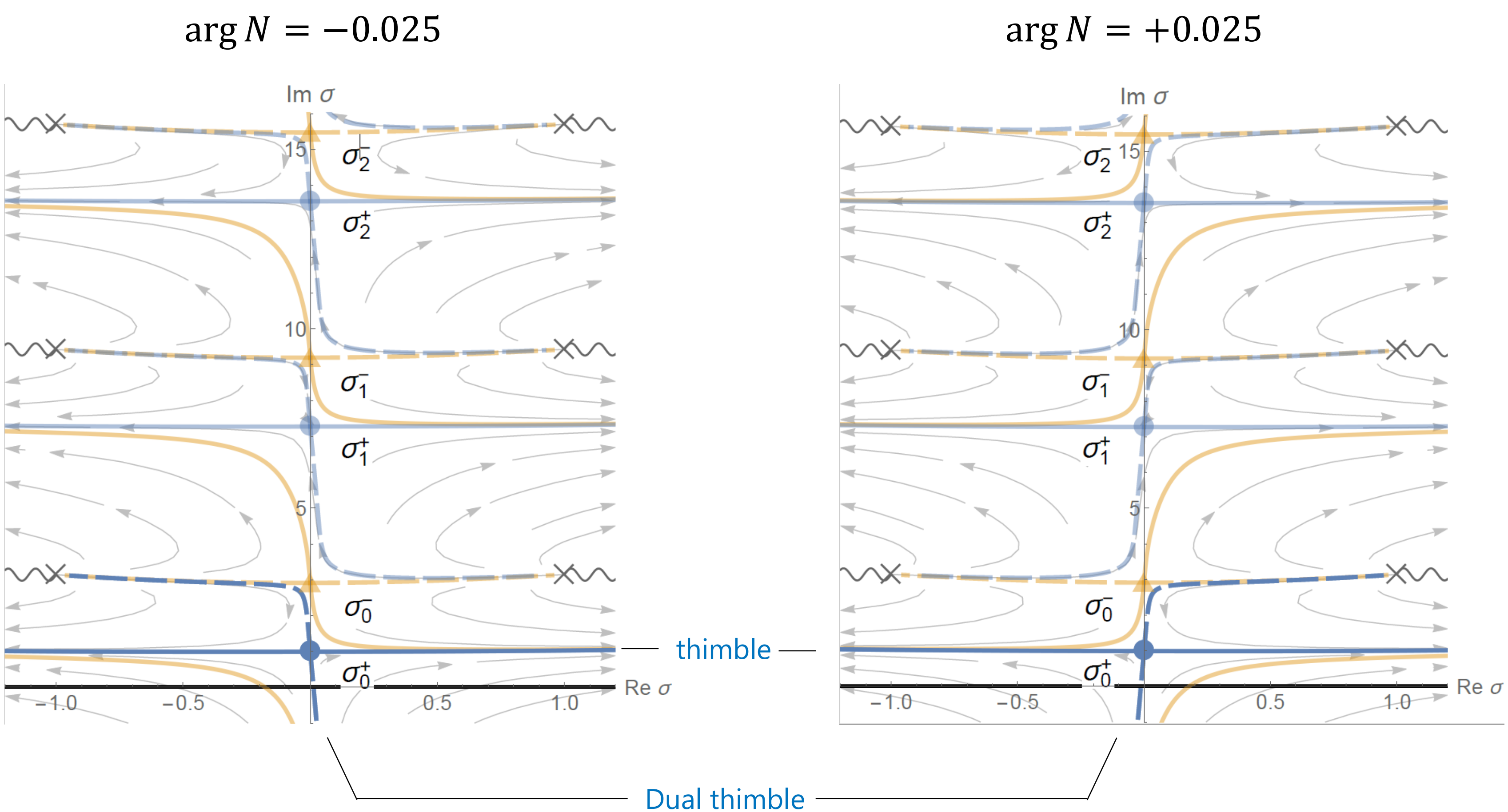


Thimble decomposition $\mathcal{I} = \sum_i n_i \mathcal{J}_i$ is ambiguous

$\longrightarrow$ Vary the phase of N_f

Lefschetz thimble structure (subcritical)

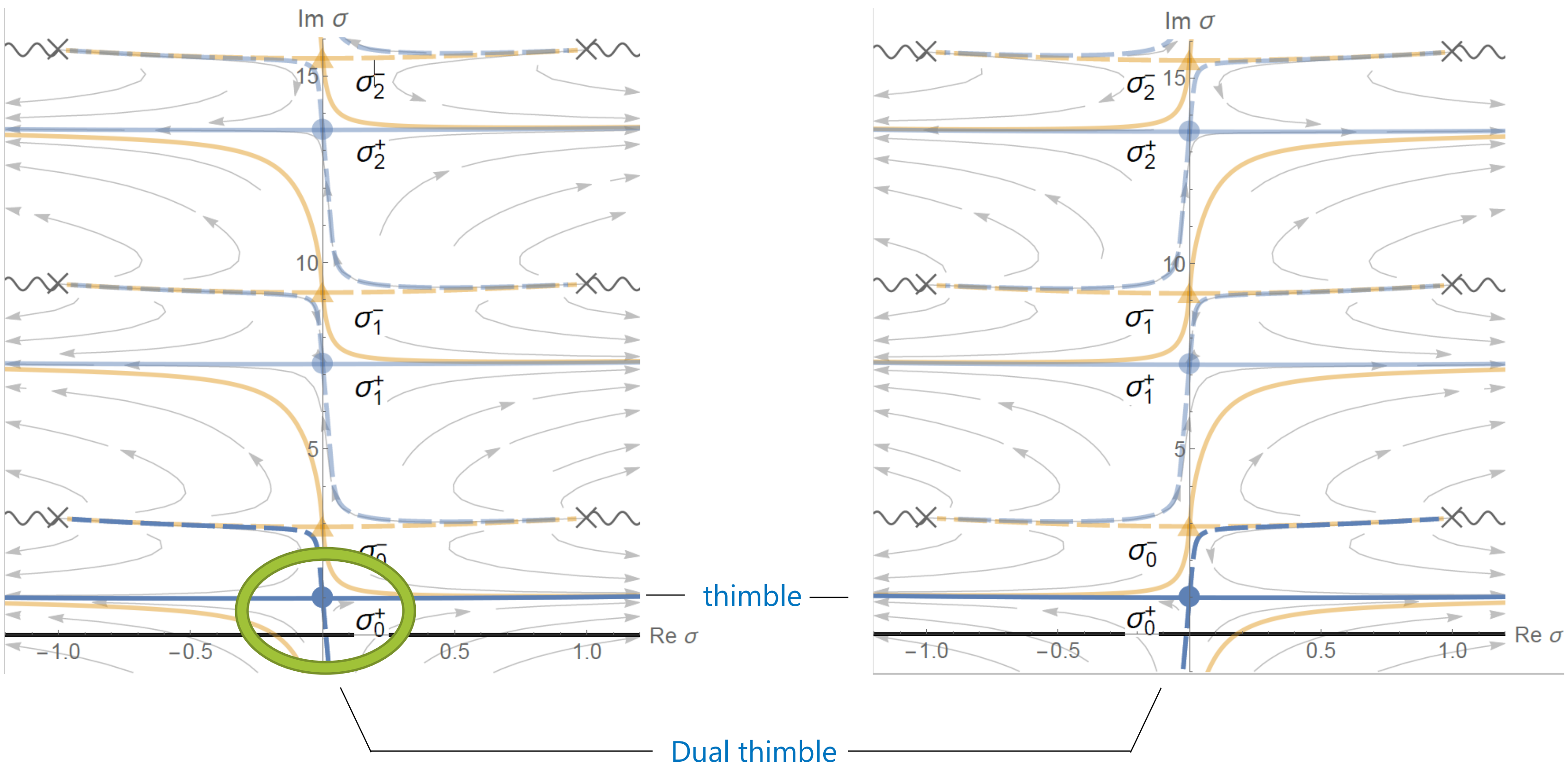
$\boxed{\lambda < \lambda_c}$



Lefschetz thimble structure (subcritical)

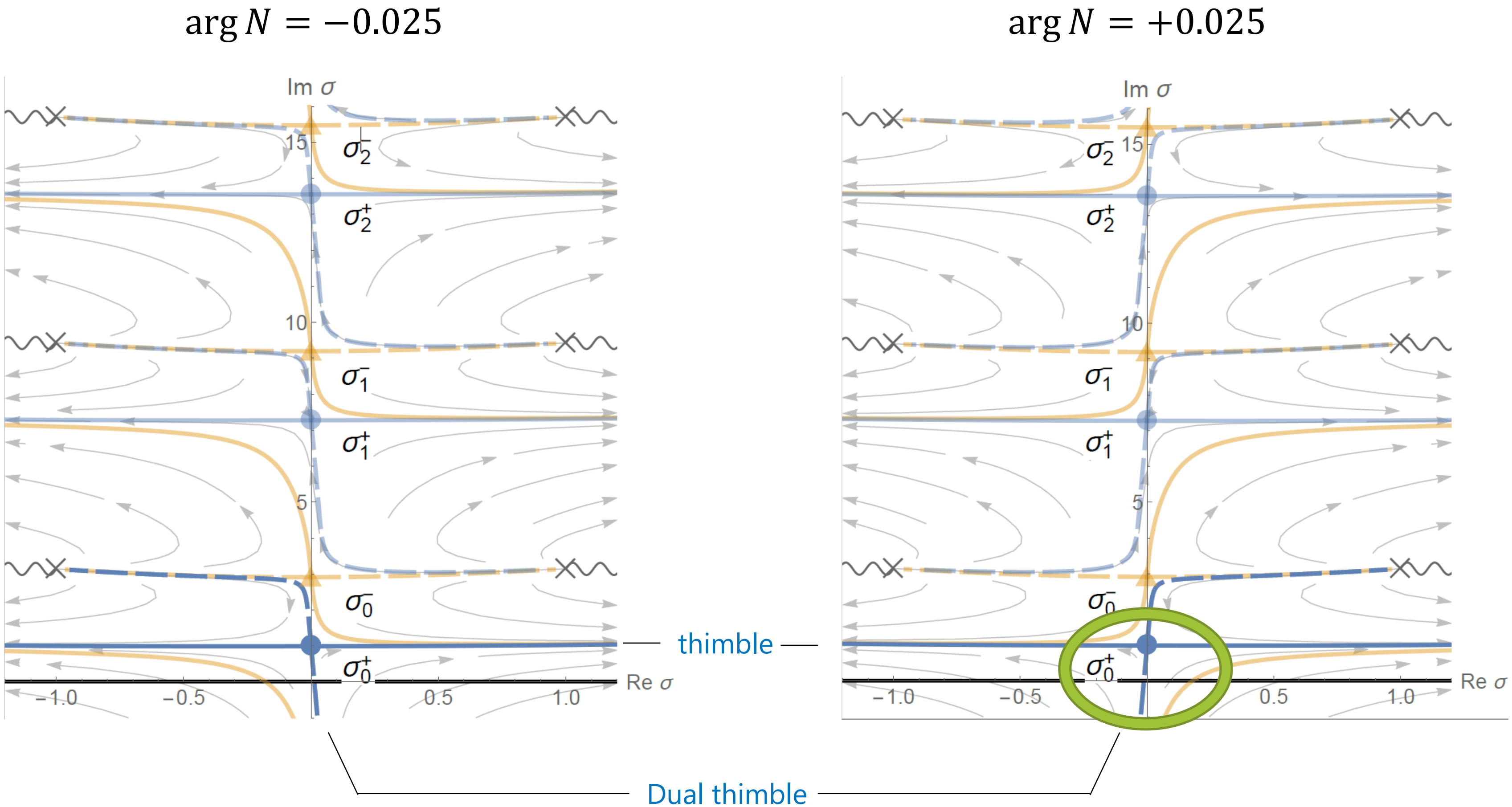
$\boxed{\lambda < \lambda_c}$

$\arg N = -0.025$ $\qquad\qquad$ $\arg N = +0.025$



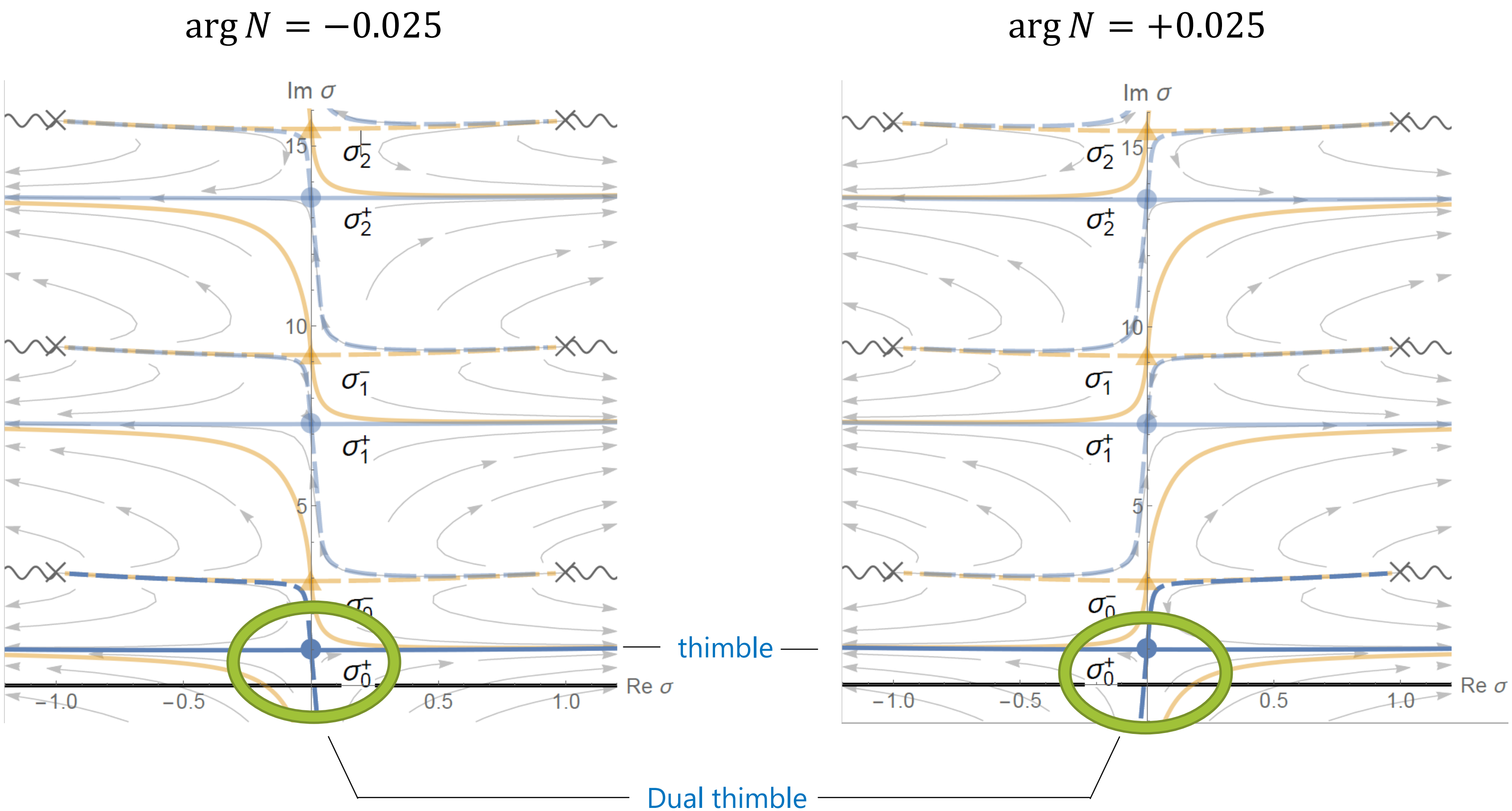
Lefschetz thimble structure (subcritical)

$\boxed{\lambda < \lambda_c}$



Lefschetz thimble structure (subcritical)

$\boxed{\lambda < \lambda_c}$

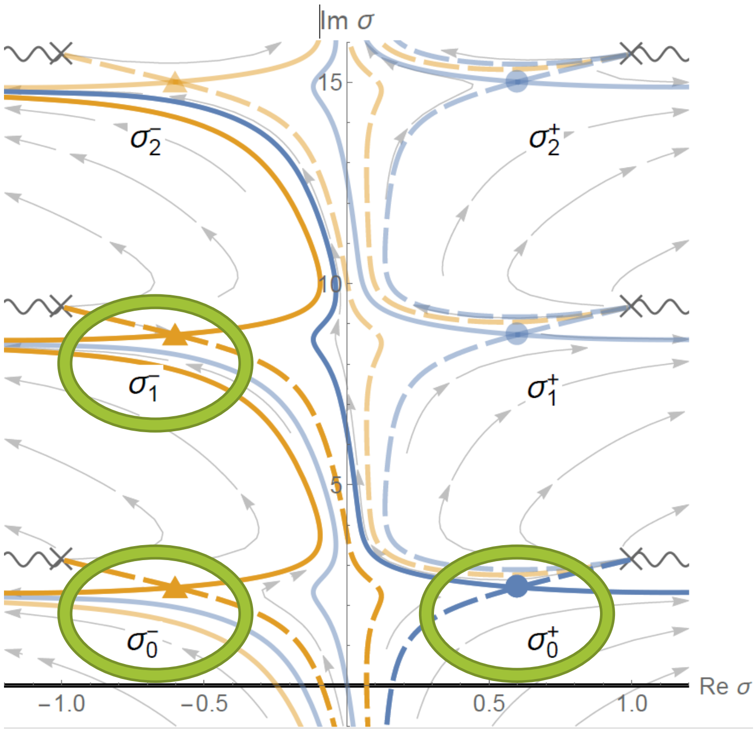


- No Stokes phenomenon
- Only the trivial saddle σ_0^+ contributes to the path integral

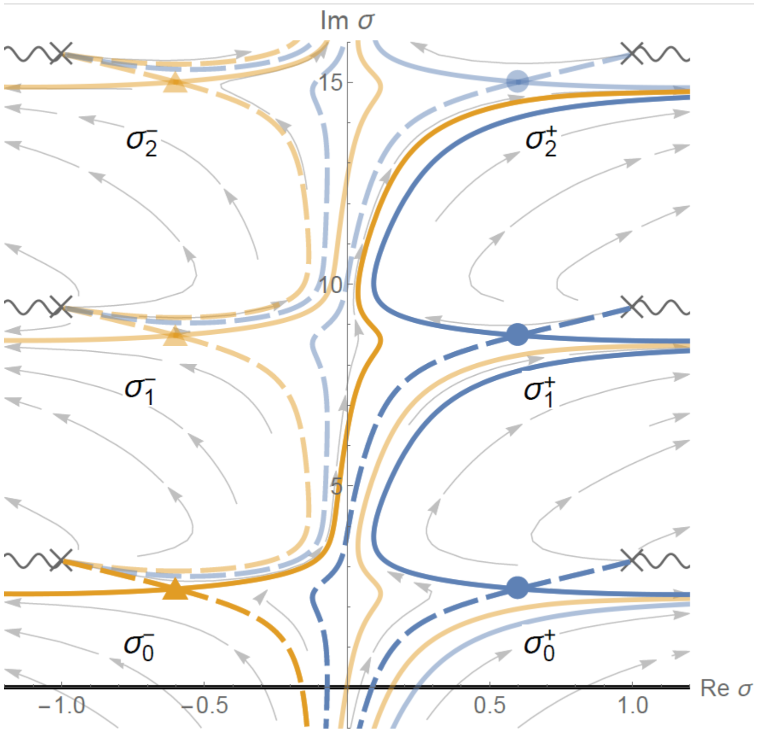
Lefschetz thimble structure (supercritical)

$\boxed{\lambda \geq \lambda_c}$

$\arg N = -0.025$



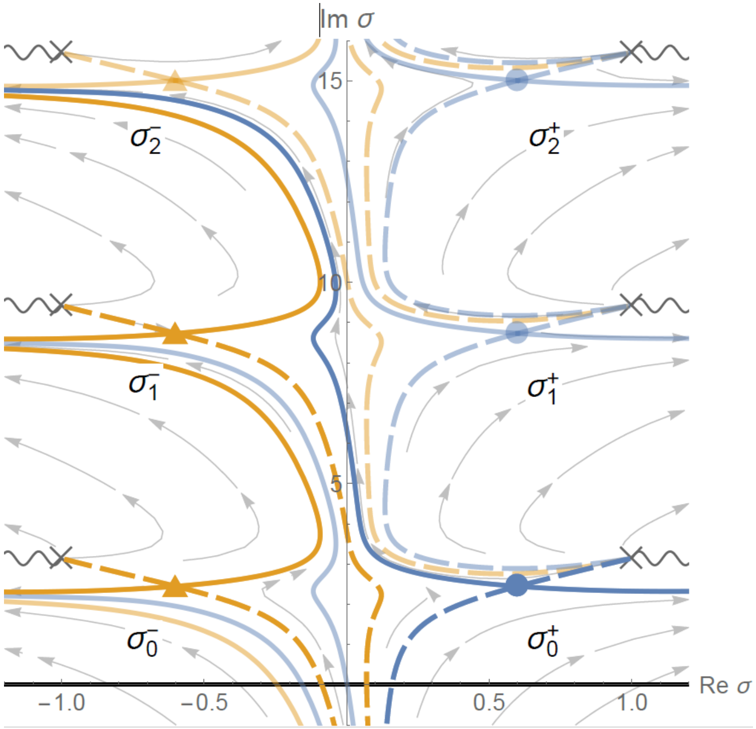
$\arg N = +0.025$



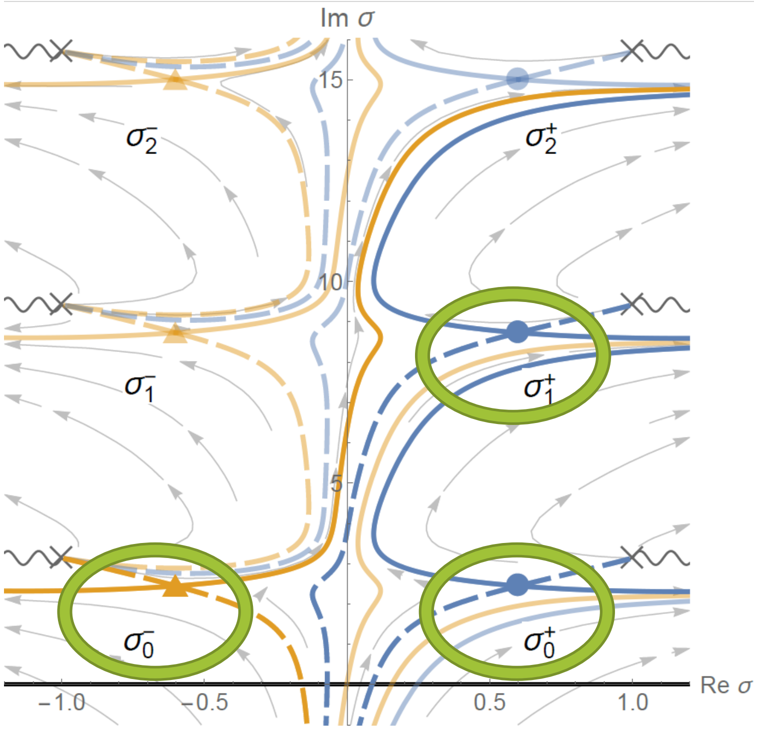
Lefschetz thimble structure (supercritical)

$\boxed{\lambda \geq \lambda_c}$

$\arg N = -0.025$



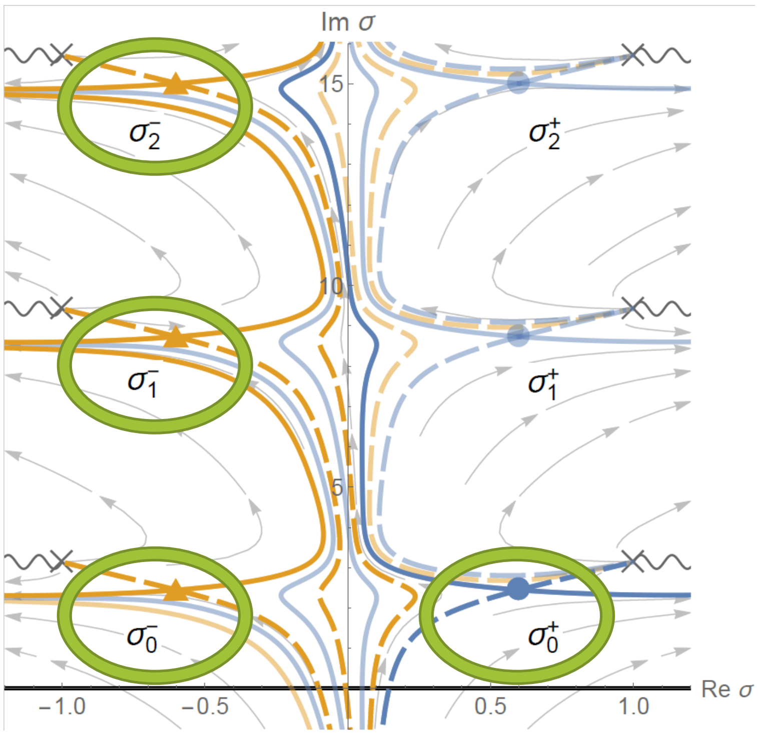
$\arg N = +0.025$



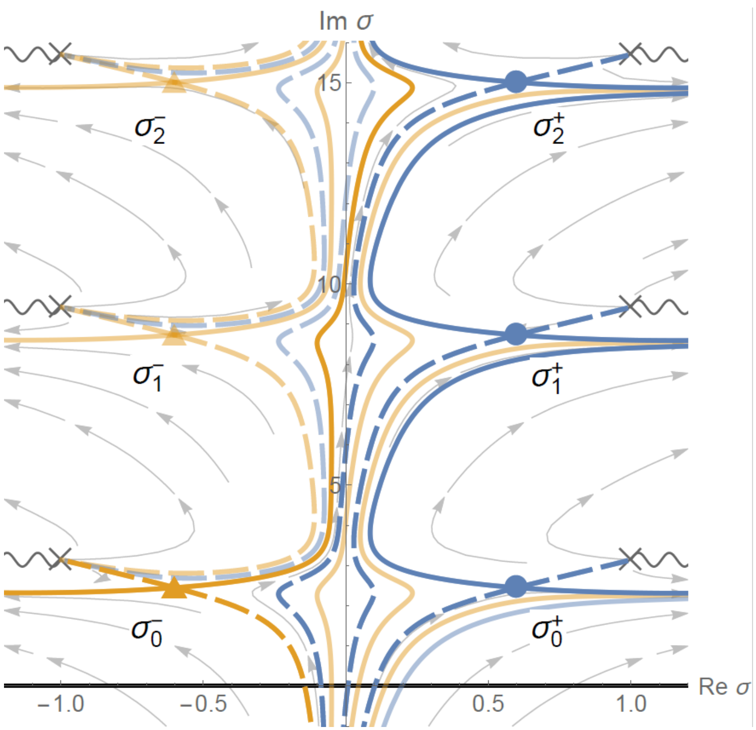
Lefschetz thimble structure (supercritical)

$\boxed{\lambda \geq \lambda_c}$

$\arg N = -0.015$



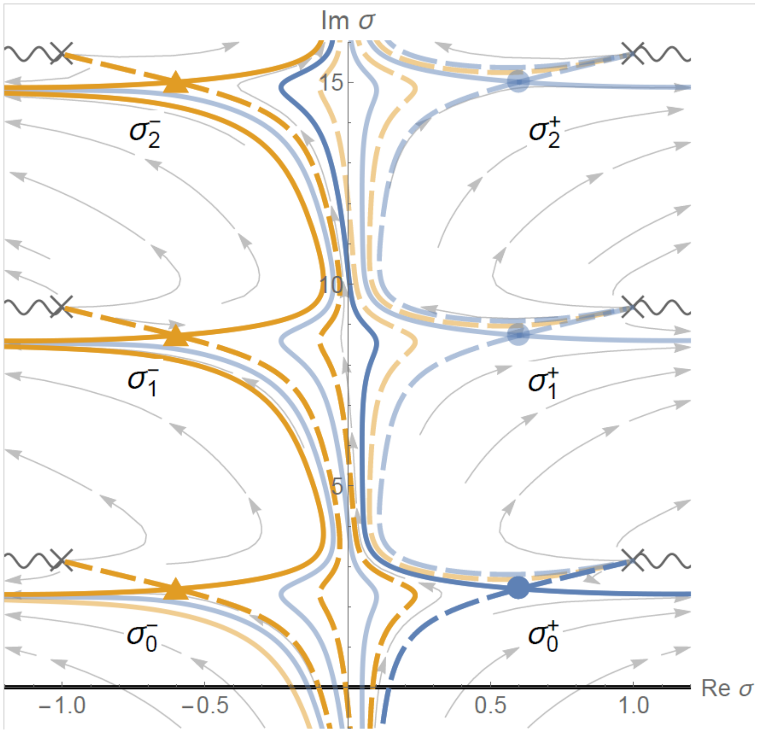
$\arg N = +0.015$



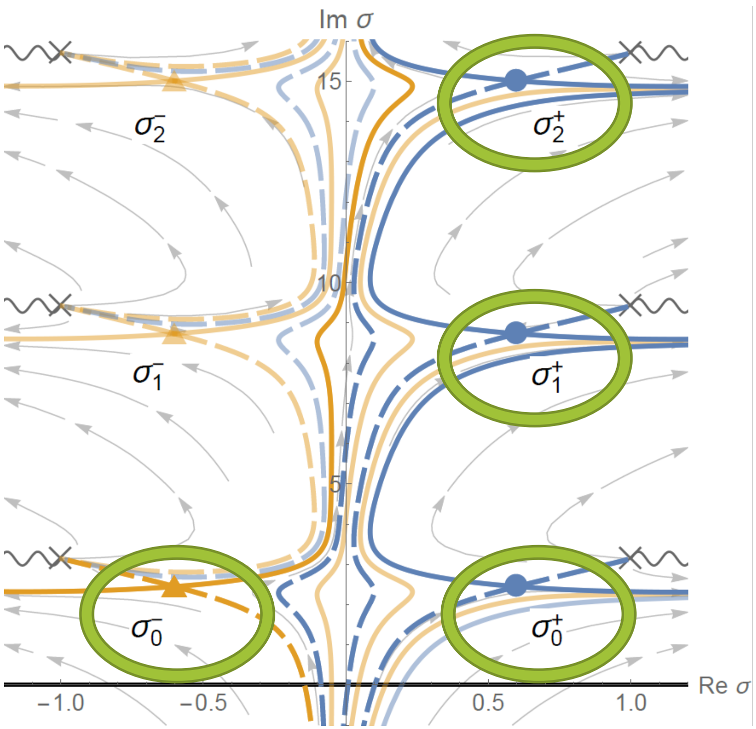
Lefschetz thimble structure (supercritical)

$\boxed{\lambda \geq \lambda_c}$

$\arg N = -0.015$



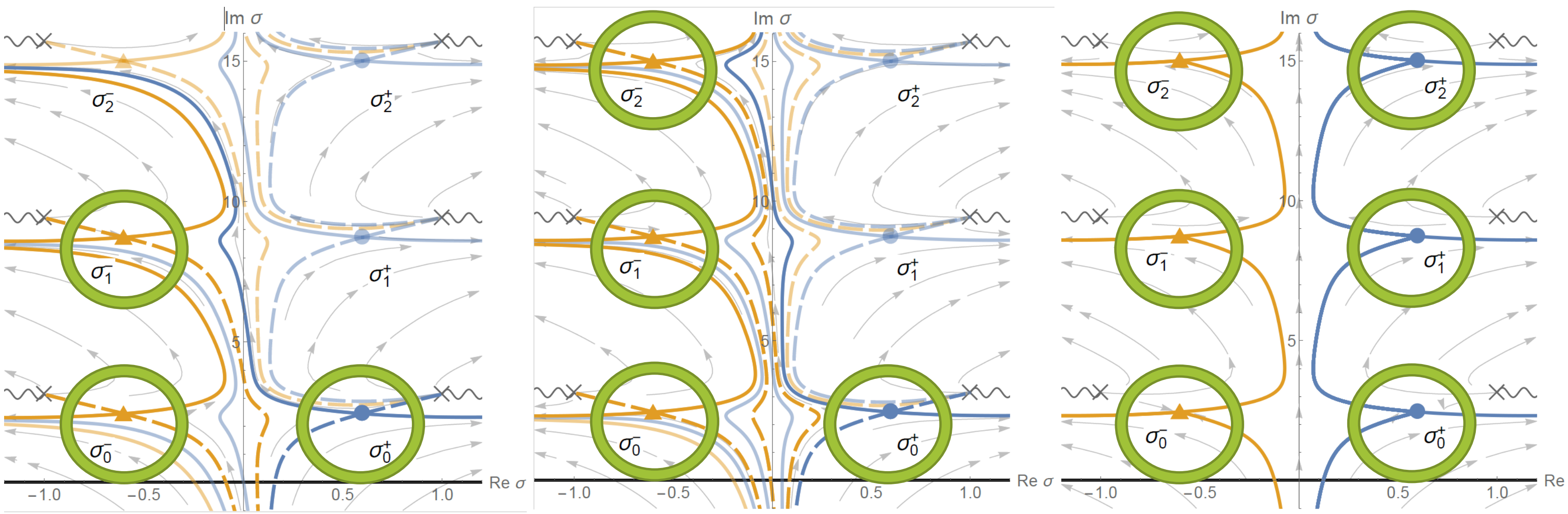
$\arg N = +0.015$



Lefschetz thimble structure (supercritical)

$\boxed{\lambda \geq \lambda_c}$

$\arg N = -0.025$ $\arg N = -0.015$ $\arg N = 0$

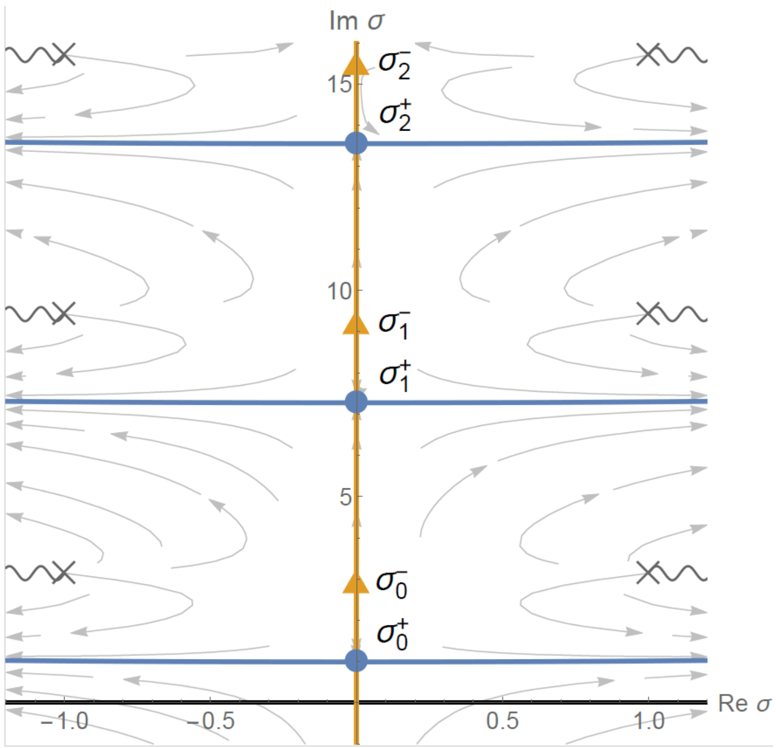


...

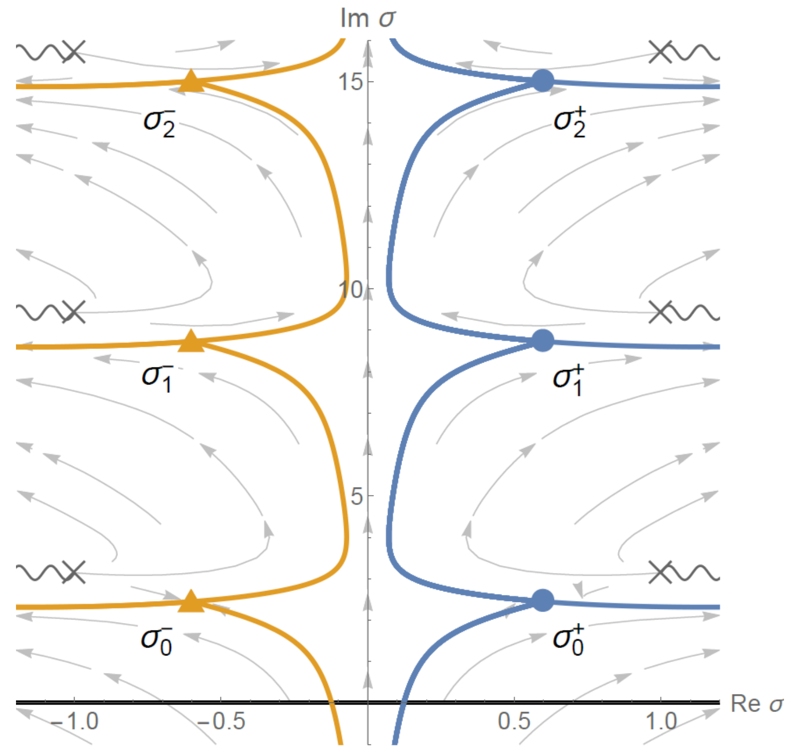
- Infinite number of Stokes phenomena occur around $\arg N_f = 0$
- Infinite number of saddles $\sigma_n^{\pm}$ contribute to the path integral

Stokes and anti-Stokes pheno. at the critical pt.

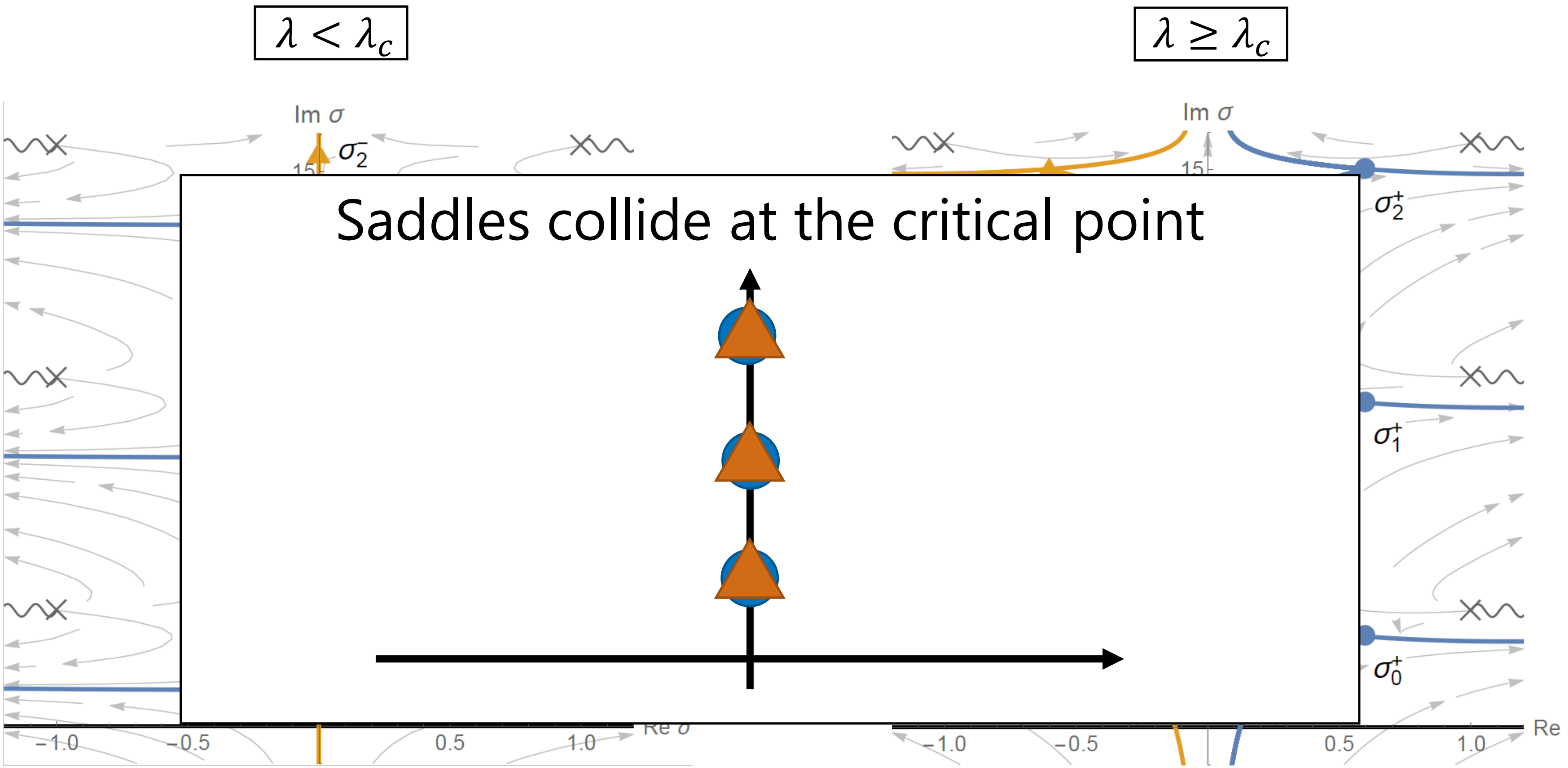
$\boxed{\lambda < \lambda_c}$



$\boxed{\lambda \geq \lambda_c}$



Stokes and anti-Stokes pheno. at the critical pt.



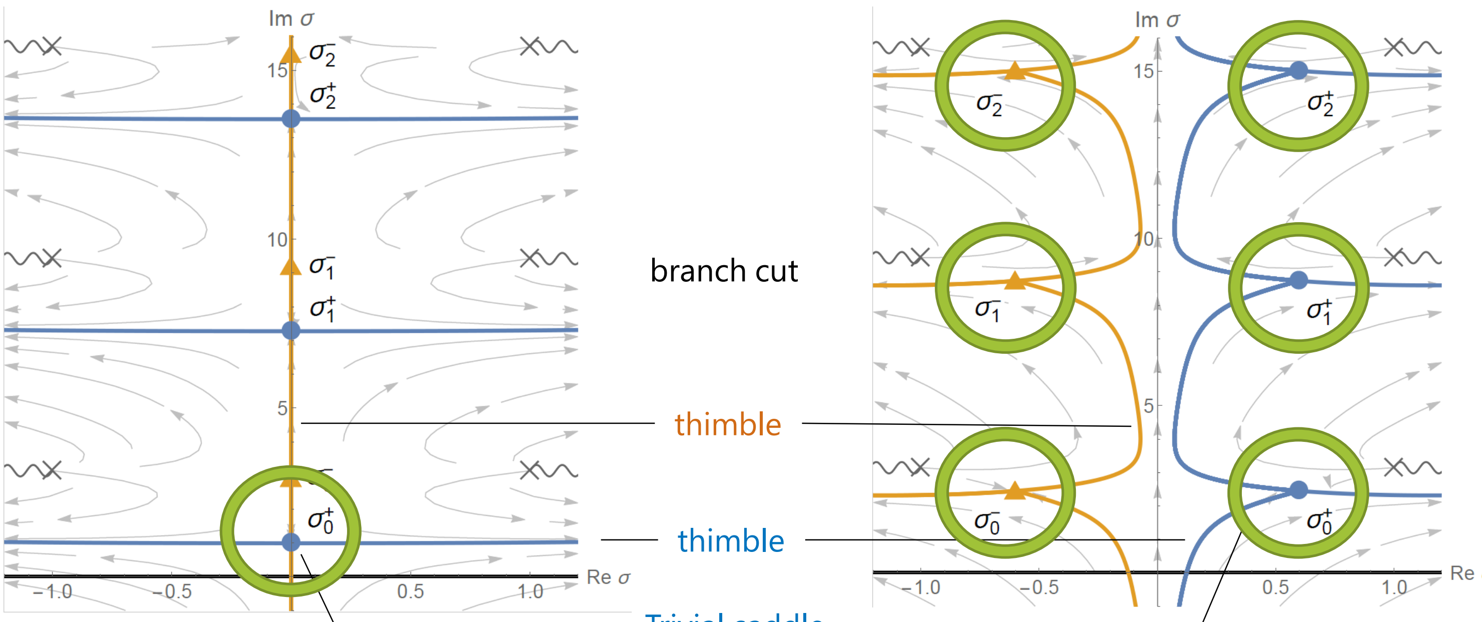
At the critical point,

- we have seen a Stokes phenomenon $\sigma_0^+ \to \sigma_n^\pm$
- anti-Stokes phenomenon occurs at the same time

Phase transition and Lefschetz thimble structure

Summary

$\boxed{\lambda < \lambda_c}$	$\boxed{\lambda \geq \lambda_c}$
• No Stokes phenomenon • Only the trivial saddle σ_0^+ contributes	• Infinite number of Stokes phenomena • Infinite number of saddles $\sigma_n^\pm$ contribute (Two of which σ_0^+, σ_0^- survive the large-flavor limit)

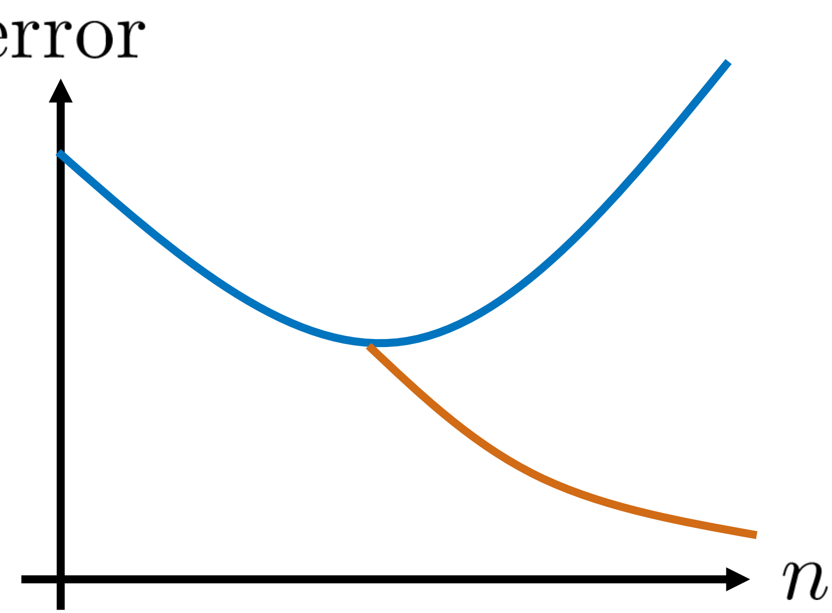


- Stokes and anti-Stokes phenomena occur at the same time

⟶ The phase transition is interpreted from the view point of thimbles

Contents

Asymptotic series and resurgence structure



$$Z(g) = \sum_{n=0}^{\infty} c_n g^{n+r}, \quad c_n \sim n!$$

Resurgence structure

$$Z(g) = \boxed{\sum_{n=0}^{\infty} c_n^{(0)} g^n} + \boxed{\sum_{n=0}^{\infty} c_n^{(1)} g^n e^{-S_1/g} + \cdots}$$

[J. Ecalle, 81]
[M. Marino, 12]
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[D. Dorigoni, 19]
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Resurgence theory:

the perturbative part knows the non-perturbative parts

Borel resummation

Resuming a asymptotic series

[J. Ecalle, 81]

Lectures and reviews, e.g.
[M. Marino, 12]
[D. Dorigoni, 19]
[I. Aniceto, G. Basar, R. Schiappa, 19]

$$Z(g) = \sum_{n=0}^{\infty} c_n g^{n+r}, \quad c_n \sim n!$$

Asymptotic series

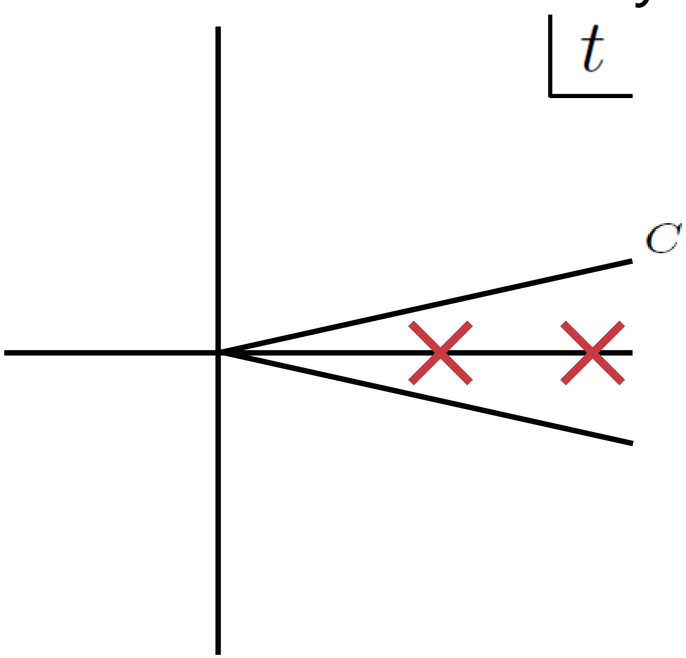
$$\mathcal{S}Z(g) = \int_C \mathrm{d}t \; e^{-t/g} \mathcal{B}Z(t)$$

Borel resummation

$$\mathcal{B}Z(t) = \sum_{n=0}^{\infty} \frac{c_n}{\Gamma(n+r)} t^{n+r-1}$$

Borel transformation

Borel transformation may have Borel singularities



Borel singularities mean
non-perturbative corrections!

Non-trivial saddle and Borel singularities

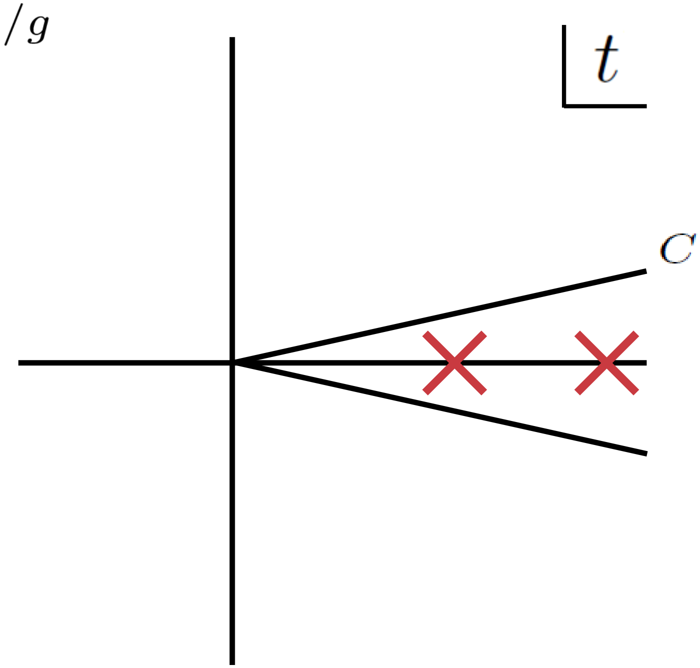
[Lipatov, 77]

$$Z(g) = \int \mathcal{D}\varphi \, e^{-S[\varphi]/g} \sim \sum_{n=0}^{\infty} c_n g^n$$

$$c_n = \frac{1}{2\pi i} \oint \frac{\mathrm{d}g}{g^{n+1}} \int \mathcal{D}\varphi \, e^{-S[\varphi]/g}$$

$$\sim e^{-S[\varphi_*]/g_* - (n+1)\ln g_*}$$

$$\sim \frac{n!}{(S[\varphi_*])^{n+1}}$$



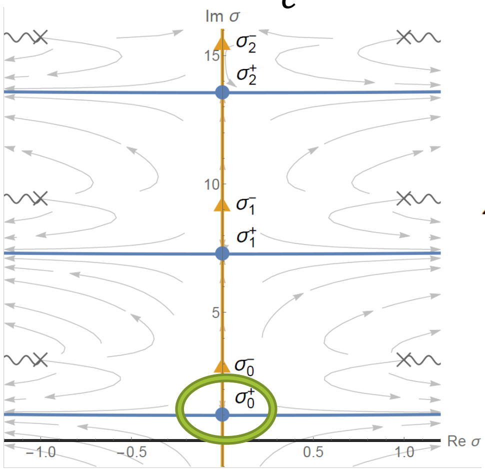
$$\mathcal{B}Z(t) \sim \sum_{n=0}^{\infty} \left(\frac{t}{S[\varphi_i]} \right)^n = \frac{1}{1 - \frac{t}{S[\varphi_i]}}$$

Non-trivial saddles are encoded in an asymptotic series

What we do

Large-flavor expansion around the trivial saddle:

$\lambda < \lambda_c$



asymptotic series!

$$Z(\lambda;N) = \frac{1}{2^N}\int d\sigma\ e^{-NS(\lambda;\sigma)}, \quad S(\lambda;\sigma) = -i\lambda\sigma + \ln(\cos \quad + \cosh m)$$

$$\overset{\text{around } \sigma_0^+}{=} \frac{1}{2^N}\sqrt{\frac{2\pi}{NS''(\lambda;\sigma_0^+)}}e^{-NS(\lambda;\sigma_0^+)}\sum_{l=0}^{\infty}\frac{a_l(\lambda)}{N^l}$$

Borel resummation:

$$\mathcal{S}Z(\lambda;N) = \frac{1}{2^N}\sqrt{\frac{2\pi}{NS''(\lambda;\sigma_0^+)}}e^{-NS(\lambda;\sigma_0^+)}\cdot N\int_C dt\ e^{-Nt}\sum_{l=0}^{\infty}\frac{a_l(\lambda)}{\Gamma(l+1)}t^l$$

Sub-questions in this part:

- Is the Borel plane structure ***consistent with the Lefschetz thimble structure***?
- Can we ***decode the phase transition*** from the perturbative series?

But please wait (1/2): Borel-Padé approximation

Exact quantities:

$$F\left(\frac{1}{N_f}\right) = \sum_{\ell=0}^{\infty} \frac{a_l}{N_f^{\ell}}$$

asymptotic series

$$\mathcal{B}F(t) = \sum_{l=0}^{\infty} \frac{a_\ell}{\Gamma(\ell+1)} t^{\ell}$$

Borel transformation

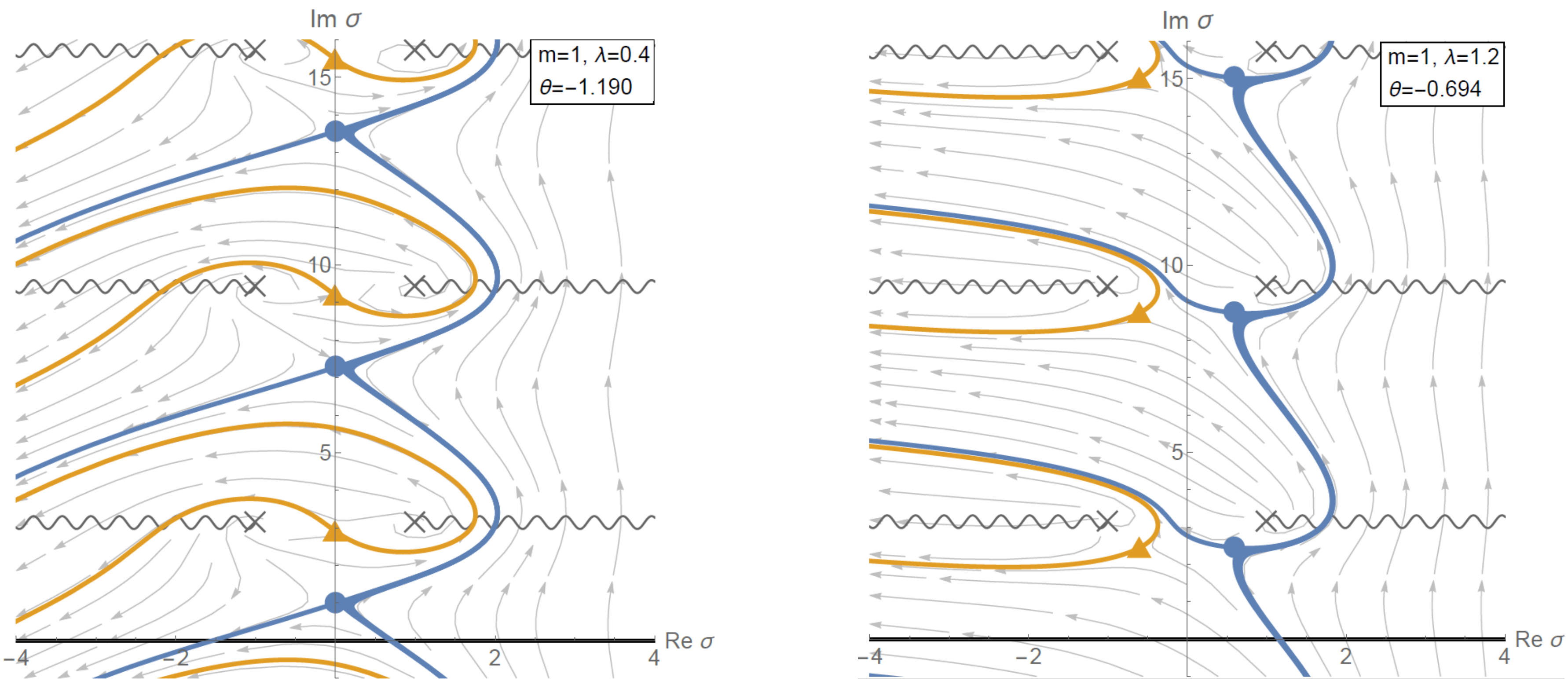
Approximate these from finite number of inputs

$$\mathcal{P}_{m,n}(t) = \frac{P_m(t)}{Q_n(t)}$$

Borel-Padé approximation

But please wait (2/2): Larger $\theta = \arg N_f$

Stokes phenomena occur on different Riemann sheets



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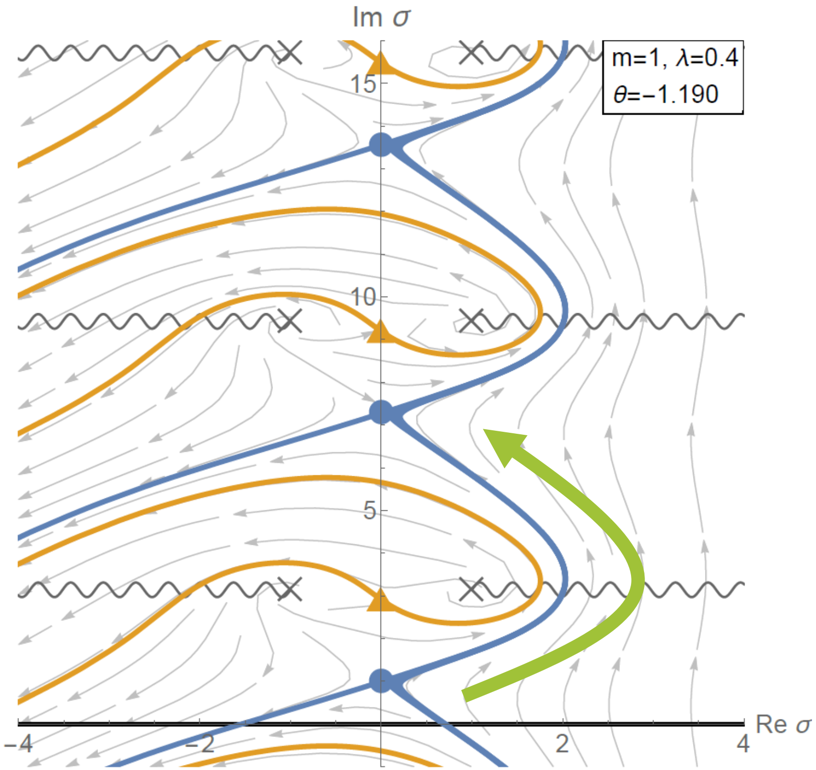
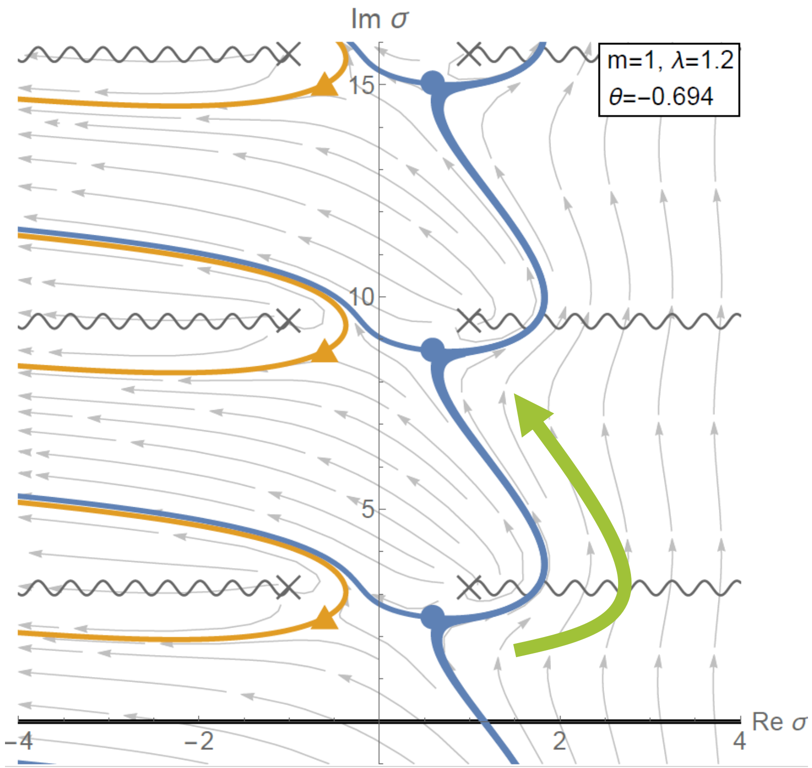


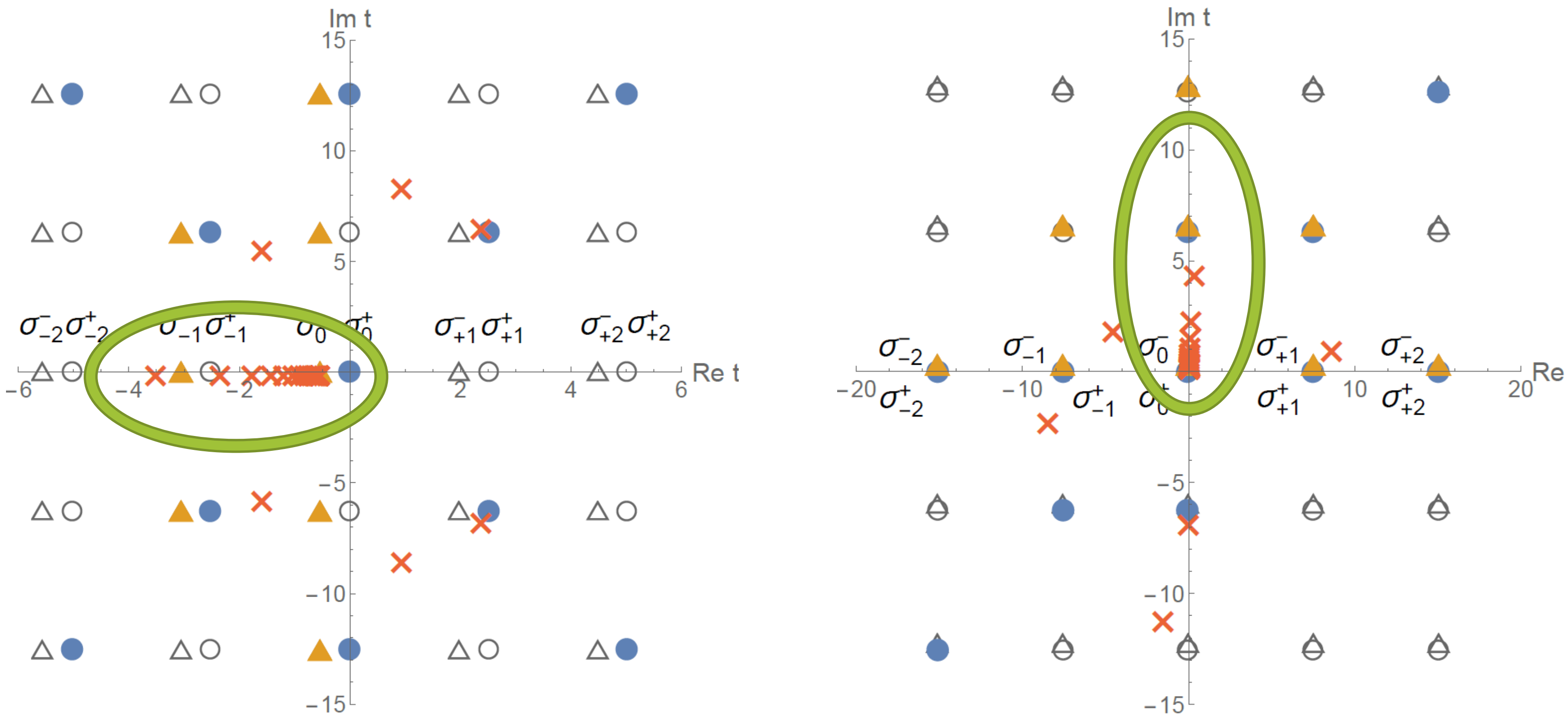


Borel plane structure

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

$\boxed{\lambda < \lambda_c}$ $\qquad$ $\boxed{\lambda \geq \lambda_c}$

$\lambda = 0.4, m = 1$, Padé (25,25) $\qquad$ $\lambda = 1.4, m = 1$, Padé (25,25)



Poles of the Padé approximant are consumed for branch cuts...

Improvement: Padé-Uniformized approximation

[Costin, Dunne, 20]

Uniformize the Borel t-plane by a map:

$$t \mapsto u(t) = -\ln\left(1 - \frac{t}{s}\right)$$

Branch cut singularity at $t = s$ is eliminated

Perform the standard Padé approximation on the u-plane

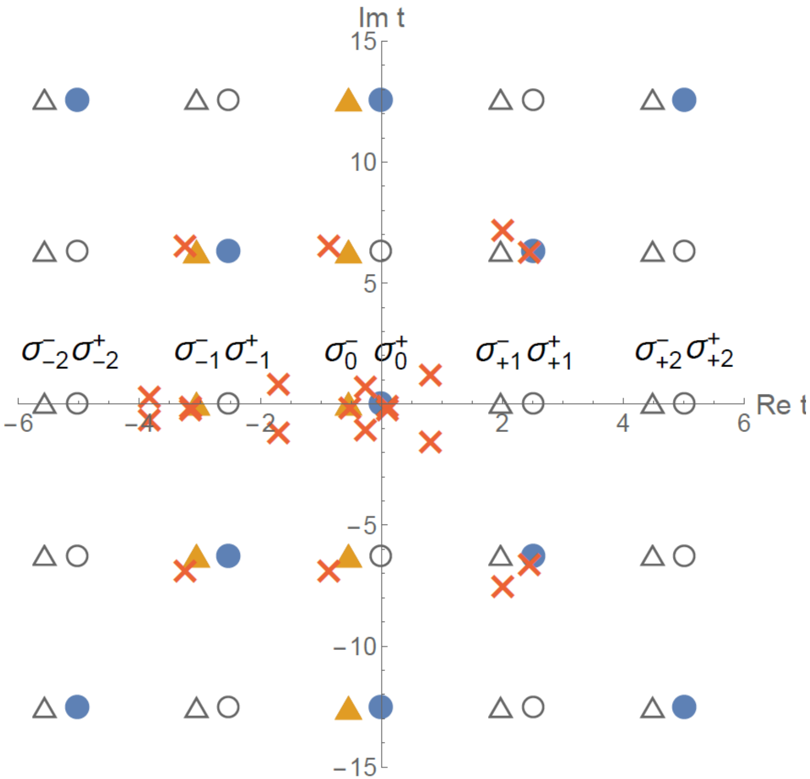
$$\widetilde{\mathcal{B}F}(t) \simeq \mathcal{P}_{m,n}(u(t))$$

Borel plane structure (improved)

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

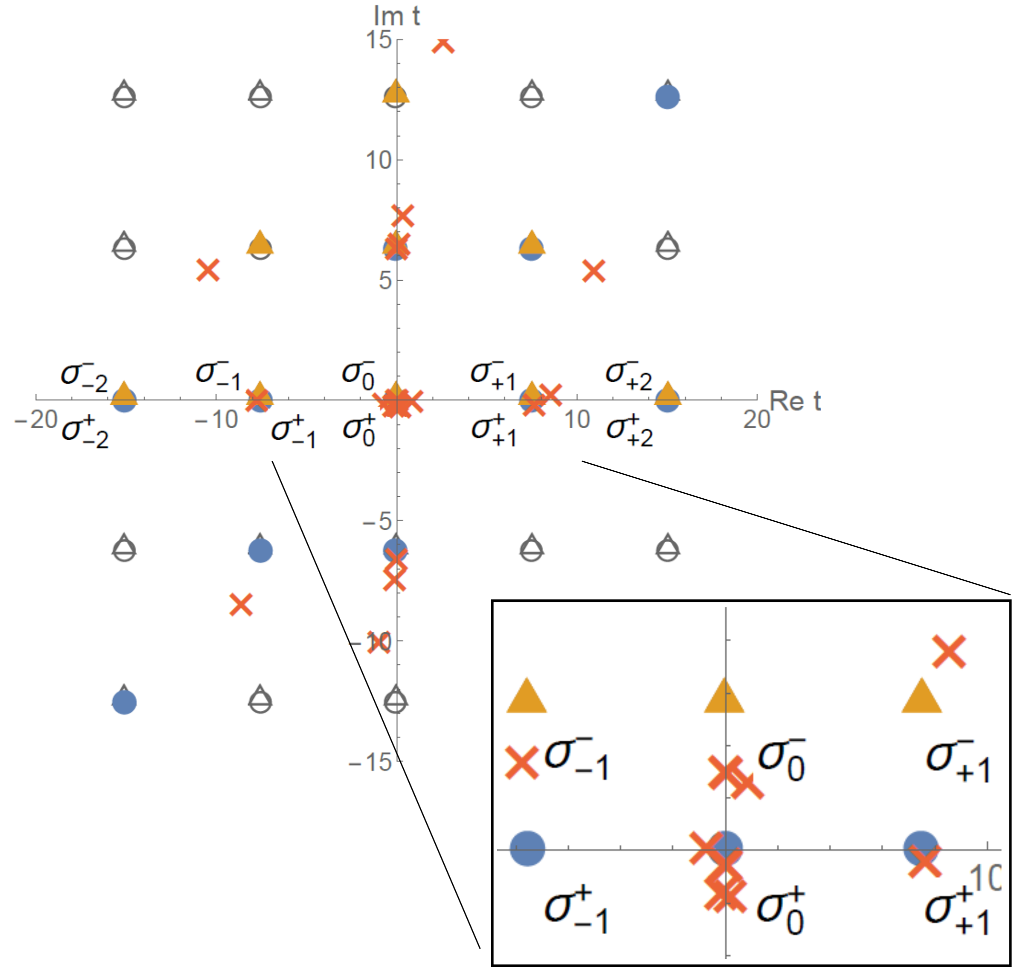
$\boxed{\lambda < \lambda_c}$

$\lambda = 0.4, m = 1$, Padé-Uniformized (25,25)



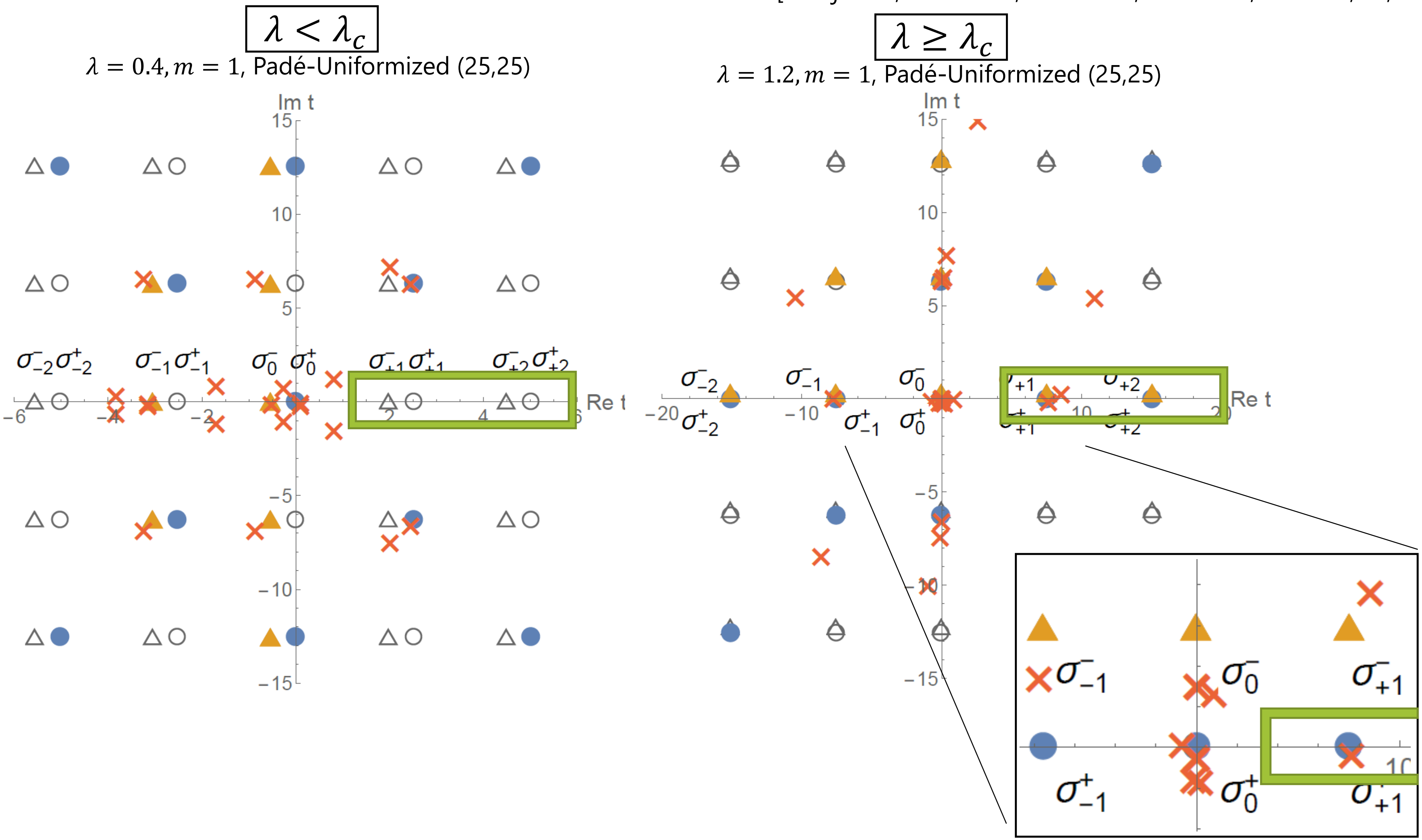
$\boxed{\lambda \geq \lambda_c}$

$\lambda = 1.2, m = 1$, Padé-Uniformized (25,25)



Borel plane structure (improved)

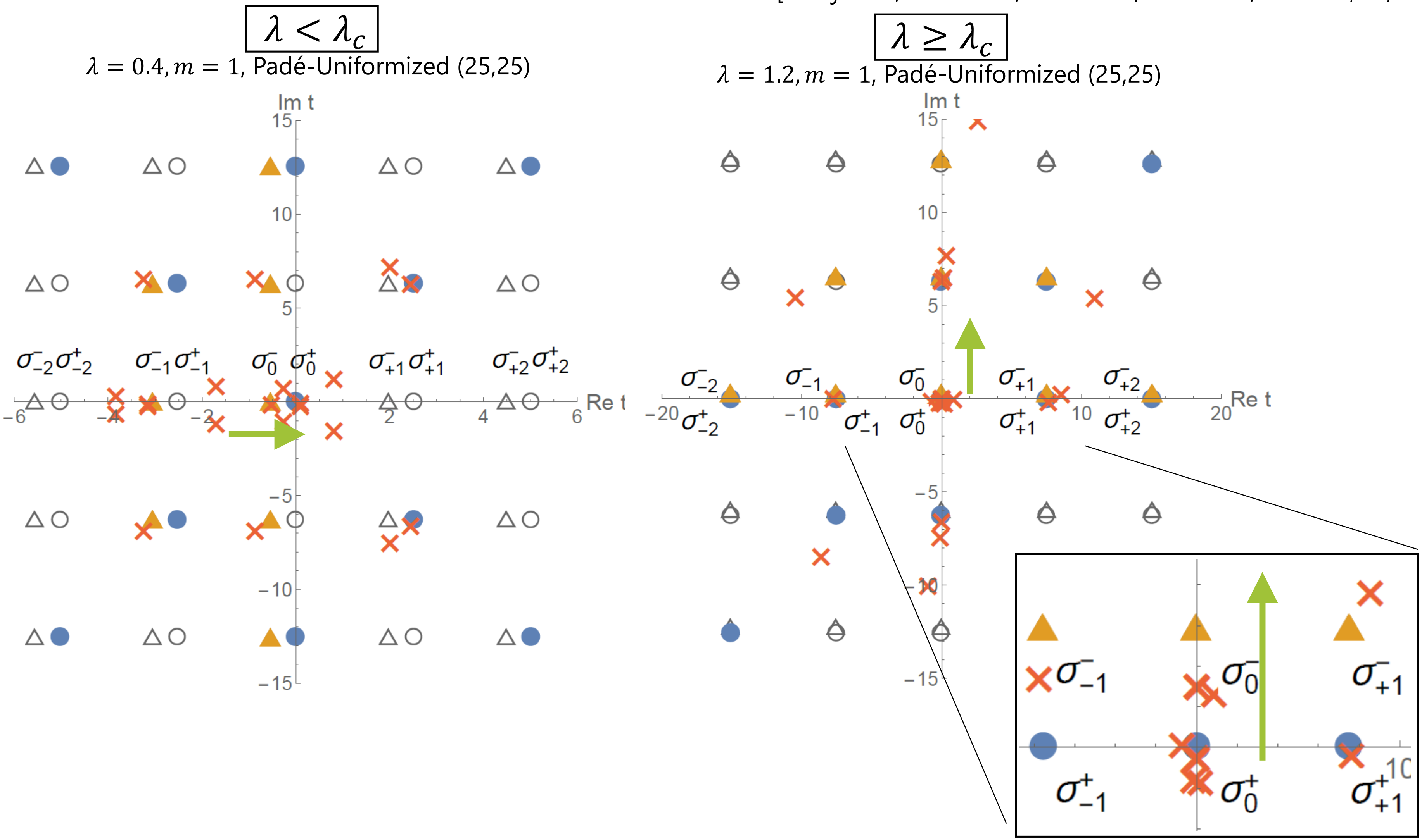
[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]



- The Stokes phenomena are encoded as Borel non-summability

Borel plane structure (improved)

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]



- The collision of saddles are encoded as collision of Borel singularities
- The anti-Stokes phenomenon is encoded as Borel singularities along the vertical axis

Analytical study for large λ

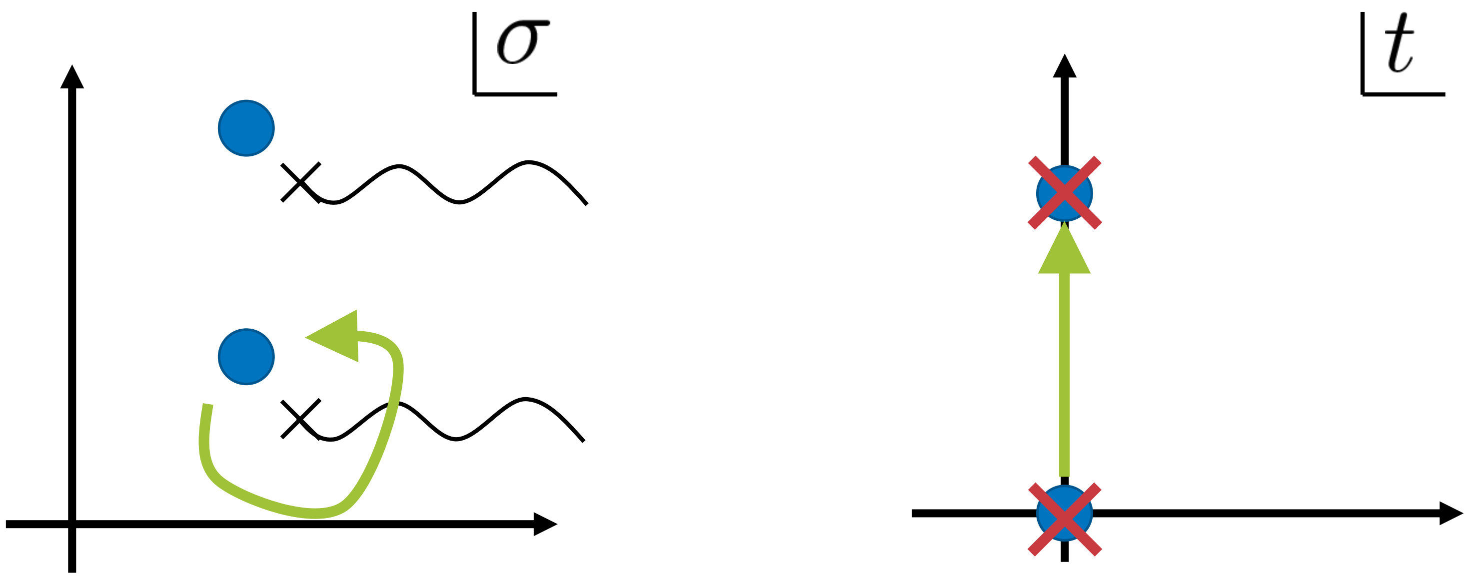
[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

Lading contribution for large λ

$$F(N_f;\lambda) = \int_{-\infty}^{\infty} \mathrm{d}\delta\sigma \; e^{-N_f(i\lambda\delta\sigma + \log(1 - i\lambda\delta\sigma))}$$

$$= \frac{1}{i\lambda}\int dt \; e^{-N_f t} \frac{W(-e^{t-1})}{1+W(-e^{t-1})}$$

σ-plane and Borel t-plane are directly related via $\delta\sigma = \dfrac{1}{i\lambda}\left(1 + W(-e^{t-1})\right)$



Contents

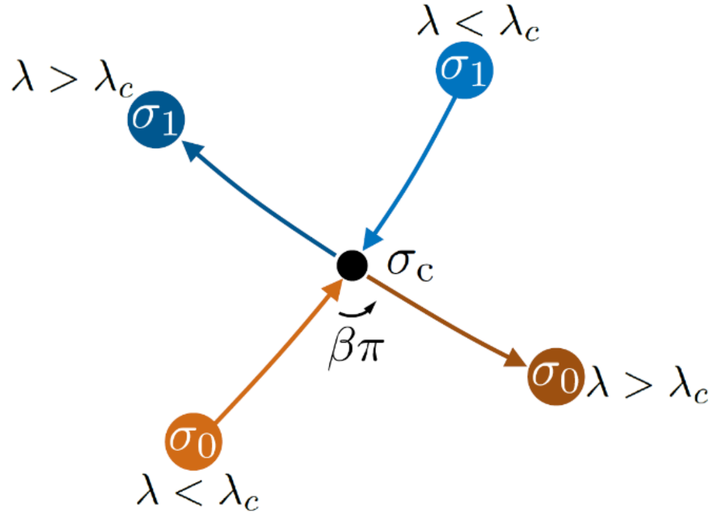
Collision of saddles

[T. Fujimori, M. Honda, S. Kamata, T. Misumi, N. Sakai, TY, 21]

Consider

$$e^{-NF(\lambda)} = \int \mathrm{d}\sigma \; e^{-N\tilde{S}(\lambda;\sigma)}$$

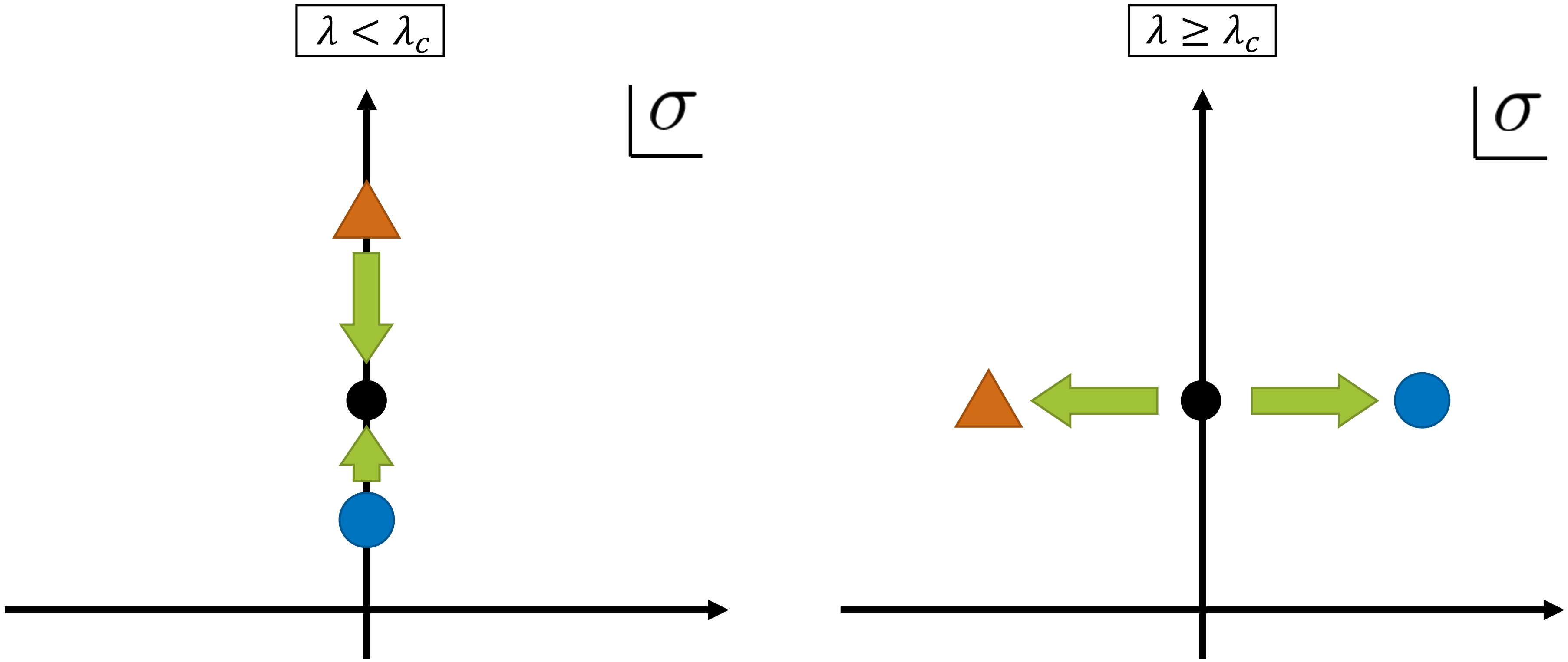
If the "action" is holomorphic, and n saddles collide as



Then, the "action" value at m-th saddle is $\tilde{S}_m \simeq c_0 + T_m(\delta\lambda)^{(n+1)\beta}$

⟶ Phase transition is of order $\lceil (n+1)\beta \rceil$

Revisiting the SQED3



$$n = 2, \beta\pi = \pi/2$$

$\longrightarrow$ Phase transition is of order $\lceil (2+1)/2 \rceil = 2$

Revisiting the SQED3

The 2nd order phase transition corresponds to followings

Stokes and
anti-Stokes phenomena
at the same time

Lefschetz thimble analysis:

i. Contributing saddles jump as $\sigma_0^+ \to \sigma_0^+, \sigma_0^-$

ii. Two saddles collide with an angle $\pi/2$

iii. Infinite number of Stokes phenomena associated with $\sigma_n^{\pm}$ occur

Borel resummation:

I. Two Borel singularities collide and line up along the vertical axis

II. Two Borel singularities collide with an angle $\pi/2$

III. Large-flavor expansion becomes Borel non-summable

Revisiting the SQED3

The 2nd order phase transition corresponds to followings

Lefschetz thimble analysis:

i. Contributing saddles jump as $\sigma_0^+ \to \sigma_0^+, \sigma_0^-$

ii. Two saddles collide with an angle $\pi/2$

iii. Infinite number of Stokes phenomena associa

The order of phase transition is decoded from "scattering angle"

Borel resummation:

I. Two Borel singularities collide and line up along the vertic l axis

II. Two Borel singularities collide with an angle $\pi/2$

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Revisiting the SQED3

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Due to SUSY

Borel resummation:

I. Two Borel singularities collide and line up alon

II. Two Borel singularities collide with an angle $\pi/2$

III. Large-flavor expansion becomes Borel non-summable

$$Z = \int_{-\infty}^{\infty} d\sigma \frac{e^{i\eta\sigma}}{\left[2\cosh\frac{\sigma+m}{2} \cdot 2\cosh\frac{\sigma-m}{2}\right]^{N_f}}$$

Revisiting the SQED3

The 2nd order phase transition corresponds to followings

Lefschetz thimble analysis:

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Borel resummation:

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II. Two Borel singularities collide with an angle $\pi/2$

III. Large-flavor expansion becomes Borel non-summable

They can be generalized as long as $e^{-NF(\lambda)} = \int \mathrm{d}\sigma \; e^{-N\tilde{S}(\lambda;\sigma)}$

Contents

Conclusion and future works

Question:

Is resurgence applicable to 2nd order phase transitions or more realistic QFTs?

Answer: resurgence is applicable!

2nd order phase transition = simultaneous Stokes and anti-Stokes phenomenon

- The order of phase transition is determined by a collision of saddles
- It is decoded from a perturbative series

$\longrightarrow$ Generalized to other systems

Future works:

- Relation to Lee-Yang zeros ?
- Expansion with respect to other parameters ?
- Physical meaning of the phase transition ?

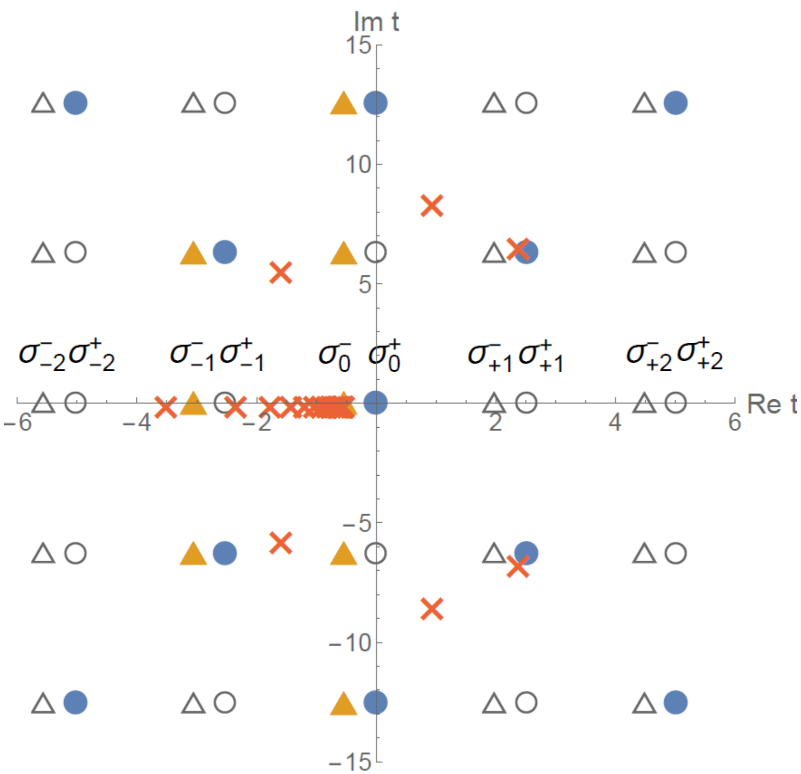
Backups

Borel singularities

- The Borel plane structure is consistent with the Lefschetz thimble structure
- Still there are artifacts, and missing singularities far from the origin

$\lambda < \lambda_c$

$\lambda = 0.4, m = 1,$ Padé (25,25)



$\lambda < \lambda_c$

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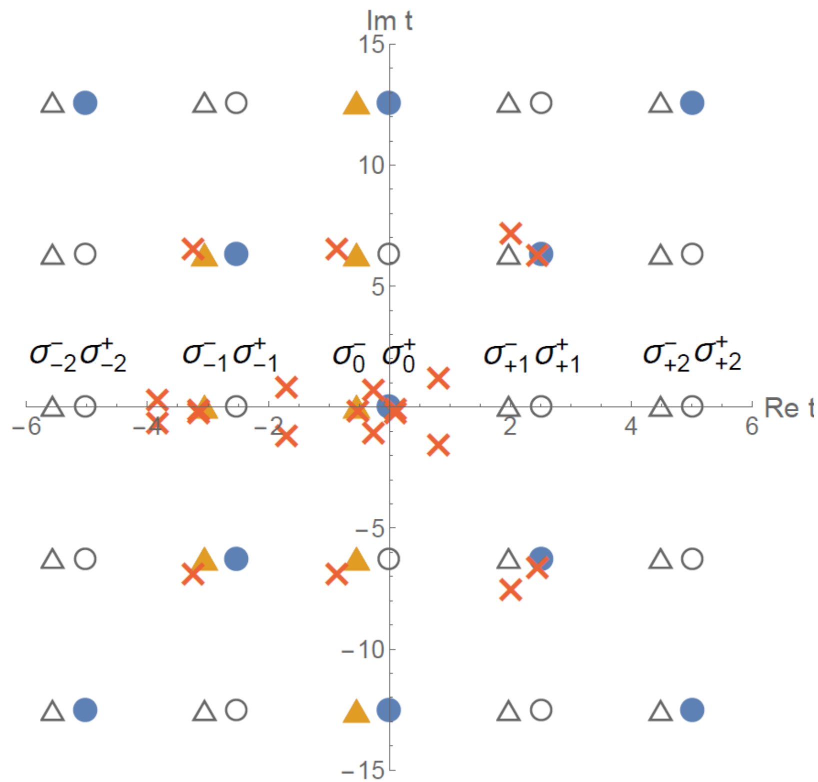

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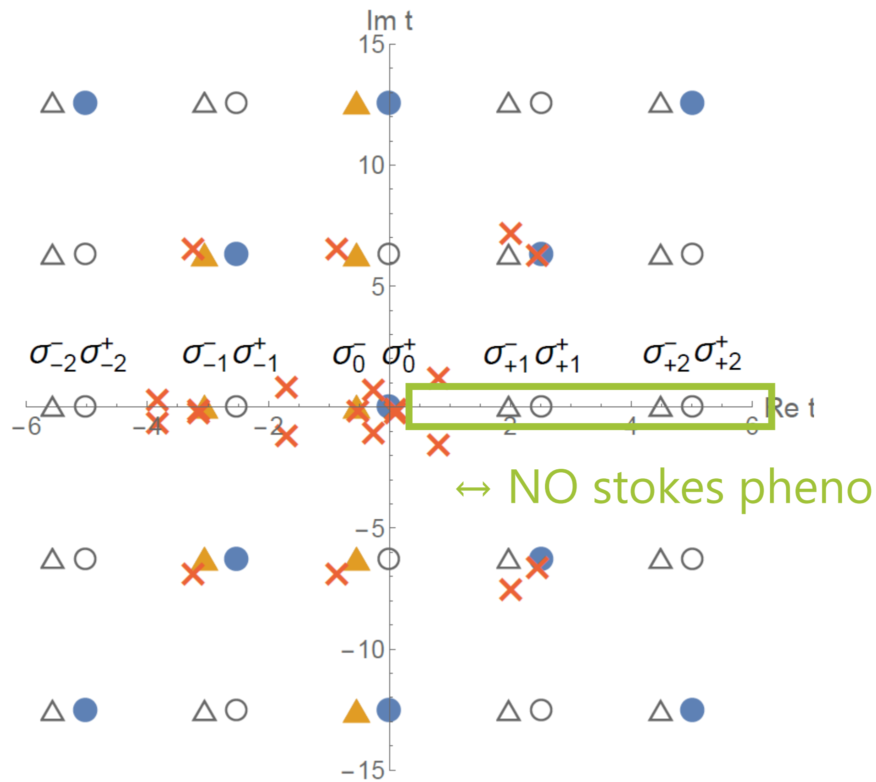

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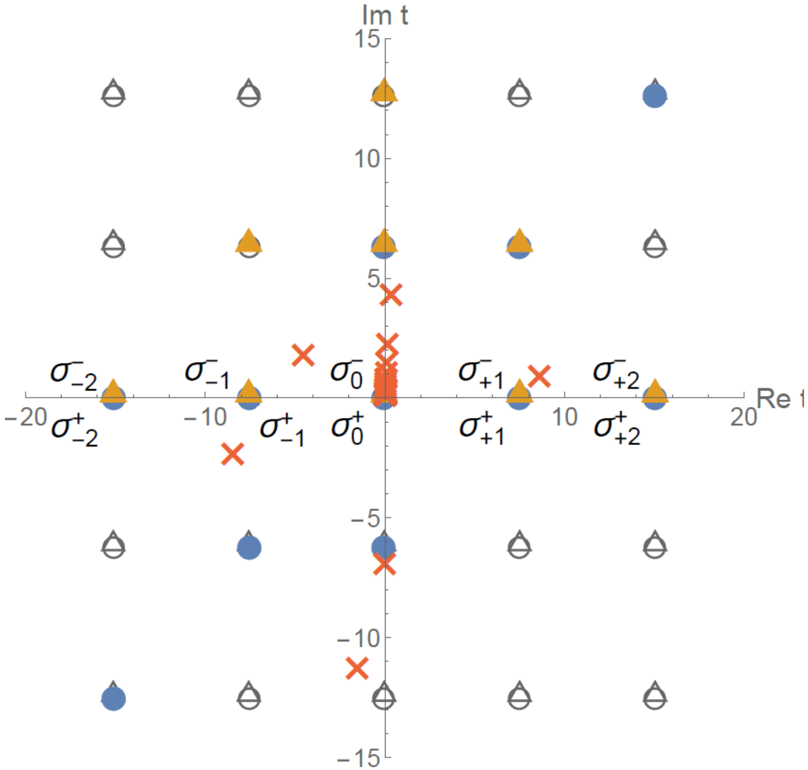

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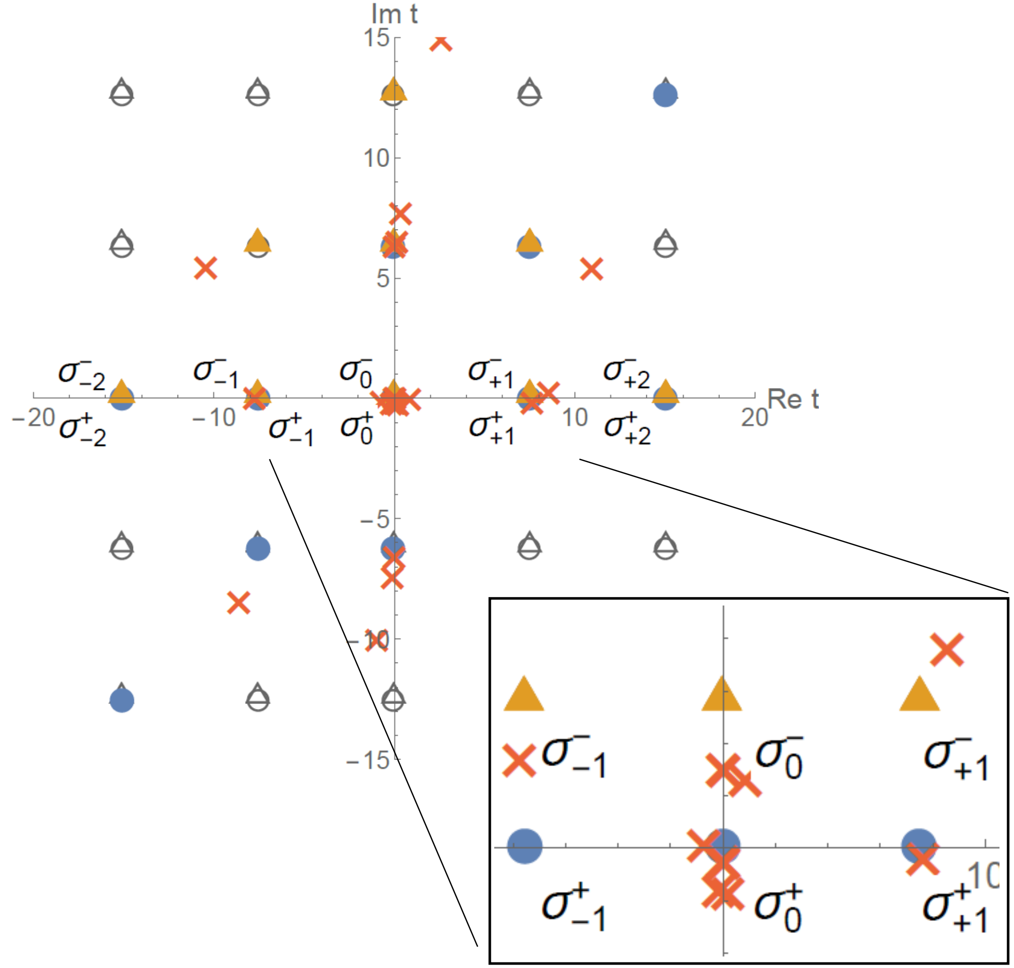
$\lambda = 1.2, m = 1$, Padé-Uniformized (25,25)



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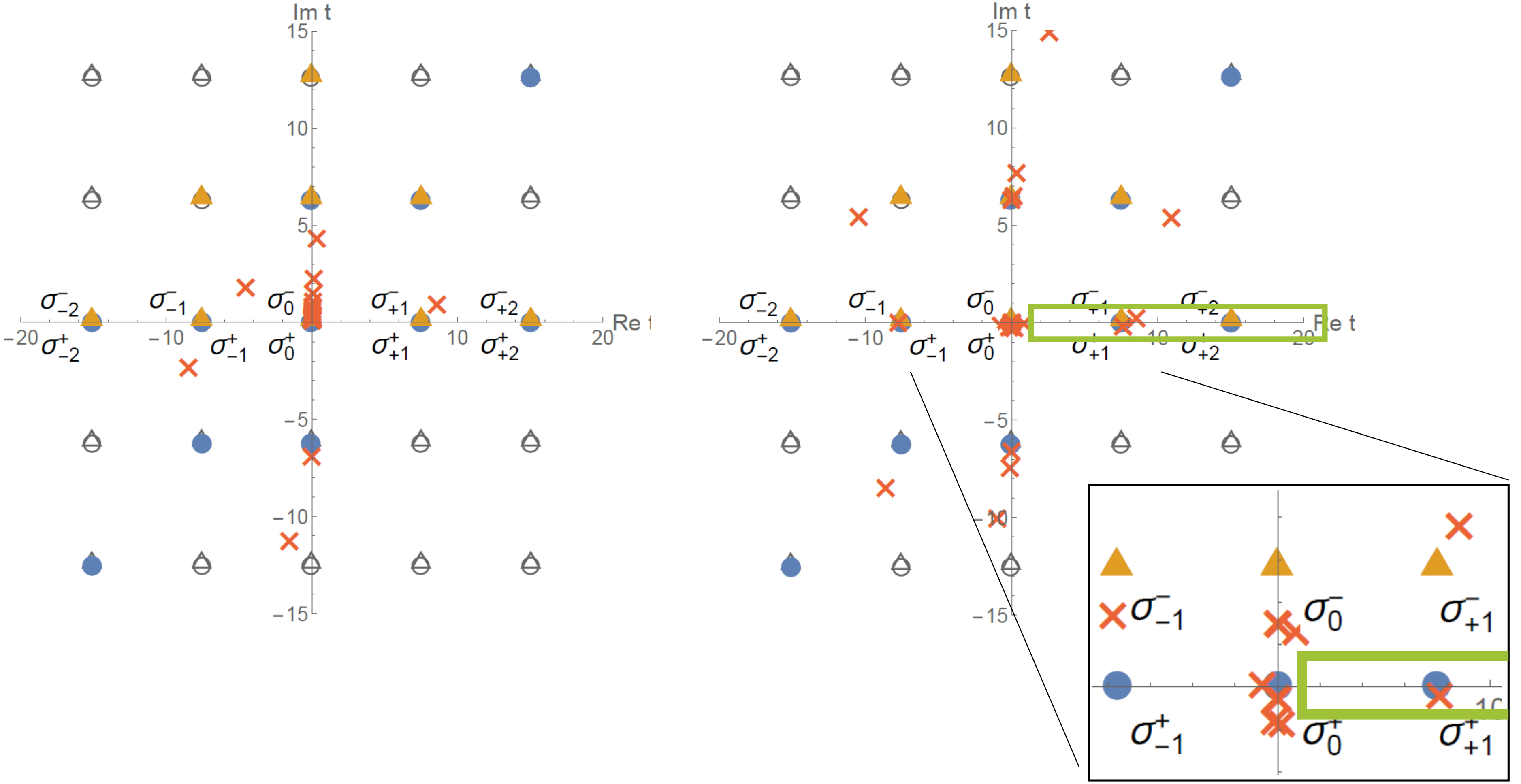
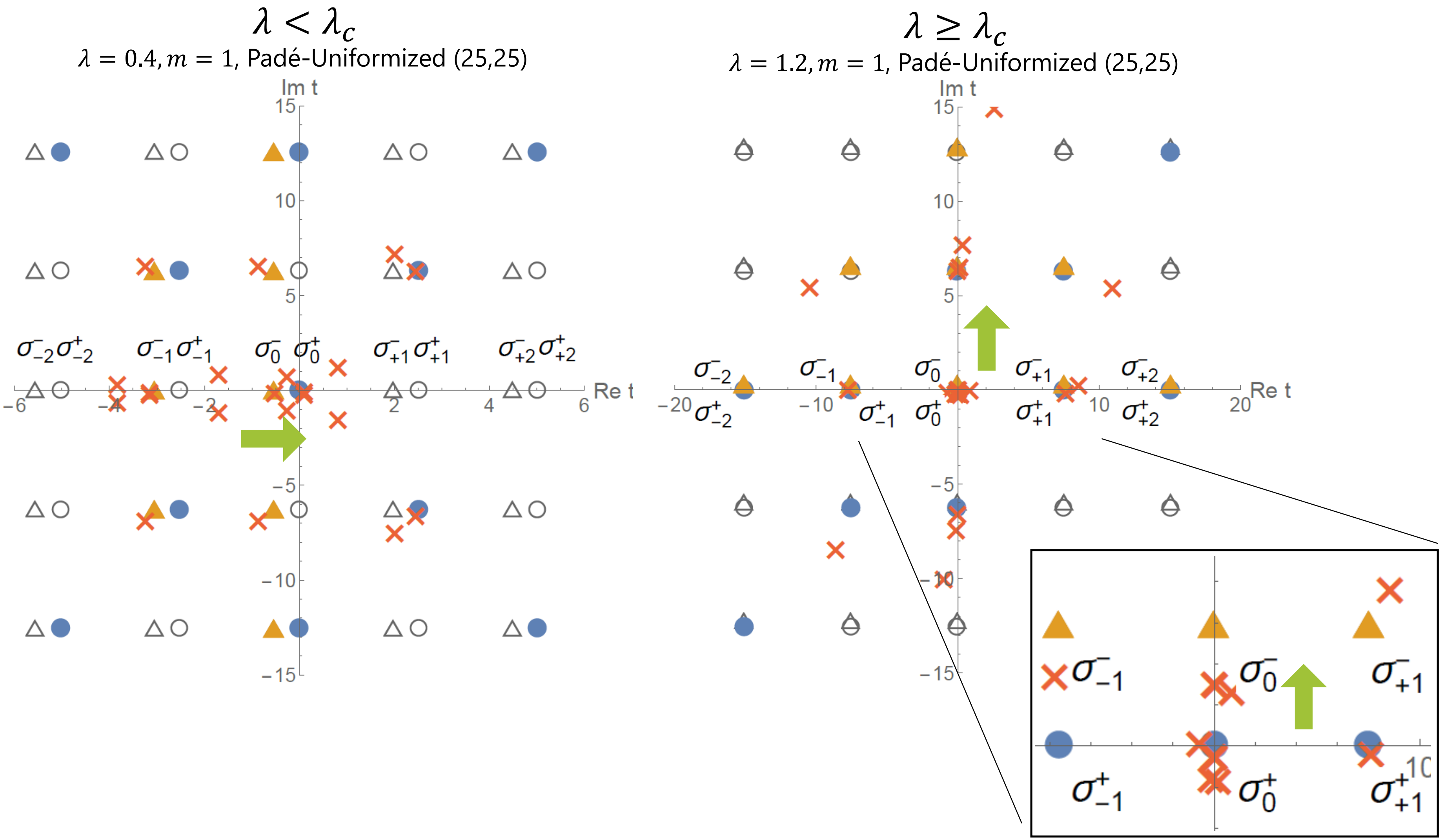


The order of the phase transition

2nd order phase transition corresponds to

collision of two saddles with the reflection angle $\pi/2$



Lefschetz thimble analysis

0dim Sine-Gordon model

[Cherman, Dorigoni, Unsal, 14]
[Cherman, Koroteev, Unsal, 14]

$$Z(g) = \frac{1}{(2\pi g)^{1/2}} \int_{-\pi/2}^{\pi/2} \mathrm{d}\phi \; e^{-S(\phi)/g}, \quad S(\phi) = \frac{1}{2}\sin^2\phi$$

Saddles and Lefschetz thimbles

$$0 = \frac{\mathrm{d}S(\phi)}{\mathrm{d}\phi} \Rightarrow \phi = 0, \pm\frac{\pi}{2}$$

Trivial saddle and non-trivial saddles

$$\mathcal{J}_i : \frac{\mathrm{d}\phi(t)}{\mathrm{d}t} = \overline{\frac{\mathrm{d}S}{\mathrm{d}\phi}}, \quad \phi(-\infty) = \phi_i \qquad \mathrm{Im}\, S(\phi(t)) = \mathrm{const.}, \quad \mathrm{Re}\, S(\phi_i) \leq \mathrm{Re}\, S(\phi(t))$$

$$\mathcal{K}_i : \frac{\mathrm{d}\phi(t)}{\mathrm{d}t} = -\overline{\frac{\mathrm{d}S}{\mathrm{d}\phi}}, \quad \phi(-\infty) = \phi_i \qquad \mathrm{Im}\, S(\phi(t)) = \mathrm{const.}, \quad \mathrm{Re}\, S(\phi_i) \geq \mathrm{Re}\, S(\phi(t))$$

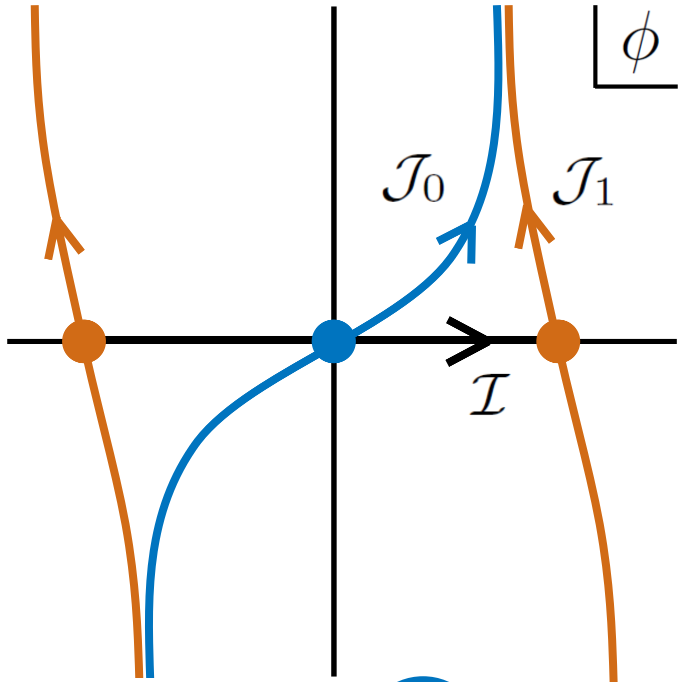
Lefschetz thimble analysis

Around $\arg g = 0$,

[Cherman, Dorigoni, Unsal, 14]
[Cherman, Koroteev, Unsal, 14]

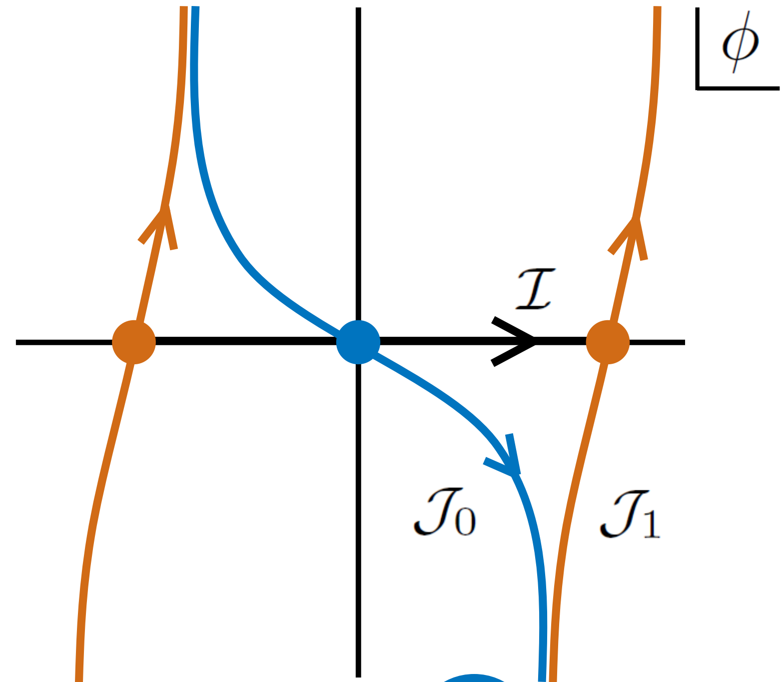
Stokes phenomenon associated with the trivial saddle

$\arg g = +0$



$$\mathcal{I} = \mathcal{J}_0 - \mathcal{J}_1$$

$\arg g = -0$



$$\mathcal{I} = \mathcal{J}_0 + \mathcal{J}_1$$

NO Stokes phenomenon associated with the non-trivial saddles

Borel resummation

[Cherman, Dorigoni, Unsal, 14]
[Cherman, Koroteev, Unsal, 14]

Perturbation theory around the trivial saddle diverges

$$Z(g) = \frac{1}{(2\pi g)^{1/2}} \int_{-\pi/2}^{\pi/2} \mathrm{d}\phi\ e^{-\frac{1}{2g}\sin^2\phi}$$

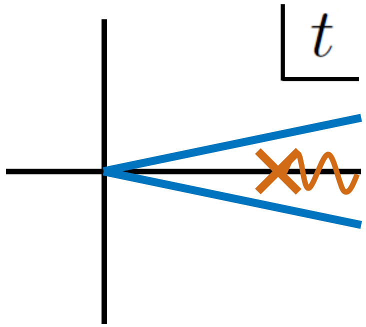
$$\overset{\text{around } \phi=0}{=} e^{-S(0)/g} \cdot \frac{1}{g} \sum_{n=0}^{\infty} \frac{(+2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1}$$

There is a Borel singularity (and a branch cut) around $\arg g = 0$

$$\mathcal{S}Z(g) = \int_C \mathrm{d}t\ e^{-t/g} \mathcal{B}Z(t)$$

$$= e^{-S(0)/g} \cdot \frac{1}{g} \int_C \mathrm{d}t\ e^{-t/g} \sum_{n=0} \frac{2^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)^2} (+t)^n$$

$$= e^{-S(0)/g} \cdot \frac{1}{g} \int_C \mathrm{d}t\ e^{-t/g}\, {}_2F_1\left(\frac{1}{2}, \frac{1}{2}, 1; +2t\right)$$



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$$Z(g) = \frac{1}{(2\pi g)^{1/2}} \int_{-\pi/2}^{\pi/2} \mathrm{d}\phi \; e^{-\frac{1}{2g}\sin^2\phi}$$

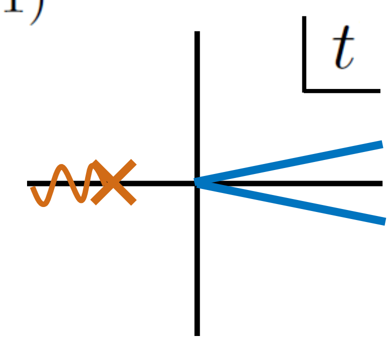
$$\overset{\text{around } \phi=\pi/2}{=} i e^{-S(\pi/2)/g} \cdot \frac{1}{g} \sum_{n=0}^{\infty} \frac{(-2)^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)} g^{n+1}$$

There is NO Borel singularity (nor branch cut) around $\arg g = 0$

$$\mathcal{S}Z(g) = \int_C \mathrm{d}t \; e^{-t/g} \mathcal{B}Z(t)$$

$$= i e^{-S(\pi/2)/g} \cdot \frac{1}{g} \int_C \mathrm{d}t \; e^{-t/g} \sum_{n=0} \frac{2^n \Gamma(n+1/2)^2}{\Gamma(1/2)^2 \Gamma(n+1)^2} (-t)^n$$

$$= i e^{-S(\pi/2)/g} \cdot \frac{1}{g} \int_C \mathrm{d}t \; e^{-t/g} {}_2F_1\left(\frac{1}{2}, \frac{1}{2}, 1; -2t\right)$$



Resurgence structure

[Cherman, Dorigoni, Unsal, 14]
[Cherman, Koroteev, Unsal, 14]

The two types of ambiguities cancel
and the location of the Borel singularity agrees with $S\left(\frac{\pi}{2}\right) = 1/2$

$$\begin{aligned}
\mathcal{S}Z(g) &= \underbrace{\mathcal{S}_{\pm}Z(g)|_{\text{around } \phi=0}} \mp \underbrace{\mathcal{S}Z(g)|_{\text{around } \phi=\pi/2}} \\
&= e^{-S(0)/g} \cdot \frac{1}{g}\int_{C^{\pm}} \mathrm{d}t \; e^{-t/g} {}_2F_1\left(\frac{1}{2}, \frac{1}{2}, 1; \boxed{+2t}\right) \\
&\qquad \mp i e^{-\boxed{S(\pi/2)}/g} \cdot \frac{1}{g}\int_{C} \mathrm{d}t \; e^{-t/g} {}_2F_1\left(\frac{1}{2}, \frac{1}{2}, 1; -2t\right) \\
&= \mathrm{Re}\, \mathcal{S}_{\pm}Z(g)|_{\text{around } \phi=0}
\end{aligned}$$

Information of non-trivial saddles is
encoded in perturbation theory around the trivial saddle

Large-flavor expansion

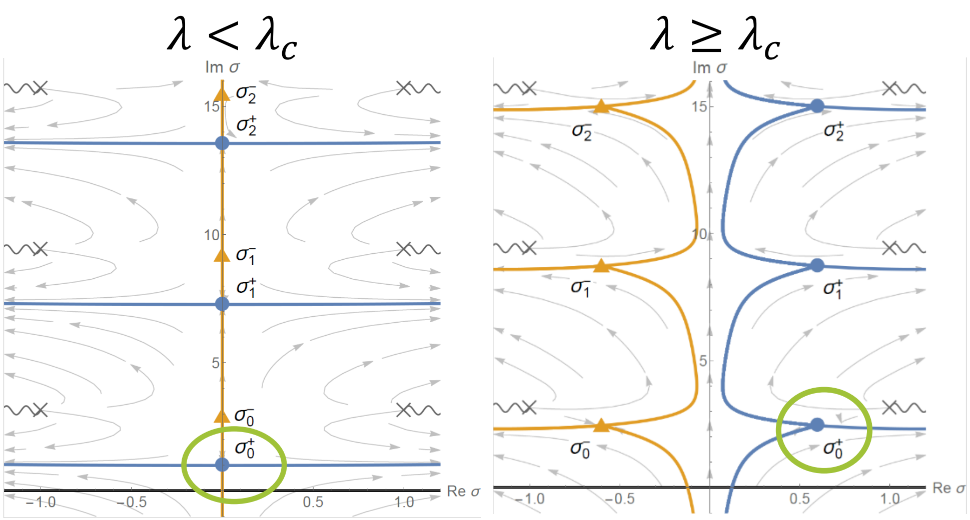
[Fujimori, Honda, Kamata, Misumi, Sakai, TY, to appear]

Consider the Borel resummation of $1/N$ expansion
to see how thimbles' structure is encoded

$$Z(\lambda;N)=\frac{1}{2^N}\int \mathrm{d}\sigma\ e^{-NS(\lambda;\sigma)},\quad S(\lambda;\sigma)=-i\lambda\sigma+\ln(\cosh\sigma+\cosh m)$$

$$\overset{\boxed{\text{around }\sigma_0^+}}{=}\frac{1}{2^N}\sqrt{\frac{2\pi}{NS''(\lambda;\sigma_0^+)}}e^{-NS(\lambda;\sigma_0^+)}\sum_{l=0}^{\infty}\frac{a_l(\lambda)}{N^l}$$

$$\mathcal{S}Z(\lambda;N)=\frac{1}{2^N}\sqrt{\frac{2\pi}{NS''(\lambda;\sigma_0^+)}}e^{-NS(\lambda;\sigma_0^+)}\cdot N\int_C \mathrm{d}t\ e^{-Nt}\underline{\sum_{l=0}^{\infty}\frac{a_l(\lambda)}{\Gamma(l+1)}t^l}$$



Does perturbation theory around the trivial saddle know non-trivial saddles and the phase transition?