

Electromagnetic Soliton Theory

an excerpt from

**A Novel Constructive Electromagnetic Quantum Theory
describes the Origin of Mass and Unifies the Forces.**

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Slide 2: What is a wave? (The d'Alembert wave equation)

Towne¹ states that the requirement for a physical condition to be referred to as a wave, is that its mathematical representation give rise to a partial differential equation of particular form, known as the wave equation. The classical form

$$\frac{\partial^2 w}{\partial p^2} - \frac{1}{u^2} \frac{\partial^2 w}{\partial t^2} = 0 \quad \text{or} \quad \nabla^2 w - \frac{1}{u^2} \frac{\partial^2 w}{\partial t^2} = 0.$$

was proposed in 1748 by d'Alembert for a one-dimensional continuum. A decade later, Euler established the equation for the three-dimensional continuum.

A pendulum is described by the pendulum equation

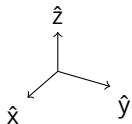
$$\frac{\partial^2 \theta}{\partial t^2} + \frac{g}{l} = 0$$

is not a wave and cannot be described as a soliton even if it is Lorentz boosted, *e. g.* taking the pendulum on a journey in an aeroplane. A pendulum equation does not, and will never describe *displacement* motion.

1 Dudley H. Towne. *Wave phenomena*. New York: Dover Publications, 1988.

Slide 3: Three orthogonal vectors in an Euclidean reference system

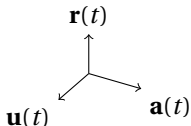
The reference system whose axis are the unit vectors \hat{x}, \hat{y} and \hat{z}



is defined by:

$$\begin{aligned} \hat{x} \cdot \hat{y} &= 0 & \hat{y} \cdot \hat{z} &= 0 & \hat{z} \cdot \hat{x} &= 0 \\ \hat{x} \times \hat{y} &= \hat{z} & \hat{y} \times \hat{z} &= \hat{x} & \hat{z} \times \hat{x} &= \hat{y} \end{aligned}$$

Consider three orthogonal vectors of function of time



gives:

$$\begin{aligned} \mathbf{u} \cdot \mathbf{a} &= 0 & \mathbf{a} \cdot \mathbf{r} &= 0 & \mathbf{r} \cdot \mathbf{u} &= 0 \\ \mathbf{u} \times \mathbf{a} &= \mathbf{r} & \mathbf{a} \times \mathbf{r} &= \mathbf{u} & \mathbf{r} \times \mathbf{u} &= \mathbf{a} \end{aligned}$$

Next consider the indefinite series

$$\mathbf{z}_1 = \mathbf{u}_0 \times \mathbf{a}_0, \quad \mathbf{u}_1 = \mathbf{a}_0 \times \mathbf{r}_1, \quad \mathbf{a}_1 = \mathbf{r}_1 \times \mathbf{u}_1, \quad \mathbf{r}_2 = \mathbf{u}_1 \times \mathbf{a}_1 \quad \dots \quad \mathbf{r}_n = \mathbf{u}_{n-1} \times \mathbf{a}_{n-1} \quad \dots \quad \dots$$

and what needs to be done so that $\mathbf{u}_0 = \mathbf{u}_n$, $\mathbf{a}_0 = \mathbf{a}_n$, $\mathbf{r}_0 = \mathbf{r}_n$ to give us a simultaneous vector cross product equation set which has defined solutions?

Slide 4: Theorem: The Soliton Equation System

We introduce normalisation: $\mathbf{u} = \frac{1}{\|\mathbf{a}\|^2} \mathbf{a} \times \mathbf{r}$, $\mathbf{a} = \frac{1}{\|\mathbf{u}\|^2} \mathbf{r} \times \mathbf{u}$, $\mathbf{r} = \mathbf{u} \times \mathbf{a}$

Theorem 1: The soliton equation system. *In a space \mathbb{C}^3 the system of simultaneous equations*

$$\mathcal{M}(\mathbf{u}, \mathbf{a}, \mathbf{r}) \xrightarrow{\text{defines}} \left\{ \mathbf{u} = \frac{1}{\mathbf{a} \cdot \mathbf{a}^*} \mathbf{a} \times \mathbf{r}, \quad \mathbf{a} = \frac{1}{\mathbf{u} \cdot \mathbf{u}^*} \mathbf{r} \times \mathbf{u}, \quad \mathbf{r} = \mathbf{u} \times \mathbf{a} \right\}$$

defines the motion of a soliton characterised by a velocity vector $\mathbf{u}(t)$ and two co-orthogonal vectors $\mathbf{a}(t)$ and $\mathbf{r}(t)$ that describe the disturbance in a homogenous and isotropic medium.

Here the vector quantities \mathbf{u} , \mathbf{a} and \mathbf{r} are complex vectors, for example $\mathbf{a} = \hat{x} a_x e^{i\alpha_x} + \hat{y} a_y e^{i\alpha_y} + \hat{z} a_z e^{i\alpha_z}$ $\mathbf{a}^* = \hat{x} a_x e^{-i\alpha_x} + \hat{y} a_y e^{-i\alpha_y} + \hat{z} a_z e^{-i\alpha_z}$ therefore $\mathbf{a} \cdot \mathbf{a}^* = a_x^2 + a_y^2 + a_z^2 = a^2 = \|\mathbf{a}\|^2$

Slide 5: Proof: $\{\mathbf{u} = \mathbf{a} \times \mathbf{r} / \mathbf{a} \cdot \mathbf{a}^*, \quad \mathbf{a} = \mathbf{r} \times \mathbf{u} / \mathbf{u} \cdot \mathbf{u}^*, \quad \mathbf{r} = \mathbf{u} \times \mathbf{a}\}$ describes a soliton.

Performing a 'left and right side' curl operation on the second and third equations of the equation-set gives

$$\nabla \times \mathbf{a} = \frac{1}{\mathbf{u} \cdot \mathbf{u}^*} \nabla \times (\mathbf{r} \times \mathbf{u}) \quad \text{and} \quad \nabla \times \mathbf{r} = \nabla \times (\mathbf{u} \times \mathbf{a}) \quad (1)$$

and to evaluate the vector triple products we use general vector analytic methods to give

$$\begin{aligned} \nabla \times (\mathbf{r} \times \mathbf{u}) &= \mathbf{r}(\nabla \cdot \mathbf{u}) - \mathbf{u}(\nabla \cdot \mathbf{r}) + (\mathbf{u} \cdot \nabla)\mathbf{r} - (\mathbf{r} \cdot \nabla)\mathbf{u} \\ \nabla \times (\mathbf{u} \times \mathbf{a}) &= \mathbf{u}(\nabla \cdot \mathbf{a}) - \mathbf{a}(\nabla \cdot \mathbf{u}) + (\mathbf{a} \cdot \nabla)\mathbf{u} - (\mathbf{u} \cdot \nabla)\mathbf{a}. \end{aligned}$$

Because the vectors \mathbf{a} and \mathbf{r} are position independent (from theorem: **vectors $\mathbf{a}(t)$ and $\mathbf{r}(t)$ describe the disturbance in a homogenous and isotropic medium.**, therefore we have

$$\nabla \cdot \mathbf{a} = 0 \quad \text{and} \quad \nabla \cdot \mathbf{r} = 0. \quad (2)$$

Slide 6: Proof: $\{\mathbf{u} = \mathbf{a} \times \mathbf{r} / \mathbf{a} \cdot \mathbf{a}^*, \quad \mathbf{a} = \mathbf{r} \times \mathbf{u} / \mathbf{u} \cdot \mathbf{u}^*, \quad \mathbf{r} = \mathbf{u} \times \mathbf{a}\}$ describes a soliton.

$$\nabla \times (\mathbf{r} \times \mathbf{u}) = \mathbf{r}(\nabla \cdot \mathbf{u}) - \cancel{\mathbf{u}(\nabla \cdot \mathbf{r})} + (\mathbf{u} \cdot \nabla)\mathbf{r} - (\mathbf{r} \cdot \nabla)\mathbf{u}$$

$$\nabla \times (\mathbf{u} \times \mathbf{a}) = \cancel{\mathbf{u}(\nabla \cdot \mathbf{a})} - \mathbf{a}(\nabla \cdot \mathbf{u}) + (\mathbf{a} \cdot \nabla)\mathbf{u} - (\mathbf{u} \cdot \nabla)\mathbf{a}.$$

Evaluating the terms containing $\mathbf{u} = \hat{x} \partial x / \partial t + \hat{y} \partial y / \partial t + \hat{z} \partial z / \partial t$ we obtain

$$\mathbf{u} \cdot \nabla = \nabla \cdot \mathbf{u} = \frac{\partial x}{\partial t} \frac{\partial}{\partial x} + \frac{\partial y}{\partial t} \frac{\partial}{\partial y} + \frac{\partial z}{\partial t} \frac{\partial}{\partial z} = \frac{\partial}{\partial t}$$

Because $\mathbf{a}(\mathbf{u} \cdot \nabla) = \mathbf{a} \partial 1 / \partial t = 0$, we are left with

$$\nabla \times (\mathbf{u} \times \mathbf{a}) = -\frac{\partial \mathbf{a}}{\partial t} \quad \text{and} \quad \nabla \times (\mathbf{r} \times \mathbf{u}) = \frac{\partial \mathbf{r}}{\partial t}.$$

Therefore, the 'left and right side' curl operations (1) generate the new relations:

$$\nabla \times \mathbf{a} = \frac{1}{u^2} \frac{\partial \mathbf{r}}{\partial t} \quad \text{and} \quad \nabla \times \mathbf{r} = -\frac{\partial \mathbf{a}}{\partial t} \quad (3)$$

Slide 7: Proof: $\{\mathbf{u} = \mathbf{a} \times \mathbf{r} / \mathbf{a} \cdot \mathbf{a}^*, \quad \mathbf{a} = \mathbf{r} \times \mathbf{u} / \mathbf{u} \cdot \mathbf{u}^*, \quad \mathbf{r} = \mathbf{u} \times \mathbf{a}\}$ describes a soliton.

Therefore, the 'left and right side' curl operations generate the new relations:

$$\nabla \times \mathbf{a} = \frac{1}{u^2} \frac{\partial \mathbf{r}}{\partial t} \quad \text{and} \quad \nabla \times \mathbf{r} = -\frac{\partial \mathbf{a}}{\partial t} \quad (3)$$

A further 'left and right side' curl operation on (3) gives

$$\nabla \times \nabla \times \mathbf{r} = -\frac{\partial(\nabla \times \mathbf{a})}{\partial t} \quad \text{and} \quad \nabla \times \nabla \times \mathbf{a} = \frac{1}{u^2} \frac{\partial(\nabla \times \mathbf{r})}{\partial t}$$

and because $\nabla \times \nabla \times \mathbf{r} = \nabla(\nabla \cdot \mathbf{r}) - \nabla^2 \mathbf{r}$ we recover the d'Alembert wave equations

$$\nabla^2 \mathbf{r} - \frac{1}{u^2} \frac{\partial^2 \mathbf{r}}{\partial t^2} = 0 \quad \text{and} \quad \nabla^2 \mathbf{a} - \frac{1}{u^2} \frac{\partial^2 \mathbf{a}}{\partial t^2} = 0. \quad (4)$$

This concludes the proof that the three vector algebraic equations of \mathcal{M} give rise to the d'Alembert wave equations (4). Therefore, the equation set \mathcal{M} is a generic bimodal-transverse soliton equation system. \square

Slide 8: One Plane of an Electromagnetic Wave

$\mathcal{M}(\mathbf{u}, \mathbf{B}, \mathbf{E})$ recovers the Maxwell equations in vacuum

from (2):

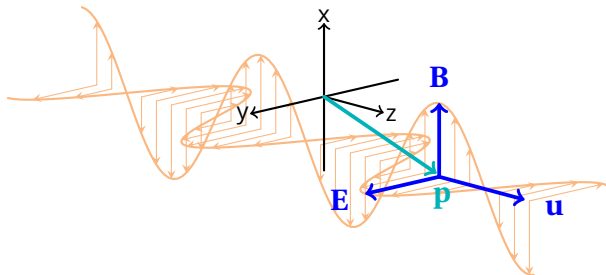
$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \cdot \mathbf{E} = 0$$

from (3):

$$\nabla \times \mathbf{B} = \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$



$$\left\{ \mathbf{E} = \mathbf{u} \times \mathbf{B}, \quad \mathbf{u} = \frac{1}{\|\mathbf{B}\|^2} \mathbf{B} \times \mathbf{E}, \quad \mathbf{B} = \frac{1}{\|\mathbf{u}\|^2} \mathbf{E} \times \mathbf{u} \right\} \quad \mathbf{p} = \int \mathbf{u} dt$$

Remark 1: Electrostatic *versus* Electromotive. I intentionally distinguish between

- the **elementary electrostatic charge** e (represented by e in italic serif font) and
- the **electromotive charge** e (represented by e in script font)

In the conventional interpretation of electromagnetic theory, both charges are considered equal in magnitude and are measured in units of coulombs.

Remark 2: Defining the *emflux*. As a consequence of above, we must distinguish between

- the **magnetic flux** ϕ (slanted ϕ) and
- the **elementary magnetic momentum** ϕ (upright ϕ)

The elementary magnetic momentum ϕ or its vector form Φ , which will henceforth be referred to as **a magnetic emflux (magnetic electromotive flux)**

Slide 10: Theorem: Elementary Electromagnetic Soliton

The equation system \mathcal{M} defines a purely mathematical and a purely generical Maxwell equation system allowing us to formulate new Maxwellian like system.

Theorem 2: Elementary Electromagnetic Soliton. *There exists an elementary length denoted as l_o and an elementary time denoted as t_o defining the speed of light c such that $l_o = ct_o$. Additionally, an elementary magnetic emflux ϕ_o represents a quantum of magnetic momentum. Furthermore, an elementary EM-soliton, defined by $\mathcal{M}(\mathbf{u}, \phi_o, \Upsilon_o)$, carries an elementary electromotive charge denoted as e and has action h , whilst propagating at the speed of light.*

Here Υ_o is the electric electromotive flux.

Slide 11: Proof: Elementary Electromagnetic Soliton

The proof is structured by demonstrating that the set of simultaneous equations

$$\mathcal{M}(\mathbf{u}, \boldsymbol{\phi}_0, \boldsymbol{\Upsilon}_0) \xrightarrow{\text{defines}} \left\{ \mathbf{u} = \frac{1}{\|\boldsymbol{\phi}_0\|^2} \boldsymbol{\phi}_0 \times \boldsymbol{\Upsilon}_0, \boldsymbol{\phi}_0 = \frac{1}{\|\mathbf{u}\|^2} \boldsymbol{\Upsilon}_0 \times \mathbf{u}, \boldsymbol{\Upsilon}_0 = \mathbf{u} \times \boldsymbol{\phi}_0 \right\}$$

together with the theorem's assertions demands the presence of ϵ_0 and μ_0 in their known forms.

Assuming that $\boldsymbol{\phi}_0 \times \boldsymbol{\Upsilon}_0$ represents wave action, we multiply the equation $\mathbf{u} = (\boldsymbol{\phi}_0 \times \boldsymbol{\Upsilon}_0) \|\boldsymbol{\phi}_0\|^{-2}$ by h and evaluate its norm, yielding

$$\|h\mathbf{u}\| = \left\| \frac{h}{\|\boldsymbol{\phi}_0\|^2} \boldsymbol{\phi}_0 \times \boldsymbol{\Upsilon}_0 \right\| \quad \text{to give}$$

$$h = \left[\frac{h}{c\phi_0^2} \right] (\|\boldsymbol{\phi}_0\| \|\boldsymbol{\Upsilon}_0\|) \quad (5)$$

where the square brackets indicate the development of a constant.

Slide 12: Proof: Electromagnetic action

- We define the elementary EM-action as $h_e = \rho h$, where $\rho = 1 \text{ C/kg}$
- Theorem 2 demands that e is transported at a velocity c
- Action is momentum times distance, we consider the elementary distance l_o

therefore

$$h_e = \rho h = \kappa e c l_o \quad (6)$$

where κ is a dimensionless proportionality constant of unknown value, scaling $e c l_o$ to the EM-action h_e .

Also, Theorem 2 states that ϕ represents a quantum of magnetic momentum. Consequently, EM-action is also proportional to the product of ϕ and the distance travelled:

$$\rho h = \chi \|\Phi_o\| l_o \quad \text{combining with (6)} \quad \|\Phi_o\| = \kappa e c \chi^{-1} \quad (7)$$

where χ is a physical quantity with units and scaling to be determined.

Slide 13: Proof: Permittivity and Permeability

Repeating equation (5): $h = [h/c\phi_0^2] (\|\phi_0\| \|\gamma_0\|)$ which we now rewrite, using $\|\phi_0\| = \kappa e c \chi^{-1}$ and the relationship $\gamma_0 = \mathbf{u} \times \phi_0$ (or $\|\gamma_0\| = c \|\phi_0\|$) giving:

$$h = \left[\frac{h}{c\phi_0^2} \right] \left[\frac{1}{\chi} \right] c^2 \kappa e \phi_0$$

We are now in a position to define the expression for $\phi_0 = \frac{h}{\kappa e}$ but only if

$$1 = \left[\frac{h}{c\phi_0^2} \right] \left[\frac{1}{\chi} \right] c^2 \quad \text{and replacing } \phi_0 \text{ gives}$$

$$1 = \left[\frac{\kappa^2 e^2}{ch} \right] \left[\frac{1}{\chi} \right] c^2 \quad \text{which requires } \frac{1}{\chi} = \frac{h}{\kappa^2 e^2 c}, \text{ hence}$$

$$1 = \left[\frac{\kappa^2 e^2}{ch} \right] \left[\frac{h}{\kappa^2 e^2 c} \right] c^2 \tag{8}$$

Slide 14: Proof: Permittivity and Permeability

Equation (8), that is $1 = \left[\frac{\kappa^2 e^2}{ch} \right] \left[\frac{h}{\kappa^2 e^2 c} \right] c^2$ defines, purely mathematical the permittivity and permeability of the medium. These are the two constants developed mathematically and enclosed in the square brackets.

$$\epsilon_0 = \frac{\kappa^2 e^2}{hc} \quad \text{and} \quad \mu_0 = \frac{h}{\kappa^2 e^2 c}$$

and mapping $e \mapsto e$ and setting $\kappa^{-2} = 2\alpha$ gives the accustomed

$$\epsilon_0 = \frac{e^2}{2\alpha hc} \quad \text{and} \quad \mu_0 = \frac{2\alpha h}{e^2 c}$$

Slide 15: The Mass Gap

Recalling $\rho h = \kappa e l_0 c$, *i.e.* (6); we are now in the position to calculate the numeric values for the elementary length and time, using $\kappa^{-2} = 2\alpha$, $\rho = 1 \text{ C/kg}$, and the 2018 CODATA values:

$\kappa = 8.277\,559\,999\,29(62)$	which I name the Heaviside constant
$l_0 = 1.666\,566\,299\,11(12) \times 10^{-24}$	elementary length in metres
$t_0 = 5.559\,066\,796\,49(42) \times 10^{-33}$	elementary time in seconds
$\Delta_0 = 3.683\,476\,656\,21(18) \times 10^{-66}$	mass gap in joules

where $\Delta_0 = h t_0$ is the least energy gap from a vacuum to the next lowest energy state.

Historic note: In the late 19th century Oliver Heaviside developed vector calculus, and rewrote the Maxwell works into the form commonly used today. The Heaviside constant κ is a coupling constant relating the electric charge momentum to mechanical momentum.

Slide 16: Syntax for describing solutions of \mathcal{M}

Introducing a new mathematical syntax, utilising a row-by-row matrix product operator \diamond , defined as follows:

$$\begin{pmatrix} Pa_{1,1} & Pa_{1,2} \\ Qa_{2,1} & Qa_{2,2} \end{pmatrix} = \begin{pmatrix} P \\ Q \end{pmatrix} \diamond \begin{pmatrix} a_{1,1} & a_{1,2} \\ a_{2,1} & a_{2,2} \end{pmatrix}$$

A wave or soliton ξ that is a solution of \mathcal{M} is precisely defined by the three vectors \mathbf{u} , Φ_0 , and Υ_0 , expressed in matrix form as

$$\xi \stackrel{\text{def}}{\text{by}} \begin{pmatrix} \mathbf{u} \\ \Phi_0 \\ \Upsilon_0 \end{pmatrix} = \begin{pmatrix} c \\ \Phi_0 \\ c\Phi_0 \end{pmatrix} \diamond \begin{pmatrix} a_{1,1} & a_{1,2} & a_{1,3} \\ a_{2,1} & a_{2,2} & a_{2,3} \\ a_{3,1} & a_{3,2} & a_{3,3} \end{pmatrix} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

This expression can be further simplified by considering only the parameters of interest:

$$\xi \stackrel{\text{par}}{\text{by}} \begin{pmatrix} c \\ \Phi_0 \\ \Upsilon_0 \end{pmatrix} \diamond \begin{pmatrix} a_{1,1} & a_{1,2} & a_{1,3} \\ a_{2,1} & a_{2,2} & a_{2,3} \\ a_{3,1} & a_{3,2} & a_{3,3} \end{pmatrix}$$

Slide 17: Solutions of $\mathcal{M}(\mathbf{u}, \mathbf{a}, \mathbf{r}) \xrightarrow{\text{defines}} \{ \mathbf{u} = \mathbf{a} \times \mathbf{r} / \mathbf{a} \cdot \mathbf{a}^*, \quad \mathbf{a} = \mathbf{r} \times \mathbf{u} / \mathbf{u} \cdot \mathbf{u}^*, \quad \mathbf{r} = \mathbf{u} \times \mathbf{a} \}$

Electromagnetic solitons of interest are described generically:

$$\begin{array}{l}
 \xi_1 \xrightarrow{\text{par by}} \left(\begin{array}{c} c \\ \phi_0 \\ \Upsilon_0 \end{array} \right) \diamond \left(\begin{array}{ccc} \cos \theta & \sin \theta & 0 \\ -\sin \theta \sin \vartheta & \cos \theta \sin \vartheta & \cos \vartheta \\ \sin \theta \cos \vartheta & -\cos \theta \cos \vartheta & \sin \vartheta \end{array} \right) \\
 \text{or} \\
 \xi_2 \xrightarrow{\text{par by}} \left(\begin{array}{c} c \\ \phi_0 \\ \Upsilon_0 \end{array} \right) \diamond \left(\begin{array}{ccc} \sin \theta \cos \vartheta & -\cos \theta \cos \vartheta & \sin \vartheta \\ \cos \theta & \sin \theta & 0 \\ -\sin \theta \sin \vartheta & \cos \theta \sin \vartheta & \cos \vartheta \end{array} \right) \\
 \text{or} \\
 \xi_3 \xrightarrow{\text{par by}} \left(\begin{array}{c} c \\ \phi_0 \\ \Upsilon_0 \end{array} \right) \diamond \left(\begin{array}{ccc} -\sin \theta \sin \vartheta & \cos \theta \sin \vartheta & \cos \vartheta \\ \sin \theta \cos \vartheta & -\cos \theta \cos \vartheta & \sin \vartheta \\ \cos \theta & \sin \theta & 0 \end{array} \right)
 \end{array}
 \left. \vphantom{\begin{array}{l} \xi_1 \\ \xi_2 \\ \xi_3 \end{array}} \right\} \text{where} \left\{ \begin{array}{l} \theta = sn\omega_0 t \\ \vartheta = r_2 m \omega_0 t \\ s_a \in \{1/2, 1, 3/2, \dots\} \\ r_a \in \{-1, 0, 1\} \\ s = s_a r_a \\ r_2 \in \{-1, 0, 1\} \\ n \in \mathbb{Q} \geq 0 \\ m \in \mathbb{Q} \geq 0 \end{array} \right.$$

Slide 18: Rotating flux vectors

We are working with rotating emflux vectors. For instance, $\boldsymbol{\phi}_0$ represents a rotating vector, which we define as the source of a north-pointing elementary magnetic emflux, denoted as $\phi_0 = l_0^2 \boldsymbol{\phi}_0$. Consequently, $-\boldsymbol{\phi}_0$ still acts as a source of a north-pointing emflux but in the opposite direction. We are now required to introduce $\bar{\boldsymbol{\phi}}_0$ as the magnetic field vector that absorbs a north-pointing emflux. This implies that $\boldsymbol{\phi}_0 + \bar{\boldsymbol{\phi}}_0 \equiv 0$, and $\boldsymbol{\phi}_0 - \bar{\boldsymbol{\phi}}_0 \equiv 2\boldsymbol{\phi}_0$ if and only if $\boldsymbol{\phi}_0 = \hat{p}\phi_0$ and $\bar{\boldsymbol{\phi}}_0 = \hat{p}\bar{\phi}_0$, where \hat{p} represents any unit vector. Below is a visual representation of this concept where the symbol \textcircled{S} signifies the source or the sink:

$$\begin{aligned} \boldsymbol{\phi}_0 &\mapsto \text{S}\textcircled{S} \longrightarrow \text{N} & \text{and} & & -\boldsymbol{\phi}_0 &\mapsto \text{N} \longleftarrow \textcircled{S}\text{S} \\ \bar{\boldsymbol{\phi}}_0 &\mapsto \text{N}\textcircled{S} \longleftarrow \text{S} & \text{and} & & -\bar{\boldsymbol{\phi}}_0 &\mapsto \text{S} \longrightarrow \textcircled{S}\text{N} \end{aligned}$$

Slide 19: Elementary EM-solitons, the *emtron*

$$m_o \xrightarrow[\text{by}]{\text{par}} \begin{pmatrix} c \\ \phi_o \\ \Upsilon_o \end{pmatrix} \diamond \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \text{and} \quad \bar{m}_o \xrightarrow[\text{by}]{\text{par}} \begin{pmatrix} c \\ \bar{\phi}_o \\ \bar{\Upsilon}_o \end{pmatrix} \diamond \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\dot{m}_o \xrightarrow[\text{by}]{\text{par}} \begin{pmatrix} c \\ \phi_o \\ \Upsilon_o \end{pmatrix} \diamond \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad \text{and} \quad \dot{\bar{m}}_o \xrightarrow[\text{by}]{\text{par}} \begin{pmatrix} c \\ \bar{\phi}_o \\ \bar{\Upsilon}_o \end{pmatrix} \diamond \begin{pmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

The overaccented bar represents an 'anti' (where the source becomes an absorber), while the underaccented dot signifies 'contra' (indicating a 180-degree rotation). At creation the recoil reactions are:

		\mathcal{E}	p	ϕ_o	Υ_o
a)	$m_o + \bar{m}_o + \dot{m}_o + \dot{\bar{m}}_o \Rightarrow$	4	0	0	0
b)	$m_o + \bar{m}_o \Rightarrow$	2	0	0	2
c)	$m_o + \dot{\bar{m}}_o \Rightarrow$	2	0	2	0
d)	$m_o + \dot{m}_o \Rightarrow$	2	2	0	0
e)	$\bar{m}_o + \dot{\bar{m}}_o \Rightarrow$	2	-2	0	0

Slide 20: *Emtrons* in circular self-orbits: Spin=0

The notable finding is that the solutions to \mathcal{M} permit circular self-orbits. As proven by Theorem 1, \mathcal{M} leads to the Maxwell equations. Thus, circular and spherular self-orbits are intrinsic features of electromagnetic phenomena. Circular self-orbits are mathematically described by the following equations:

$$m_o^\odot \xrightarrow{\text{def by}} \begin{pmatrix} c \\ \Phi_o \\ \Upsilon_o \end{pmatrix} \diamond \begin{pmatrix} \cos \omega_o t & \sin \omega_o t & 0 \\ -\sin \omega_o t & \cos \omega_o t & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

$$m_o^\odot \xrightarrow{\text{def by}} \begin{pmatrix} c \\ \Phi_o \\ \Upsilon_o \end{pmatrix} \diamond \begin{pmatrix} \cos \omega_o t & \sin \omega_o t & 0 \\ 0 & 0 & 1 \\ \sin \omega_o t & -\cos \omega_o t & 0 \end{pmatrix} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

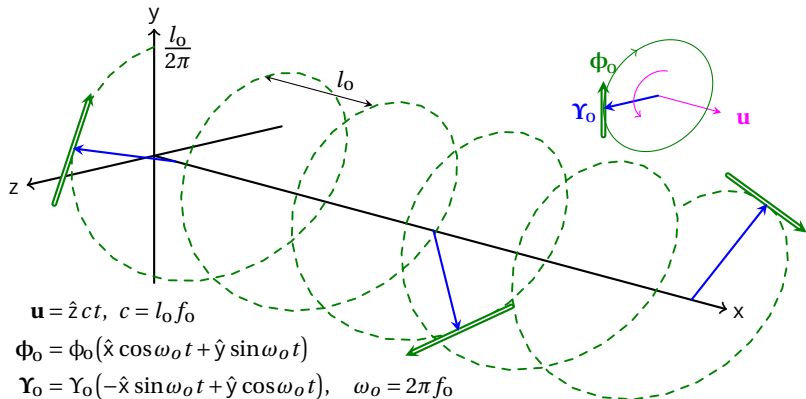
These emtrons are classified as spin-zero because either Φ_o or Υ_o remains static.

These would be responsible for establishing the electromotive field between capacitor plates, or the electromotive magnetic fields of permanent magnets.

Slide 21: *Emtrons* with Rotation and Linear Motion; Spin=1

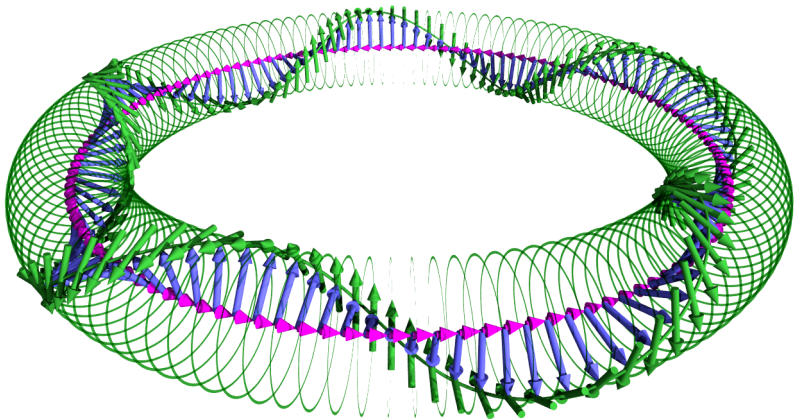
$$m\dot{\gamma} \stackrel{\text{def}}{\text{by}} \begin{pmatrix} c \\ \Phi_0 \\ n\Upsilon_0 \end{pmatrix} \diamond \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos r_a n \omega_0 t & \sin r_a n \omega_0 t \\ 0 & -\sin r_a n \omega_0 t & \cos r_a n \omega_0 t \end{pmatrix}$$

$$m\dot{\gamma} \stackrel{\text{def}}{\text{by}} \begin{pmatrix} c \\ \Phi_0 \\ n\Upsilon_0 \end{pmatrix} \diamond \begin{pmatrix} 1 & 0 & 0 \\ 0 & -\cos r_a n \omega_0 t & -\sin r_a n \omega_0 t \\ 0 & \sin r_a n \omega_0 t & -\cos r_a n \omega_0 t \end{pmatrix}$$



Slide 22: Toroidal EM-eddy — Emtrons in Circular Self-Orbits: Spin=1

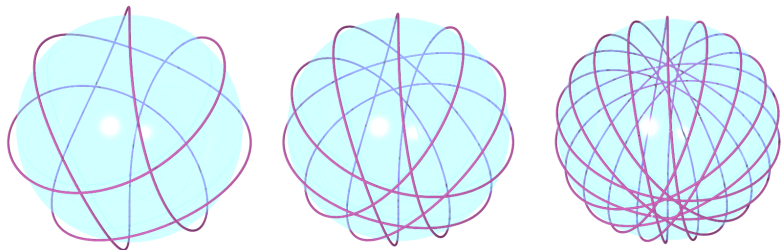
$$m^T \stackrel{\text{def}}{\text{by}} \begin{pmatrix} c \\ \Phi_0 \\ \Upsilon_0 \end{pmatrix} \diamond \begin{pmatrix} \cos \omega_0 t/n & \sin \omega_0 t/n & 0 \\ -\cos \omega_0 t/n \sin \omega_0 t/n & \cos \omega_0 t/n \cos \omega_0 t/n & \sin \omega_0 t/n \\ \sin \omega_0 t/n \sin \omega_0 t/n & -\sin \omega_0 t/n \cos \omega_0 t/n & \cos \omega_0 t/n \end{pmatrix}$$



Slide 23: Particles as EM-Solitons — Emtrons in Spherular Orbits: Spin=1

$$m^{\oplus} \xrightarrow[\text{by}]{\text{def}} \begin{pmatrix} c \\ \Phi_0 \\ Y_0 \end{pmatrix} \diamond \begin{pmatrix} \cos 2\omega_0 t/m & -\sin 2\omega_0 t/m \sin \omega_0 t/mn & \sin 2\omega_0 t/m \cos \omega_0 t/mn \\ 0 & \cos \omega_0 t/mn & \sin \omega_0 t/mn \\ -\sin 2\omega_0 t/m & -\cos 2\omega_0 t/m \sin \omega_0 t/mn & \cos 2\omega_0 t/m \cos \omega_0 t/mn \end{pmatrix}$$

where $n \in \{2, 3, 5, \dots \text{prime}\}$, and m an integer scaling value. The integral $\mathbf{p} = \int \mathbf{u} dt$ determines the path shape which has a length $2mnl_0$ and encloses a sphere of radius $r_s = ml_0/(2\pi)$.



Slide 24: Summary and Conclusion

Please refer to my full paper "A Novel Constructive Electromagnetic Quantum Theory describes the Origin of Mass and Unifies the Forces" available <https://hnp.onl/1882>

- Theorem 1 and 2 provide the mathematical framework to define electromagnetic solitons.
- We are required to separate the electrostatic fields from the electromotive fields. They are two different phenomena that combine into the electromagnetic phenomenon.
- The various solutions of $\mathcal{M}(\mathbf{u}, \phi_0, \mathbf{Y}_0)$ give explanation to electric currents, photons and particles.
- The Origin of Mass is an EM-phenomenon, $E = mc^2$ is derived from energies of these solitons in \mathbb{C}^3
- The paper provides explanation to unify the forces.
- The paper provides explanation to all quantum phenomena
- The paper provides a method for algorithmic nucleus packing of all elements and their isotopes giving the correct atomic mass number.

Unfortunately Theorem 1 and 2 require \mathbb{C}^3 and will not work in the four dimensional space-time constructs. I predict that a paradigm revolution in physics is upon us.

Thank You

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Full paper: <https://hnp.onl/1882>