

E_7 and E_8 ALH* Tesserons as the moduli space of Doubly Periodic Monopoles

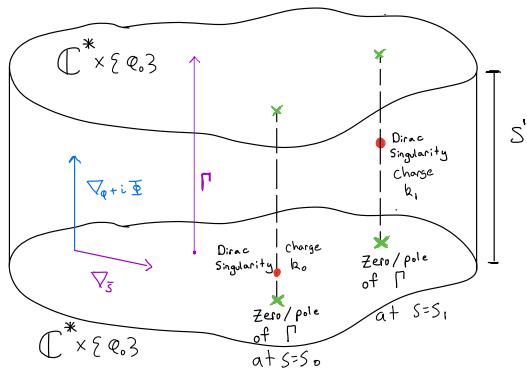
Thomas Harris

University of Arizona

Doubly Periodic Monopoles

- Hermitian Vector Bundle $E \rightarrow \mathbb{R} \times S^1 \times S^1$ with coordinates (x, θ, φ)
- Higgs field Φ and connection D_A satisfying Bogomolny equations
 $*D_A\Phi = -F_A$
- Allow Dirac style singularities
- At most linear growth of Φ as $x \rightarrow \pm\infty$.

Reinterpreting the Bogomolny Equations

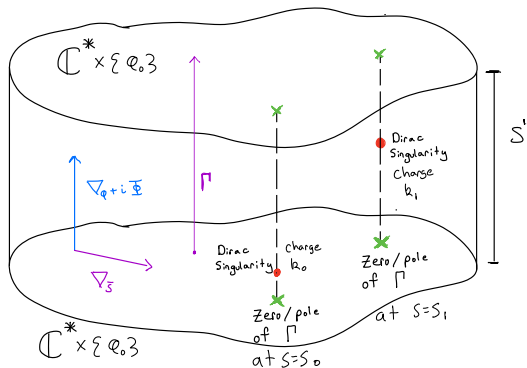


$$\mathbb{R} \times S^1 \times S^1 \simeq \mathbb{C}^* \times S^1 \quad (1)$$

$$(x, \theta, \varphi) \mapsto (s = \exp(x + i\theta), \varphi) \quad (2)$$

$$D_A = \nabla_x dx + \nabla_\theta d\theta + \nabla_\varphi d\varphi \implies \nabla_{\bar{s}} = \frac{1}{2\bar{s}} \left(\nabla_x + i\nabla_\theta \right) \quad (3)$$

Holomorphic view of Bogomolny Equations

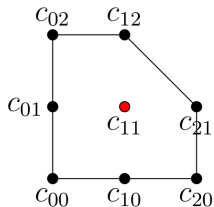


- $\nabla_{\varphi} + i\Phi$ induces parallel transport in φ direction
- Denote resulting holonomy operator $\Gamma(s) : E_{\varphi_0} \rightarrow E_{\varphi_0}$
- $\nabla_{\bar{s}}$ is holomorphic structure on E_{φ_0} .

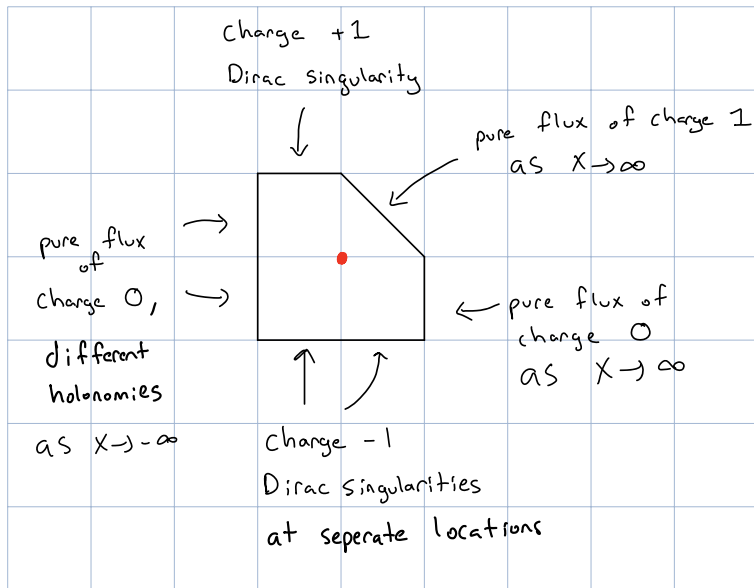
- $\Sigma := \det(\Gamma(s) - t)^{-1}(0)$, curve of eigenvalues of Γ .
- Monopole boundary conditions ensure $\Sigma = p^{-1}(0)$ for polynomial $p(s, t)$.
- From eigenvectors of Γ , get an eigenline bundle $L \rightarrow \Sigma$
- Bogomolny equations imply $\nabla_{\bar{s}}\Gamma = 0$, hence $L \rightarrow \Sigma$ has a holomorphic structure
- Holomorphic line bundle $L \rightarrow \Sigma$ along with parabolic structure at marked points can recover the monopole [**mochizuki20**] [**harris23**]

Newton Polygon

- Consider the polynomial
$$p(s, t) = c_{00} + c_{10}s + c_{20}s^2 + t(c_{10} + c_{11}s + c_{21}s^2) + t^2(c_{02} + sc_{12})$$
- We can represent it by a Newton polygon

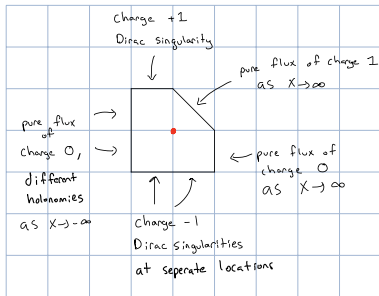
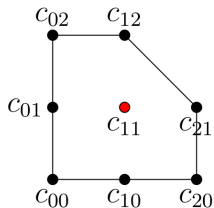


Newton Polygon Interpretation for E_4



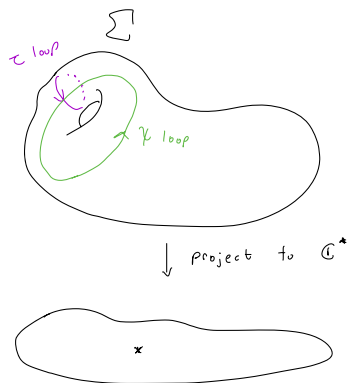
Moduli Space, coefficients

- L^2 deformations of doubly periodic monopoles induce certain deformations in spectral data, meaning L^2 pairs $(\Phi + \phi, D_A + \alpha)$ such that (ϕ, α) satisfy linearized Bogomolny equations
- $p(s, t) = c_{00} + c_{10}s + c_{20}s^2 + t(c_{10} + c_{11}s + c_{21}s^2) + t^2(c_{02} + sc_{12})$
- Boundary terms of p fixed, but c_{11} can change.



Moduli Space, holonomy

- L^2 deformations of doubly periodic monopoles induce certain deformations in spectral data
- Part of the doubly periodic monopole was encoded into a holomorphic structure on a line bundle $L \rightarrow \Sigma$
- Spectral curves holonomies also get changed.



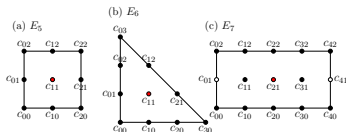
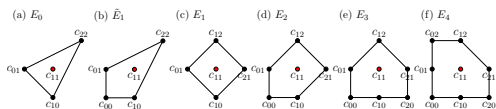
Moduli Space, combined

- L^2 doubly periodic monopole deformations correspond to interior coefficient deformations of Newton polygon and holonomy deformations on spectral curve
- Moduli space of doubly periodic monopoles understood by thinking of holomorphic line bundles over spectral curves

- Tesseracton is a complete, noncompact hyperkähler manifold of real dimension 4 with L^2 Riemann curvature tensor
- ALH^* Tesserons have geodesic ball volume growth $\sim r^{4/3}$.
- Classified by intersection forms, types E_{9-b}
- Model ALH^* metric = $\frac{Hb}{2\pi}(dH^2 + d\rho^2 + d\tau^2) + \frac{2\pi}{Hb}(d\chi - \frac{b\rho}{2\pi}d\tau)^2$
[collins21] [hein21] [sun21]

ALH* Tesseron as moduli space

- With regular boundary conditions, interior points represent 4 times the number of L^2 deformations
- Cherkis, Cross, Ward used this to find E_0, E_1, \dots, E_6 case



E_7 and E_8

- E_7 and E_8 arise from degenerate boundary conditions, [kim17]
- Some interior points of Newton polygon are now fixed

$$\Phi_j = \frac{i}{2|r-r_j|} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

$$\Phi_{-\infty} = \begin{bmatrix} iM_{1,-\infty} & 0 & 0 & 0 \\ 0 & iM_{2,-\infty} & 0 & 0 \\ 0 & 0 & iM_{3,-\infty} & 0 \\ 0 & 0 & 0 & iM_{4,-\infty} \end{bmatrix}$$

$$\Phi_{\infty} = \begin{bmatrix} iM_{1,\infty} & 0 & 0 & 0 \\ 0 & iM_{2,\infty} & 0 & 0 \\ 0 & 0 & iM_{3,\infty} & 0 \\ 0 & 0 & 0 & iM_{4,\infty} \end{bmatrix}$$

$$A_{k,-\infty} = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & ia_{k,-\infty} & 0 & 0 \\ 0 & 0 & ia_{k,-\infty} & 0 \\ 0 & 0 & 0 & ia_{k,-\infty} \end{bmatrix}$$

$$A_{k,\infty} = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & ia_{k,\infty} & 0 \\ 0 & 0 & ia_{k,\infty} & 0 \\ 0 & 0 & 0 & ia_{k,\infty} \end{bmatrix}$$

$$\Phi_j = \frac{-i}{2|r-r_j|} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

$$\Phi_j = \frac{i}{2|r-r_j|} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

$$\Phi_{-\infty} = iM_{-\infty} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$A_{k,-\infty} = ia_{k,-\infty} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$A_{k,\infty} = ia_{k,\infty} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

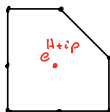
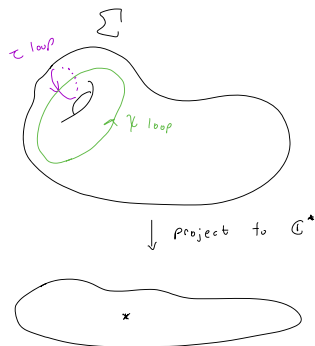
$$\Phi_{\infty} = iM_{\infty} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$A_{k,\infty} = ia_{k,\infty} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\Phi_j = \frac{-i}{2|r-r_j|} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

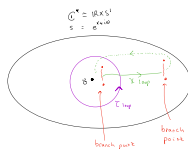
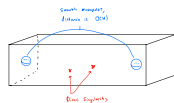
ALH^* , spectral side

- Interior coefficient is $e^{H+i\rho}$
- Holomorphic line bundle parameterized by holonomies around 1-cycles, coordinates of (τ, χ)
- 4 coordinates of (H, ρ, τ, χ)
- Model ALH^* metric = $\frac{Hb}{2\pi}(dH^2 + d\rho^2 + d\tau^2) + \frac{2\pi}{Hb}(d\chi - \frac{b\rho}{2\pi}d\tau)^2$



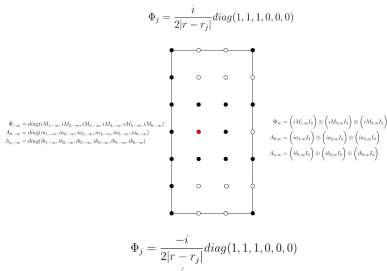
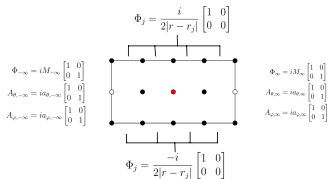
ALH^* , monopole side

- L^2 metric on monopole side should approximate model ALH^* metric
- Φ, D_A split into abelian solutions between monopoles
- Branch points of Σ are like smooth monopoles
- Distance between monopoles is $O(H)$.
- Coordinates (H, ρ, τ, χ)
- H represents size of Higgs field between monopoles
- ρ represents φ holonomy between monopoles
- τ represents θ holonomy between monopoles
- χ represents holonomy by transporting from one monopole to the other, changing sheets, and transporting back.

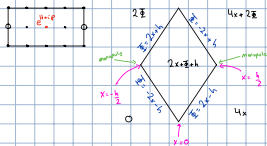


E_7 and E_8

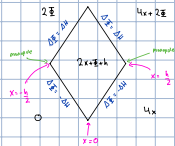
- E_7 and E_8 arise from degenerate boundary conditions
- Some interior points of Newton polygon are now fixed



Dual Polygons



$$\{0\}^1 \sim 2\pi(2\pi)^1$$



- For Long E_7 , deformations between monopoles is essentially constant

- $\Delta\Phi \sim \Delta H \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

- $\Delta A_\varphi \sim \Delta\rho \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

- $\Delta A_\theta \sim \Delta\tau \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

- $\Delta A_x \sim \frac{\Delta\chi}{H} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

- Squaring an integrating over a region of length H yields model ALH^* metric

- Model ALH^* metric = $\frac{Hb}{2\pi}(dH^2 + d\rho^2 + d\tau^2) + \frac{2\pi}{Hb}(d\chi - \frac{b\rho}{2\pi}d\tau)^2$

Thank you for your time.



Sung-Soo Kim and Futoshi Yagi

5d En Seiberg-Witten curve via toric-like diagram

arXiv:1411.7903v4 [hep-th], 11 Jul 2017.



Hans-Joachim Hein, Song Sun, Jeff Viaclovsky, and Ruobing Zhang

Gravitational instantons and del pezzo surfaces, 2021



Song Sun and Ruobing Zhang

Collapsing geometry of hyperkaehler 4-manifolds and applications, 2021



Tristan Collins, Adam Jacob, and Yu-Shen Lin

The Torelli theorem for alh^* gravitational instantons, 2021



Takuro Mochizuki

Kobayashi-Hitchin correspondence for analytically stable bundles

Trans. Amer. Math. Soc., 373(1):551–596, 2020



Thomas Harris

All ALH^* Tesserons from Doubly Periodic Monopoles

In progress

Meaning of b



b.8 $\tau_{11} \circ \tau_{12}$



b.7 $\tau_{11} \circ \tau_{13}$



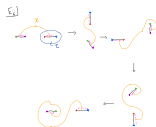
b.6 $\tau_{11} \circ \tau_{14}$



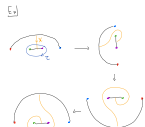
b.5 $\tau_{11} \circ \tau_{15}$



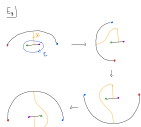
b.4 $\tau_{11} \circ \tau_{16}$



b.3 $\tau_{11} \circ \tau_{18}$



b.2 $\tau_{11} \circ \tau_{19}$



b.1 $\tau_{11} \circ \tau_{20}$