

# Tesserons from Monopoles

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Topological Solitons  
Nagoya Online Workshop  
September 2023

# Tesserons

We require all Riemannian four-manifolds below to be [complete](#) with  $L^2$  Riemann curvature:

$$\int_{M^4} |Rm|^2 dVol < \infty$$

- **Gravitational Instanton** (defined by Hawking '77) is any solution of the vacuum Einstein equations:

$$Rm = \Lambda g.$$

- **Self-dual Gravitational Instanton** has self-dual Riemann curvature tensor

$$Rm = * Rm.$$

Locally, these are hyperkähler!

- Main interest: **Hyperkähler gravitational instanton**:

“Complete hyperkähler manifold of real dimension 4 with  $L^2$  Riemann curvature tensor”

Based on the underlying quaternionic structure:

Q: What should we call it?

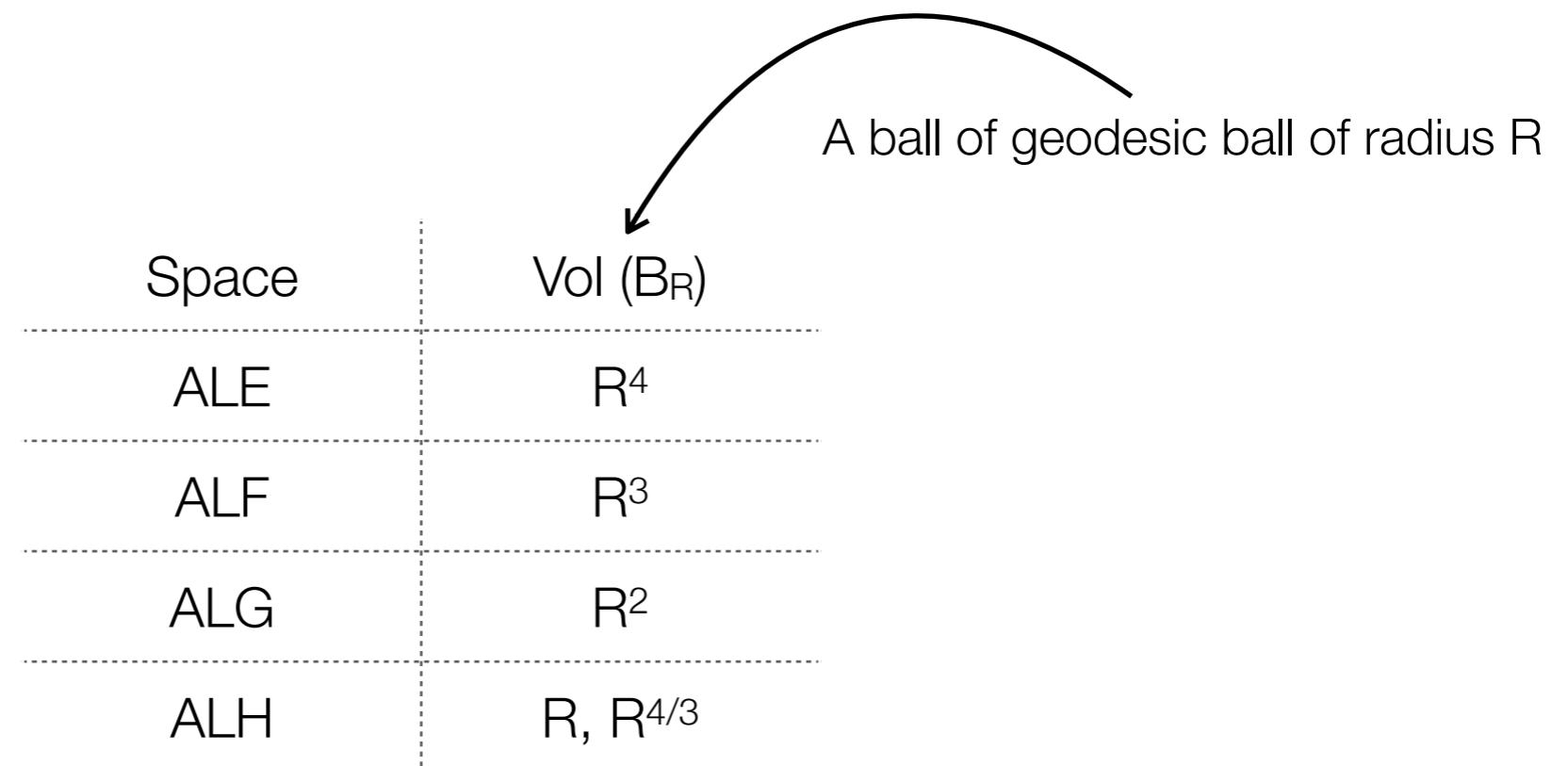
Quaternion? Quatron? Twiston? Hamilton?

- A **Tesserón** is a “Complete noncompact hyperkähler manifold of real dimension 4 with  $L^2$  Riemann curvature tensor.”

Classification of Tesserons was recently completed:

- ALE: Kronheimer '89
- ALF: Minerbe '07, '08
- ALG & ALH: G. Chen and X.-X. Chen '15; G.Chen and Viaclovsky '21
- ALG\*: G. Chen and Viaclovsky '21
- ALH\*: Hein, Sun, Viaclovsky, Zhang '21; Collins, Jacob, Lin '21; Lee, Lin '22

Tesserons are distinguished by their asymptotic **Volume Growth**:



## Asymptotic Model

As it happens, all tesseron metrics locally have asymptotic triholomorphic isometry.

According to Gibbons-Hawking, a metric with such an isometry locally has the form

$$g = V \vec{x}^2 + \frac{(d\tau + \omega)^2}{V}, \text{ where } {}^*_3 dV = d\omega,$$

Tesserons' model ends have (locally):

- ALE  $V = \frac{N}{2|\vec{x}|}$

- ALF  $V = \ell + \frac{N}{2|\vec{x}|}$

- ALG  $V = C + \frac{N}{2} \ln(x_1^2 + x_2^2)$

- ALH  $V = C + Nx_1$

Current literature distinguishes:

ALG\* and ALH\* are spaces with  $N \neq 0$ , and

ALG and ALH are with  $N = 0$  (locally constant fiber).

- ALE  $V = \frac{N}{2|\vec{x}|}$

### Prototypical example:

$\mathbb{R}^4$  metric in ‘radial coordinates’

$$g = \frac{1}{2x} d\vec{x}^2 + 2x(d\theta + \omega)^2$$

- ALF  $V = \ell + \frac{N}{2|\vec{x}|}$

### The Taub-NUT:

$$g = (\ell + \frac{1}{2x}) d\vec{x}^2 + \frac{(d\theta + \omega)^2}{\ell + \frac{1}{2x}}$$

- ALG  $V = C + \frac{N}{2} \ln(x_1^2 + x_2^2)$

(ALG\* if  $N \neq 0$ )

### Elliptic Fibrations:

$$g = \tau_2 dz d\bar{z} + \frac{|d\theta_a + \tau d\theta_b|^2}{\tau_2},$$

$$\tau = \tau_1 + i\tau_2 = C + N \frac{i}{2\pi} \ln z$$

- ALH  $V = C + Nx_1$

(ALH\* if  $N \neq 0$ )

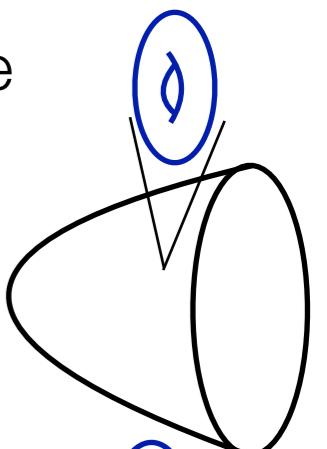
$$\tau = \tau_1 + i\tau_2 = C + iNz$$

### Asymptotic metric:

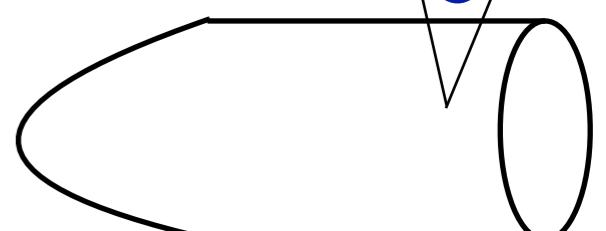
Circle fibration  
(growing circle):  
Quotient:  $\mathbb{R}^4/\Gamma$ ,  
 $\Gamma \subset SU(2)$ .

Circle fibration  
(with stabilizing circle):

Over cone base



Over cylinder base



# Classification

ALE:  $\mathbb{R}^4$ ,  $A_{k \geq 1}$ ,  $D_{k \geq 1}$ , and  $E_6, E_7, E_8$      $A_0 = \mathbb{R}^4$ ,  $A_1 = \text{Eguchi-Hanson}$

ALF:  $\mathbb{R}^3 \times S^1$ ,  $A_{k \geq 0}$  and  $D_{k \geq 0}$

- $A_0 = \text{Taub-NUT}$
- $A_k = (k+1)\text{-centered multi-Taub-NUT}$
- $D_0 = \text{Atiyah-Hitchin,}$
- $D_1 = \text{deformation of double cover of } D_0$
- $\mathbf{D}_2 = \text{deformation of } (\mathbb{R}^3 \times S^1)/\mathbb{Z}_2$

**ALG:**

& ALG\*  $E_6$ ,  $E_7$ ,  $E_8$

$$\mathbb{R}^2 \times T^2, \ D_0, D_1, D_2, D_3, \ \mathbf{D}_4, \ \mathbf{\Psi}_2, \mathbf{\Psi}_1, \mathbf{\Psi}_0$$

**ALH:**  $E_0, E_1,$

$$\& \text{ALH}^* \quad \mathbb{R} \times T^3, \quad E_2, \ E_3, \ E_4, \ E_5, \ E_6, \ E_7, \ E_8, \ \frac{1}{2}\text{K3} \\ \tilde{E}_1,$$

# Naive Parameter Count

$m_n$  denotes  $n$  real parameters specifying the form of infinity and  $m$  “interior” parameters.

**ALE:**  $\mathbb{R}^4, A_{k \geq 1}, D_{k \geq 1}$ , and  $E_{k=6,7,8}$ .  
 $(0)_0 \quad (3k)_0 \quad (3k)_0 \quad (3k)_0$

Note: additional isometries reduce this Naive count,  
e.g.  $A_1$  ALE –  $(1)_0$   
and  $A_1$  ALF –  $(1)_1$

**ALF:**  $\mathbb{R}^3 \times S^1, A_{k \geq 0}$  and  $D_{k \geq 0}$   
 $(0)_1 \quad (3k)_1 \quad (3k)_1$

<b>ALG:</b>				$IV^*$	$III^*$	$II^*$
& ALG*		$I_{4-k}^*$	$I_0^*$	$E_{k=6,7,8}$		
				$(3k)_1$	i.e.	
	$\mathbb{R}^2 \times T^2,$	$D_{k=0,1,2,3}$	$D_4,$	$\mathbf{D}_{4,}$	$IV$	$III$
	$(0)_3$	$(3k)_3$	$(3k)_3$	$\mathbf{W}_{k=0,1,2}$		$II$
				$(2+3k)_1$ (?)		

<b>ALH:</b>	$E_0, E_1,$			
& ALH*	$\mathbb{R} \times T^3,$	$E_2, E_3, E_4, E_5, E_6, E_7, E_8,$	$\frac{1}{2}\mathbf{K3}$	
	$(0)_7$	$\tilde{E}_1,$	$(3k)_3$	$(3x8)_7$

# Discrete Painlevé Eqs.

Noumi et al 1998  
 Hidetaka Sakai 2001  
 (Table 2)

The list of ALG<sup>(\*)</sup> and ALH<sup>(\*)</sup> tesserons coincides with Discrete Painlevé Equations (and associated rational surfaces)!

Note: As a result all ALG and ALG\* tesserons are moduli spaces of Hitchin systems!

**ALG & ALG\***    **Rational** Painlevé:

$$\begin{array}{ccccccccc}
 E_8^{(1)} & E_7^{(1)} & E_6^{(1)} & D_4^{(1)} & A_3^{(1)} & (A_1 + A_1)^{(1)} & A_1'^{(1)} & A_0^{(1)} \\
 E_8 \rightarrow E_7 \rightarrow E_6 \rightarrow D_4 \rightarrow D_3 \rightarrow D_2 \rightarrow D_1 \rightarrow D_0 \\
 & & & \searrow & & & & \\
 & & \mathbb{W}_0 \longleftarrow \mathbb{W}_1 \longleftarrow \mathbb{W}_2 \\
 & & A_0^{(1)} & A_1^{(1)} & A_2^{(1)} & & &
 \end{array}$$

**ALH**

**Trigonometric** Painlevé:

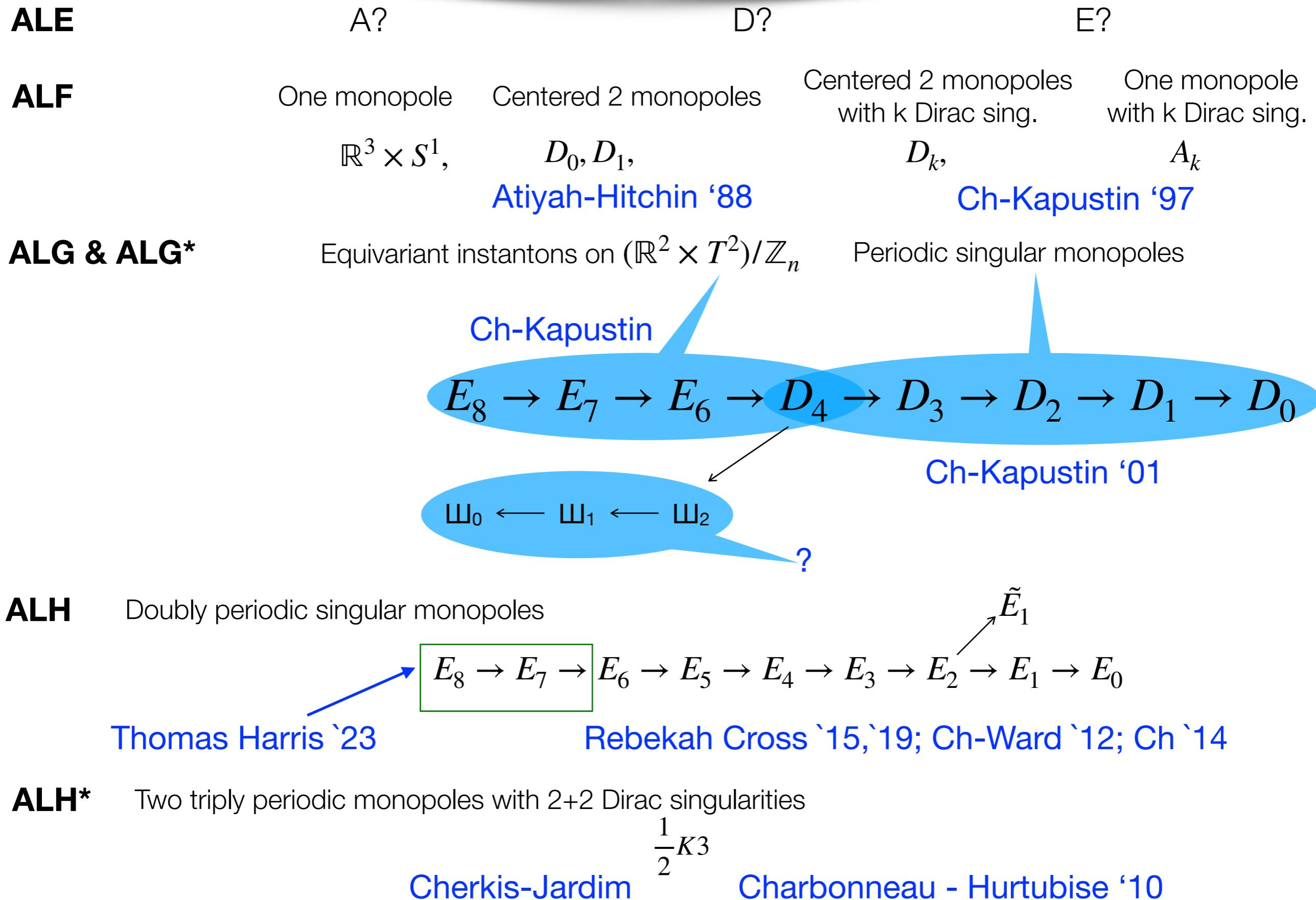
$$\begin{array}{ccccccccccc}
 & & D_5^{(1)} & & (A_2 + A_1)^{(1)} & & \nearrow \tilde{E}_1^{(1)} \\
 E_8^{(1)} \rightarrow E_7^{(1)} \rightarrow E_6^{(1)} \rightarrow E_5^{(1)} \rightarrow E_4^{(1)} \rightarrow E_3^{(1)} \rightarrow E_2^{(1)} \rightarrow E_1^{(1)} \rightarrow E_0^{(1)} \\
 & & A_4^{(1)} & & (A_1 + A_1')^{(1)} & & A_1'^{(1)} & & A_0^{(1)}
 \end{array}$$

**ALH\***

**Elliptic** Painlevé:

$$E_8^{(1)} = \frac{1}{2}K3$$

# Tesserons as Moduli Spaces of Monopoles



# Monopole

A **monopole** on a three-dimensional Riemannian manifold  $(X, g)$  is

- 1) a Hermitian vector bundle  $E \rightarrow X$  and
- 2) a pair  $(A, \Phi)$  of a connection  $A$  on  $E$  and  
an skew-hermitian endomorphism  $\Phi$  of  $E$ ,  
satisfying the Bogomolny Equation

$$F_A = {}^* d_A \Phi,$$

and appropriate asymptotic conditions.

## **Monopole on**

$$X = \mathbb{R}^3$$

Monopole

## **Moduli Space type:**

ALF

$$X = \mathbb{R}^2 \times S_R^1$$

Periodic Monopole

ALG

$$X = \mathbb{R}^1 \times T_{A,\tau}^2$$

Doubly Periodic Monopole =  
Monopole Wall = Monowall

ALH

# Singular Monopoles

Simple Dirac singularities at marked points  $p_1^-, \dots, p_{v_-}^-$  and  $p_1^+, \dots, p_{v_+}^+$ :  $\Phi = i \begin{pmatrix} \pm \frac{1}{2|\vec{x} - \vec{p}_\sigma^\pm|} & 0 \\ 0 & 0_{n-1, n-1} \end{pmatrix} + O(1)$

Note: More generally the charge is any cocharacter of the gauge group. (Talk by Thomas Harris on Friday.)

## Boundary conditions:

$\mathbb{R}^3$

$$\Phi^g(\vec{x}) = \frac{i}{2\pi} \begin{pmatrix} \lambda - \frac{\varrho}{2|\vec{x}|} & 0 \\ 0 & -\lambda + \frac{\varrho}{2|\vec{x}|} \end{pmatrix} + O(r^{-2}) \quad \begin{array}{l} \text{Finite energy} \Rightarrow \\ \text{center of mass is a "modulus"} \end{array}$$

$\mathbb{R}^2 \times S^1$

Coordinates

$z = x + iy, \varphi$

$$\Phi^g = \frac{i}{2\pi} \text{diag}(v_j + q_j \log |z| + \text{Re} \frac{\mu_j}{z}) + O(1/|z|^2) \quad \begin{array}{l} \text{Infinite energy} \Rightarrow \\ \text{center of mass is fixed} \end{array}$$

$$A^g = \frac{1}{2\pi} \text{diag}((q_j \arg z + [b_j + \text{Im} \frac{\mu_j}{z}]) d\theta + \alpha_j d\arg z) + O(|z|^{-2})$$

$\mathbb{R} \times T^2$

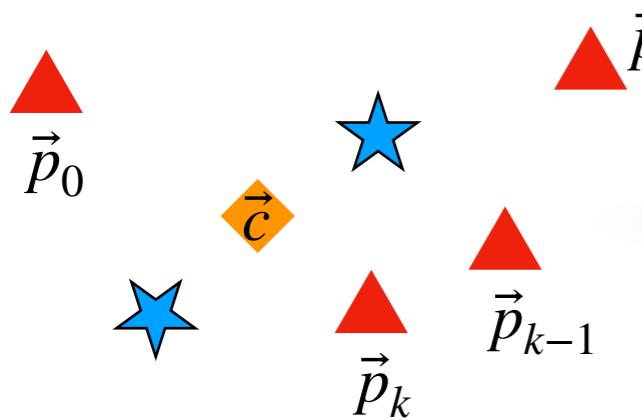
Coordinates

$x, \theta, \varphi$

$$\Phi^g = \frac{i}{2\pi} \text{diag}(Q_j^\pm x + M_j^\pm) + O(1/x) \quad \begin{array}{l} \text{Infinite energy} \Rightarrow \\ \text{center of mass is fixed} \end{array}$$

$$A^g = -\frac{i}{2\pi} \text{diag}(Q_j^\pm \theta d\varphi + \chi_{j,\theta}^\pm d\theta + \chi_{j,\varphi}^\pm d\varphi) + O(1/x)$$

# Monopoles in $\mathbb{R}^3$



**Moduli Space:**

$$\mathbb{R}^3 \times S^1_{\frac{1}{\sqrt{\lambda}}}$$

- A single monopole moduli: position in  $\mathbb{R}^3$  and phase in  $S^1$

$$\Phi^g(\vec{x}) = \frac{i}{2\pi} \begin{pmatrix} \lambda - \frac{1}{2|\vec{x}|} & 0 \\ 0 & -\lambda + \frac{1}{2|\vec{x}|} \end{pmatrix} + O(r^{-2})$$

- A single monopole with  $k+1$  simple Dirac singularities

$$\vec{p}_0, \dots, \vec{p}_k \in \mathbb{R}^3$$

$A_k$  ALF = multi-Taub-NUT  
with NUTs at  $\vec{p}_0, \dots, \vec{p}_k$

- Two monopoles in  $\mathbb{R}^3$ : two positions and two phases  
 $\Rightarrow 8$  dim moduli space with triholomorphic isometry

Two centered monopoles in  $\mathbb{R}^3$  with center at  $\vec{c} \in \mathbb{R}^3$

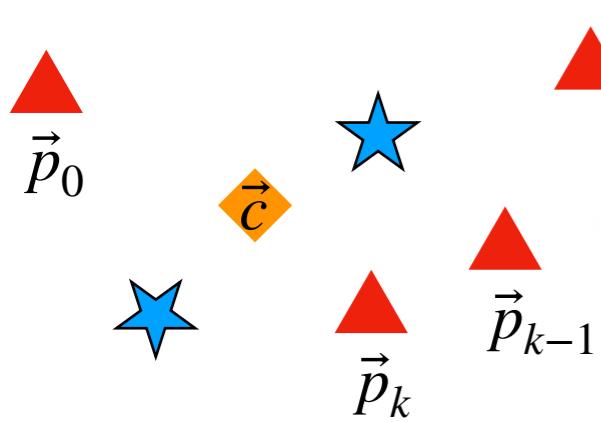
$D_0$  ALF = Atiyah-Hitchin

Two centered monopoles with  $k$  simple Dirac singularities  $\vec{p}_0, \dots, \vec{p}_{k-1} \in \mathbb{R}^3$

$D_k$  ALF

This picture leads to direct relations between these spaces!

## Relations between ALF Spaces

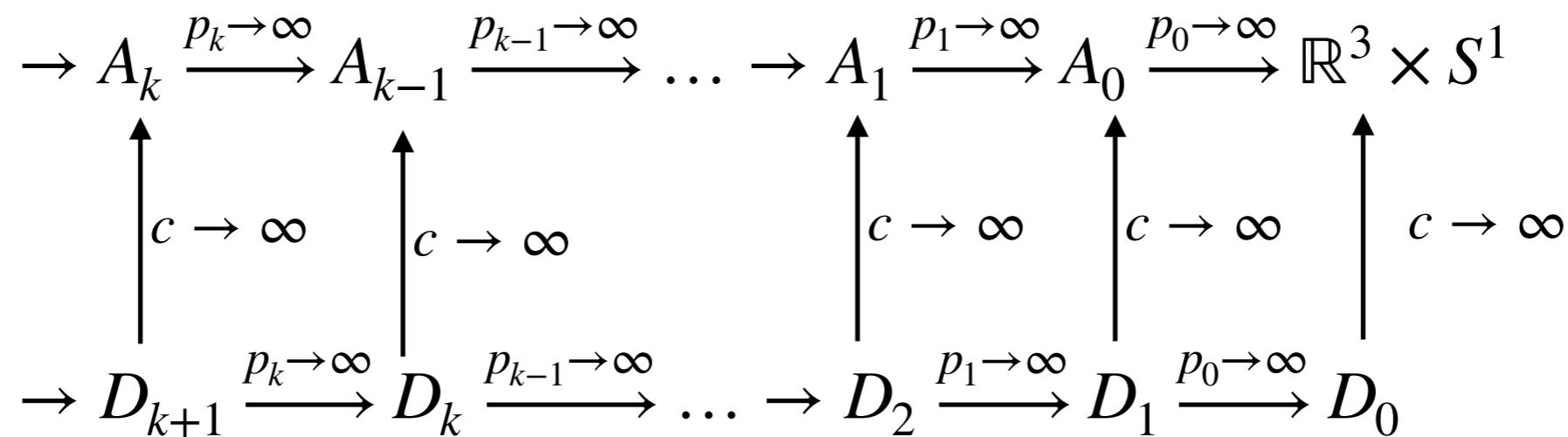


A single with  $k + 1$  simple Dirac singularities  $\vec{p}_0, \dots, \vec{p}_k \in \mathbb{R}^3$

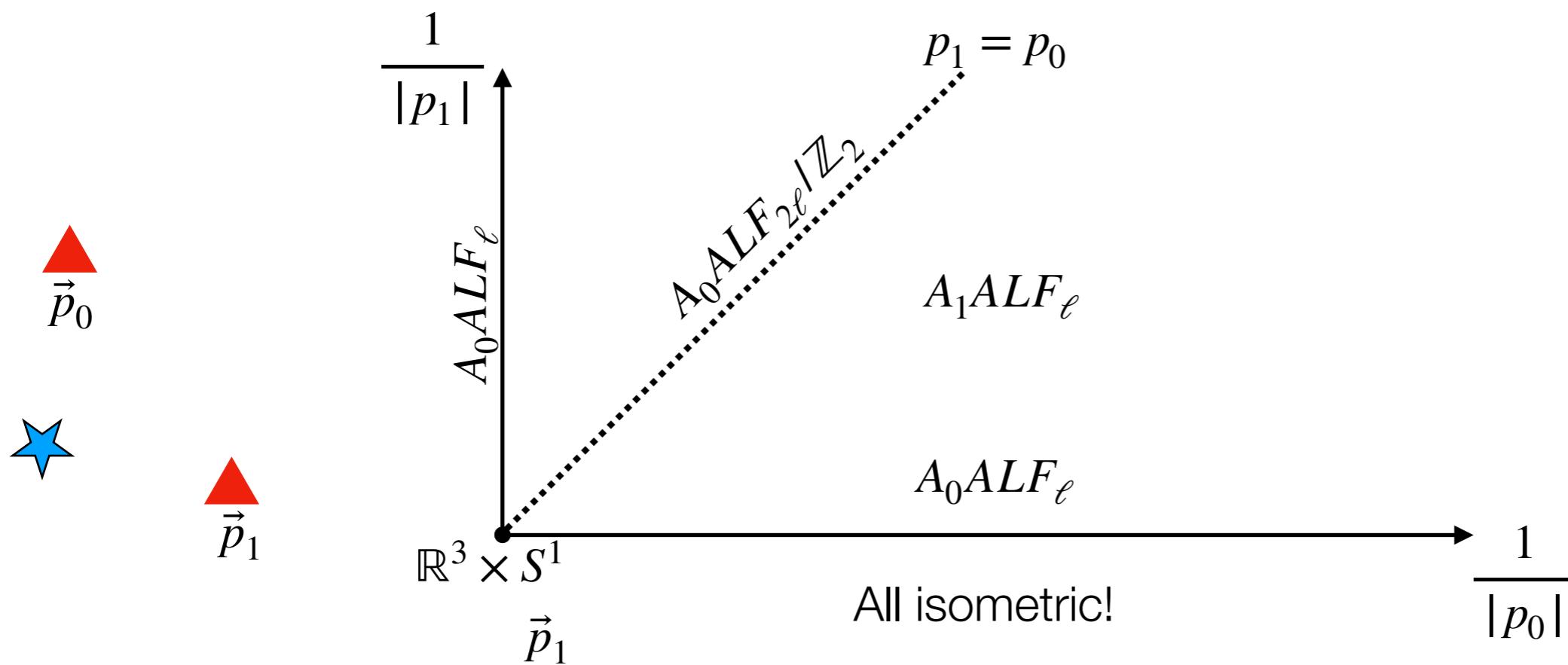
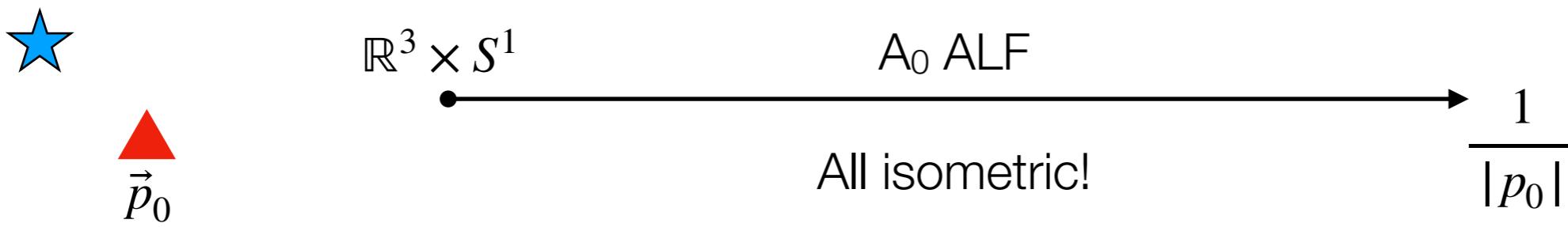
$A_k$  ALF

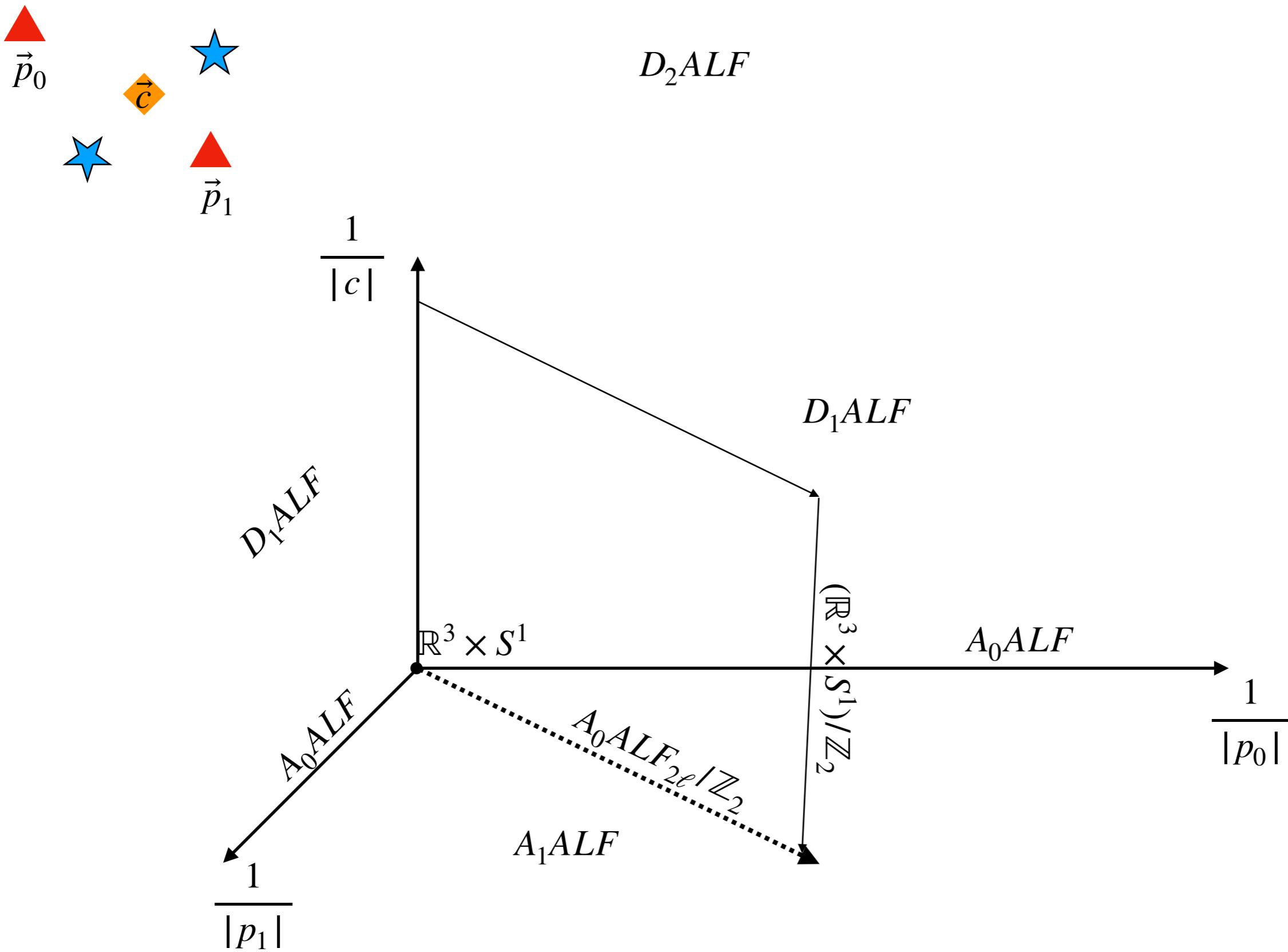
Two centered monopoles with  $k$  simple Dirac singularities  $\vec{p}_0, \dots, \vec{p}_{k-1} \in \mathbb{R}^3$

$D_k$  ALF



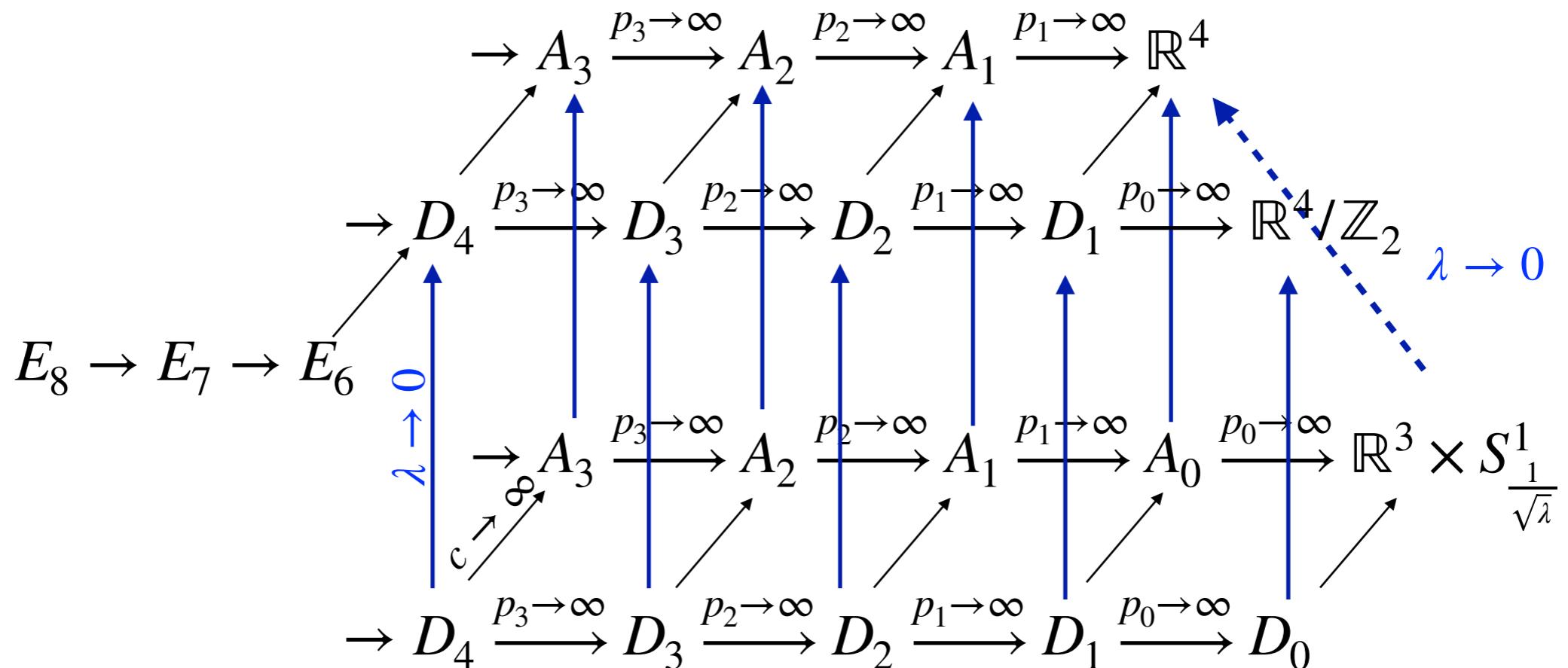
Very schematically:



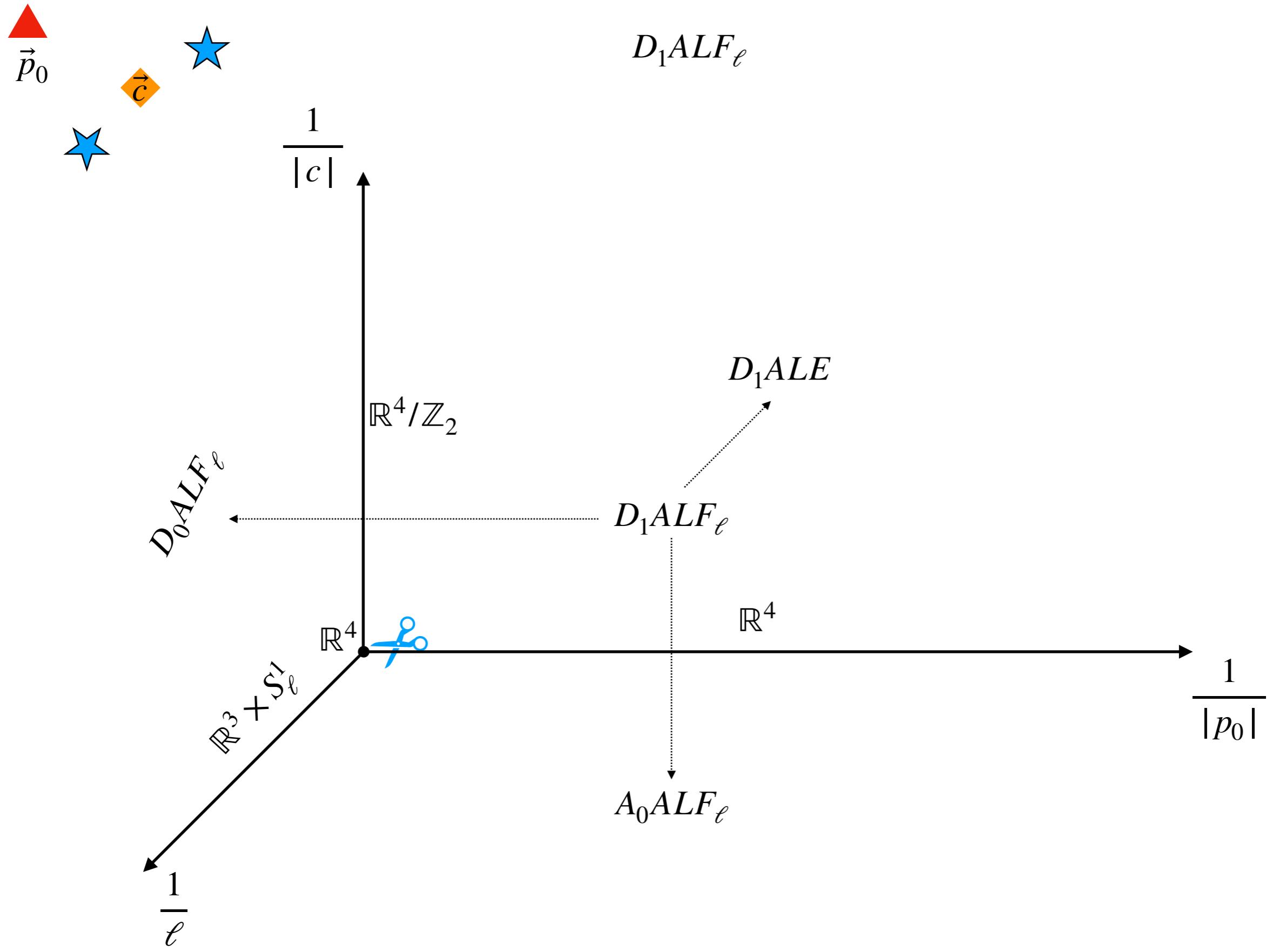


# Relation to ALE

**ALE**



Moral: A- and D-ALE spaces are moduli spaces of monopoles of holomorphic charge 1 or centered monopoles of holomorphic charge 2.



$\mathbb{R}^2 \times S_R^1$

## Periodic Monopoles

$$\Phi^g = \frac{i}{2\pi} \text{diag}(v_j + q_j \log|z| + \text{Re} \frac{\mu_j}{z}) + O(1/|z|^2)$$

$$A^g = \frac{1}{2\pi} \text{diag}((q_j \arg z + [b_j + \text{Im} \frac{\mu_j}{z}]) d\theta + \alpha_j d\arg z) + O(|z|^{-2})$$

Of the six parameters  $\mu_j \in \mathbb{C}, v_j, b_j, q_j, \alpha_j \in \mathbb{R}$  three are true parameters of the moduli space and three determine the center of mass of the periodic monopole.

Two periodic monopoles can have up to 4 singularities.

Moduli space of two periodic monopoles with  $k$  singularities =  $D_k$  ALG space.

This realization makes it clear that as  $R \rightarrow \infty$ , a periodic monopole tends to a monopole in  $\mathbb{R}^3$  and  $D_k$  ALG space approaches  $D_k$  ALF space.

Question: What is the place of E- and SU-ALG spaces?

I do not have a realization of these spaces as moduli spaces of monopoles, only as limits of such spaces;

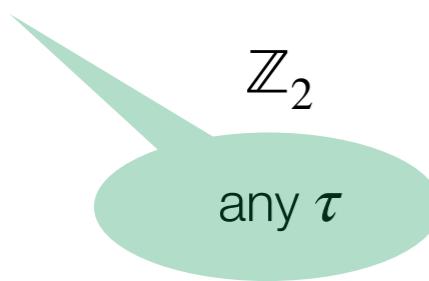
instead, let us realize them as moduli spaces of instantons on  $\mathbb{R}^2 \times T^2$  equivariant under the action of the cyclic group  $\mathbb{Z}_n, n = 3, 4, 6$ .

# Equivariant DP Instanton

Consider YM Instantons on  $\mathbb{R}^2 \times T^2$ , where  $T^2 = \mathbb{C}/(\mathbb{Z} + \tau\mathbb{Z})$

This admits cyclic action  $(z, v) \mapsto (\omega z, \omega^{\pm 1}v)$ , with  $\omega = e^{\frac{2\pi}{n}i} \in \mathbb{Z}_n$

3 real parameters:  
size and shape of  $T^2$



$$A(\omega z, \omega v) = V^{-1} A(z, v) V$$

$\mathbb{Z}_3$

$$\tau = \rho := e^{\frac{2\pi}{6}i}$$

$\mathbb{Z}_4$

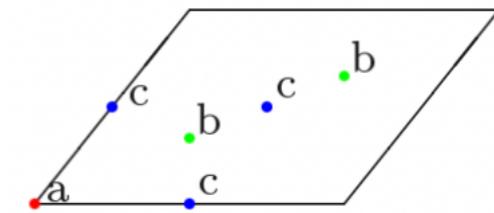
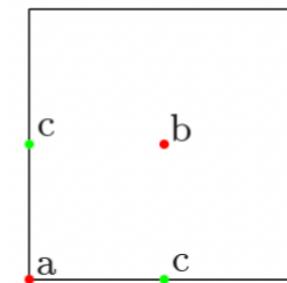
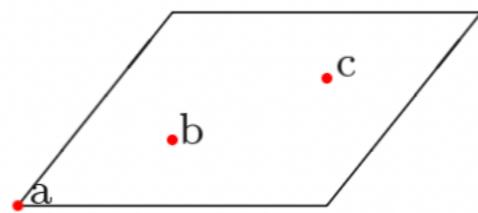
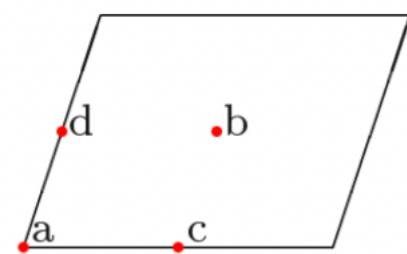
$$\tau = i$$

Only 1 real parameter:  
fixed shape of  $T^2$

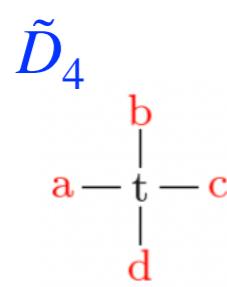
$\mathbb{Z}_6$

$$\tau = \rho := e^{\frac{2\pi}{6}i}$$

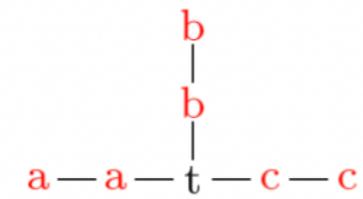
Fixed points  $\rightarrow$  Asymptotic of spectral curve:



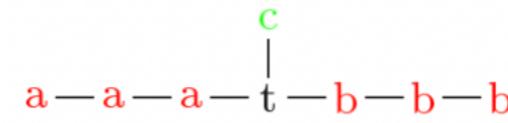
Intersection diagram of the resolution:



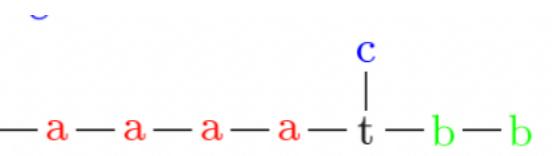
$\tilde{E}_6$



$\tilde{E}_7$



$\tilde{E}_8$



Instanton gauge groups:

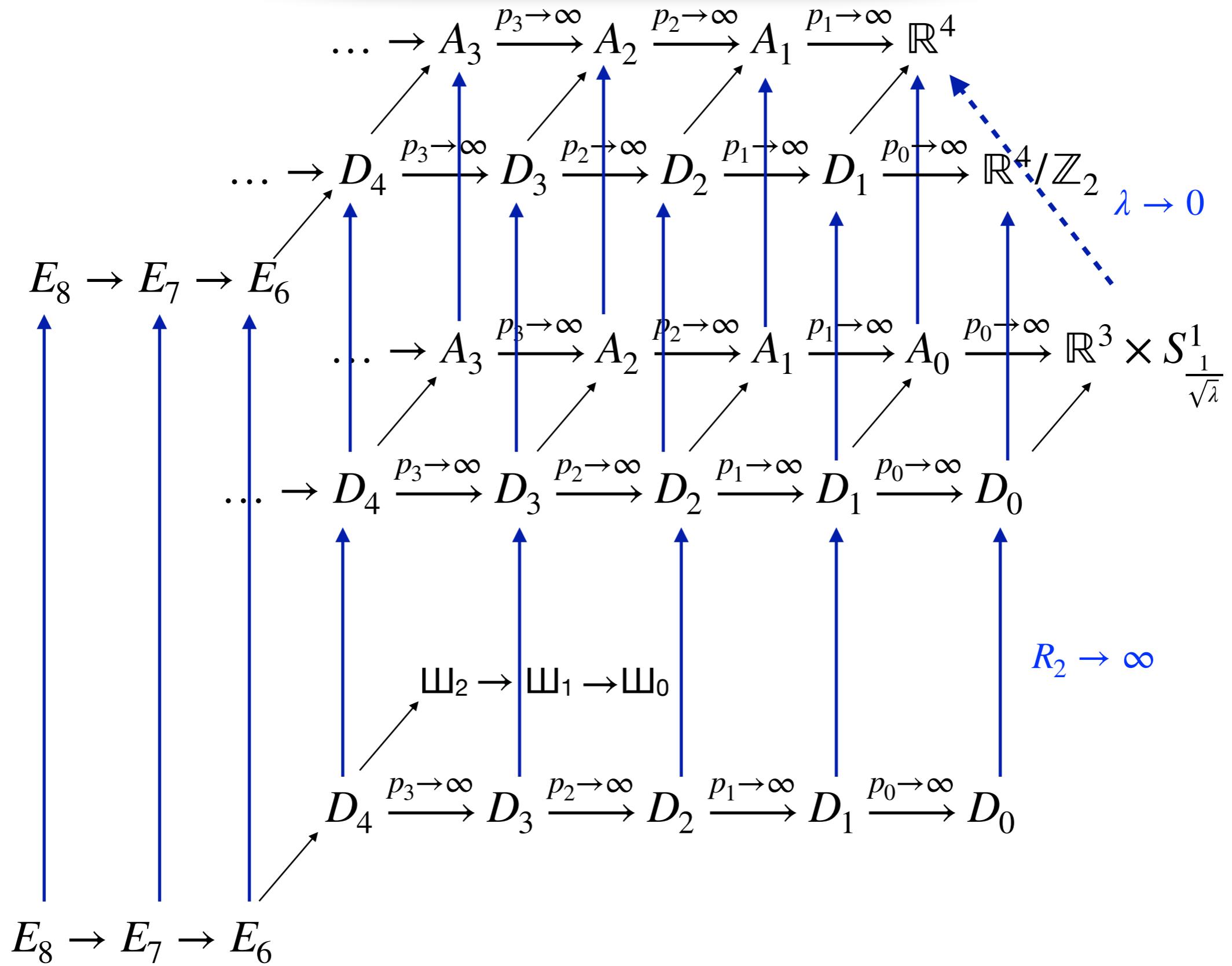
$U(4)$

$U(3)$

$U(4)$

$U(6)$

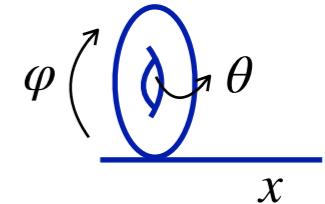
# Including ALG



# Doubly Periodic Monopoles

Ki-Myeong Lee '98  
 Fujimori, Nitta, Ohta, Sakai  
 and Yamazaki '08

Monowall = Monopole Wall = Monopole on  $\mathbb{R} \times S^1_\theta \times S^1_\varphi$



- Abelian monowall
- $d\alpha = {}^*d\phi \Rightarrow d{}^*d\phi = 0$

thus  $\phi$  is harmonic on  $\mathbb{R} \times S^1_\theta \times S^1_\varphi$ , (excluding exponential growth)  $\phi = Qx + M$  with  $Q \in \mathbb{Z}$ ,  $M \in \mathbb{R}$ .

- In general
- $F_A = {}^*d_A \Phi$

a) with boundary conditions  $\Phi = \frac{i}{2\pi} \text{diag}(Q_j^\pm x + M_j^\pm) + O(1/x)$

$$A = -\frac{i}{2\pi} \text{diag}(Q_j^\pm \theta d\varphi + \chi_{j,\theta}^\pm d\theta + \chi_{j,\varphi}^\pm d\varphi) + O(1/x)$$

and

b) simple Dirac singularities at marked points  $p_1^-, \dots, p_{v_-}^-$  and  $p_1^+, \dots, p_{v_+}^+$ .

$$\Phi = i \begin{pmatrix} \pm \frac{1}{2|\vec{x} - \vec{p}_\sigma^\pm|} & 0 \\ 0 & 0_{n-1,n-1} \end{pmatrix} + O(1)$$

The charges  $Q_j^\pm$  are rational, with the denominator equal to the multiplicity of  $(Q_j^\pm, M_j^\pm)$ .

Ch-Ward '12  
Ch '14  
Ch-Cross '19

Monopole charges + singularities  $\rightarrow$  Newton polygon N

Number of moduli =  $4 \times$  Internal integer points of N

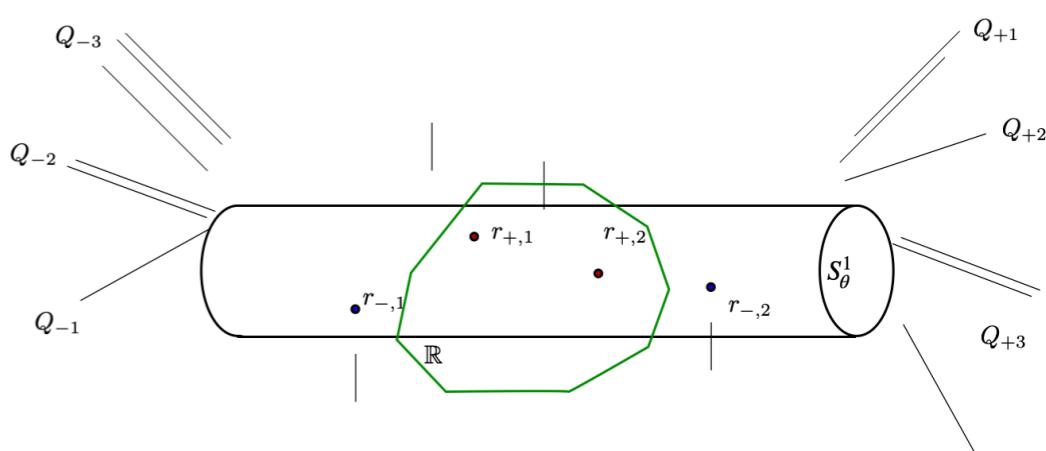
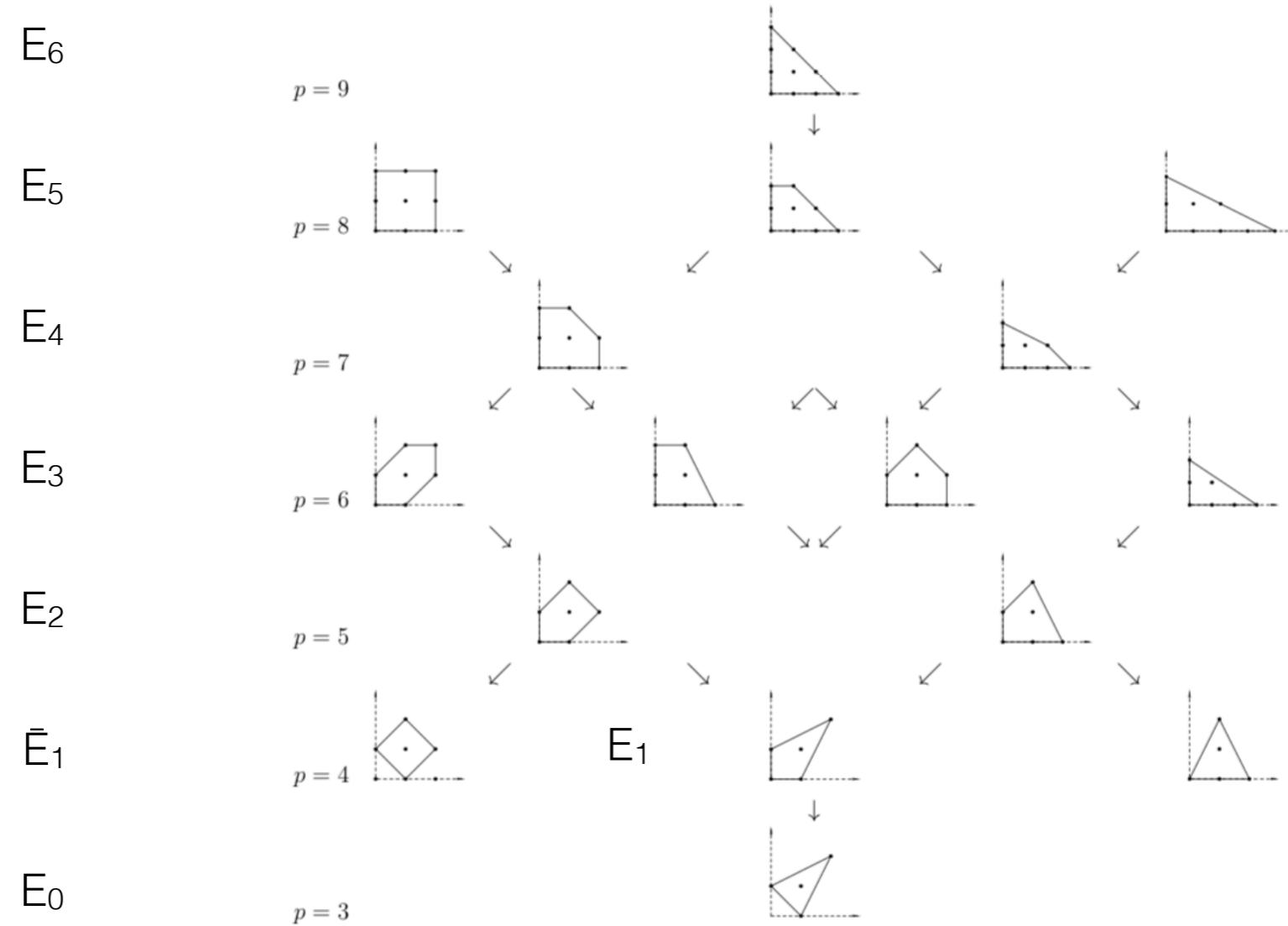
Number of parameters =  $3 \times [(\text{Perimeter integer points of } N) - 3]$

**SL(2, Z)** moduli space isometry generated by

S = Nahm transform and

T = Adding constant magnetic field  $(A, \Phi) \mapsto (A - \theta d\varphi, \Phi + x_1)$ .

All integer Newton polygons with a single internal point up to  $SL(2, \mathbb{Z})$ :

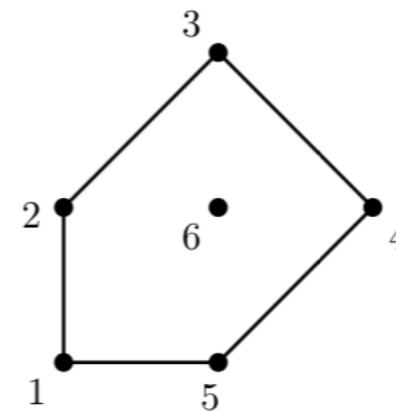


# Secondary Polyhedron

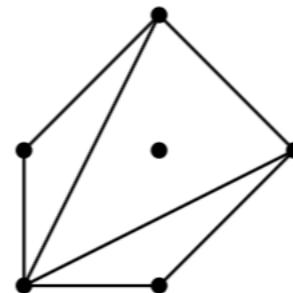
Gelfand, Kapranov,  
Zelevinsky circa '90

Organizes the phase space into sectors.

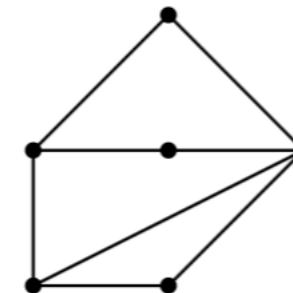
Given a Newton polygon, label its vertices:



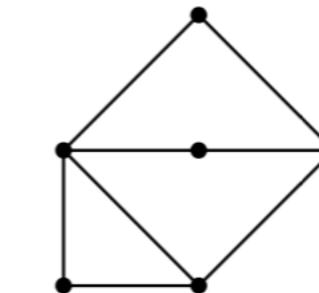
Each regular subdivision gives a  $|N|$ -dimensional vector:



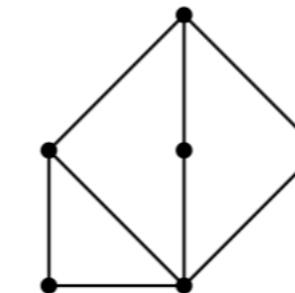
(a)  $(5,1,4,4,1,0)$



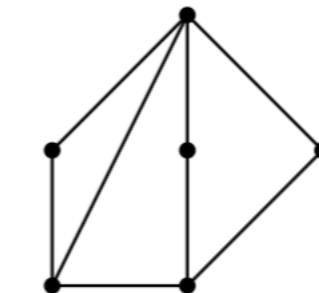
(b)  $(3,4,2,5,1,0)$



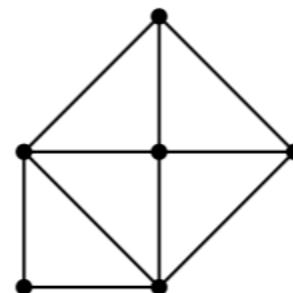
(c)  $(1,5,2,4,3,0)$



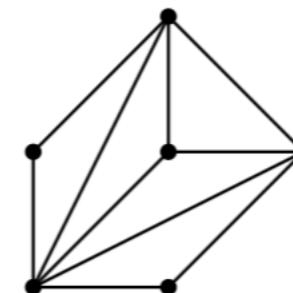
(d)  $(1,3,4,2,5,0)$



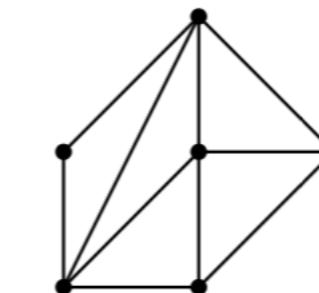
(e)  $(3,1,5,2,4,0)$



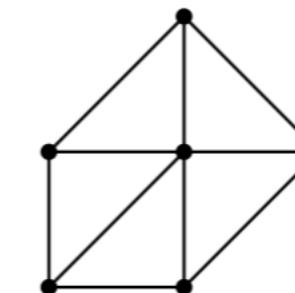
(f)  $(1,3,2,2,3,4)$



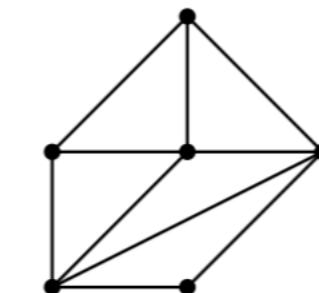
(g)  $(4,1,3,3,1,3)$



(h)  $(3,1,3,2,2,4)$



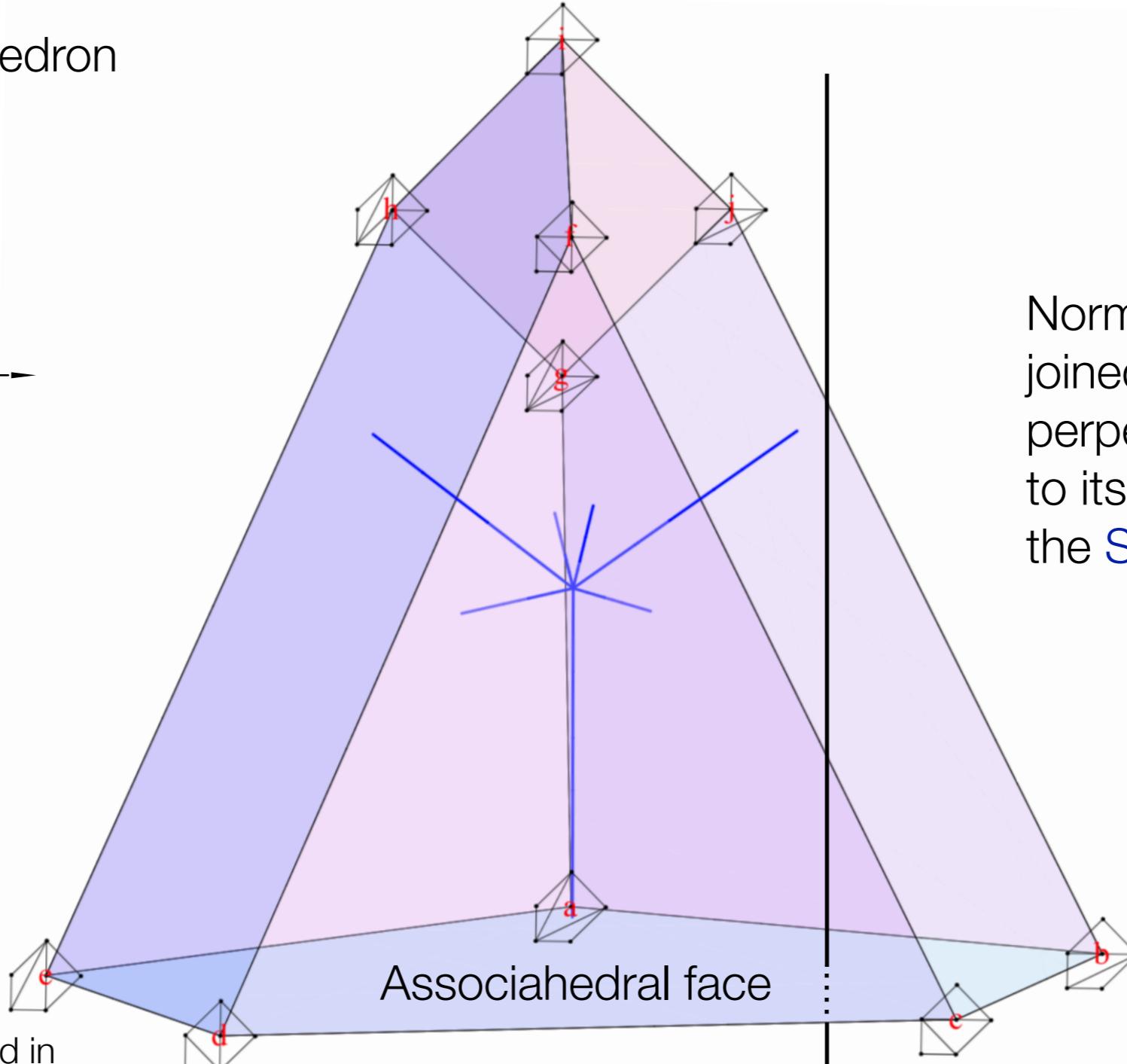
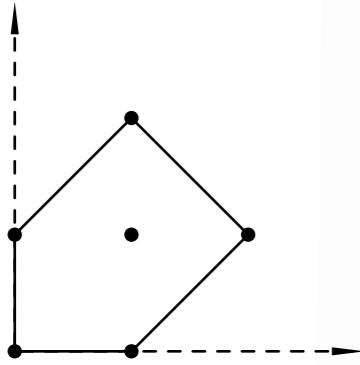
(i)  $(2,2,2,2,2,5)$



(j)  $(3,2,2,3,1,4)$

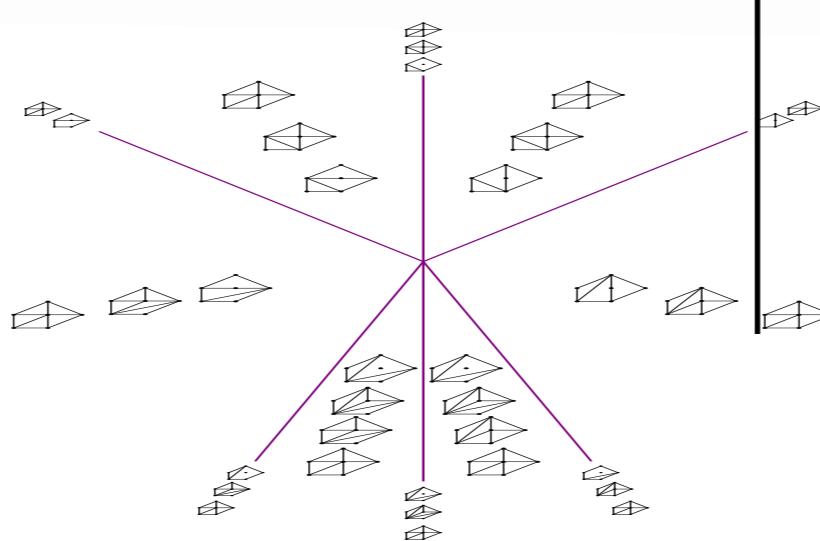
Vector from triangulations are the vertices of the minimal convex hull of all these vectors.

Secondary polyhedron  
of



Normals to its faces,  
joined by wedges  
perpendicular  
to its edges form  
the **Secondary Fan**.

No internal points involved in  
subdivisions on the associahedral  
face

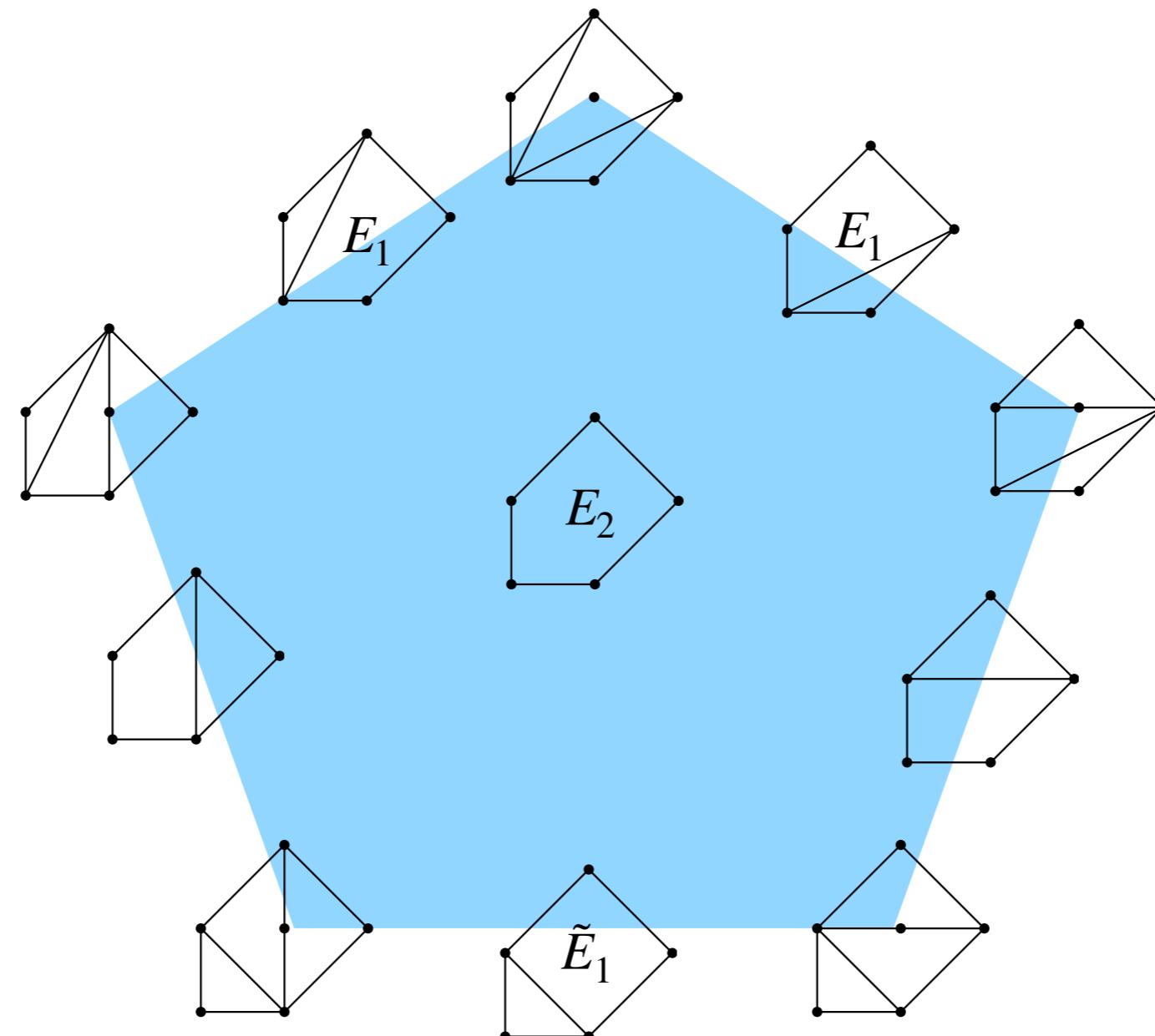


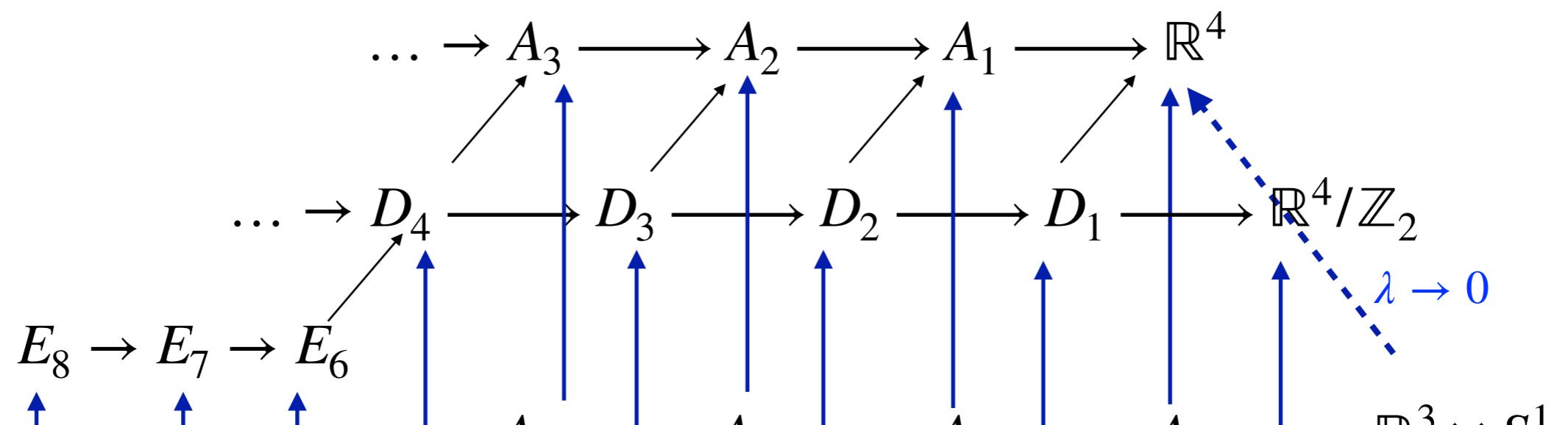
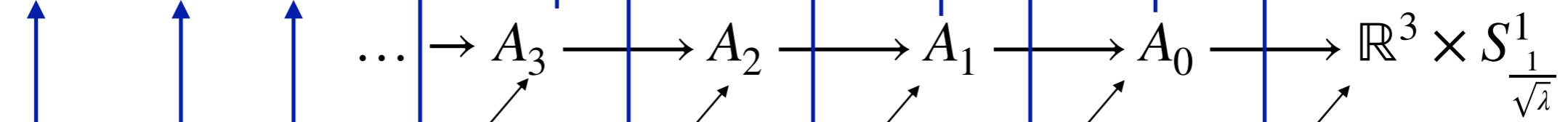
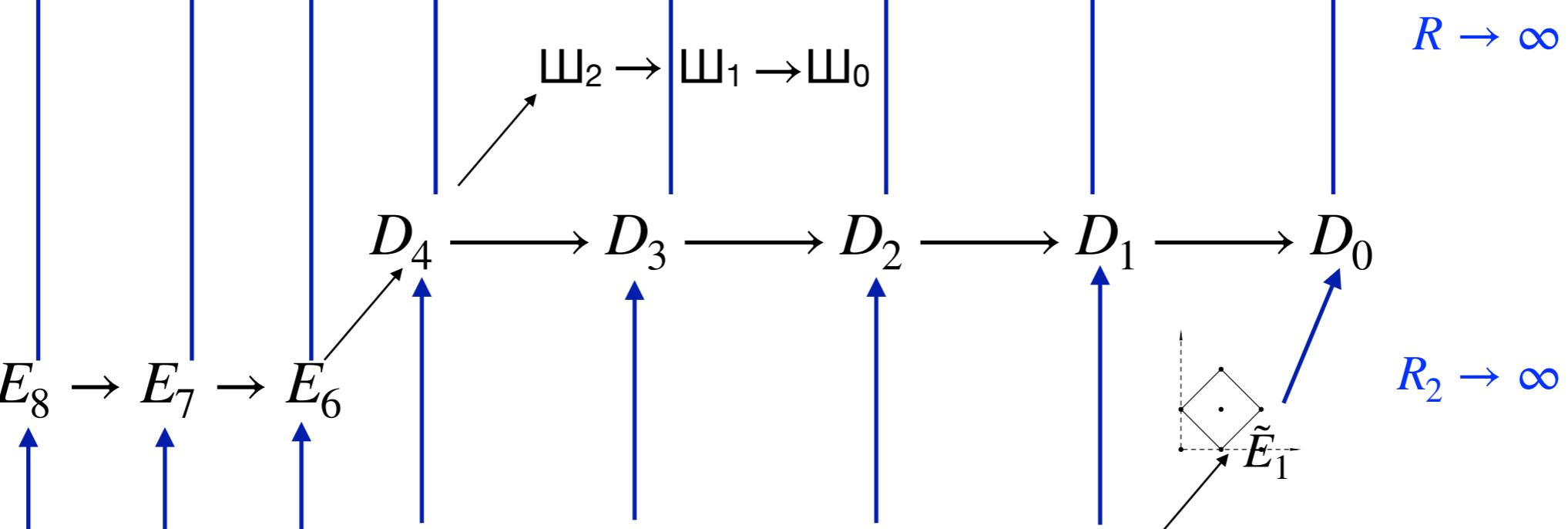
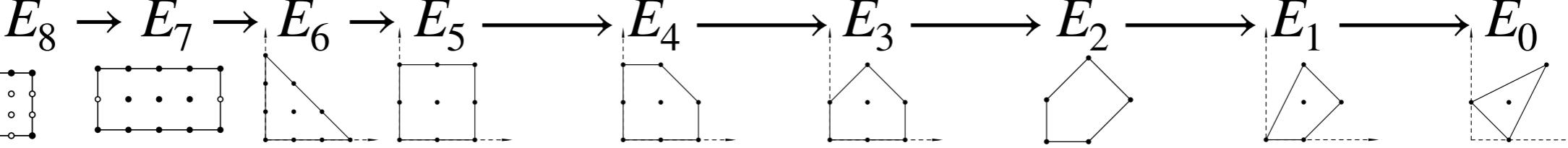
Projection of  
the **Secondary Fan**  
gives the phase space  
of monowall moduli  
spaces.

# Space of all ALH metrics

The parameter space of ALH metrics is fibered over the “universal ALH associahedron”.

For example:



**ALE****ALF****ALG****ALH**

## Summary

- Tesserons (non-compact 4 real-dim. hyperkähler manifolds with finite Pontrjagin number) are classified.
- This classification matches the classification of (discrete) Painlevé equations.
- All of tesserons are moduli spaces of monopoles (or their limits).
- The foremothers of all of these spaces are
  1.  $A_k$  ALF =  $\mathcal{M}$ (one monopole with  $k+1$  simple Dirac singularities),
  2.  $D_k$  ALF =  $\mathcal{M}$ (two centered monopoles with  $k$  simple Dirac singularities),
  3.  $E_8$  ALG =  $\mathcal{M}$ (certain doubly periodic monopole), [Thomas Harris's talk](#)
  4.  $1/2K3$  =  $\mathcal{M}$ (2 monopoles on  $T^3$  with 2+2 simple Dirac singularities).
- This monopole picture leads to the cell structure of the parameter space of all tesserons.