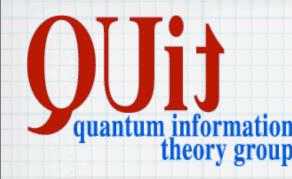
G. M. D'Ariano, arXiv:0807.4383



Probabilistic Theories: what is special about Quantum Mechanics

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BACKROUND

QM is a probabilistic theory + *something more*

OBJECTIVE

to understand what is the *something more*, and derive QM solely from operational principles



Operational framework

Axioms: (primitive notions)

- 🏺 probability ...
- 🏺 events
- independent systems

General principles for mathematical representation

- all mathematical objects must be defined operationally
- mathematical completion is for convenience (e.g. algebraic closure, norm closure, linear span, etc.)

In this talk w.l.g. we consider only:

finite dimension only one kind of "system"





- **NSF:** No signaling from the future.
- **NS:** No signaling (=existence of independent systems)
- **PFAITH:** There exists preparationally faithful states

- AE: Atomicity of evolution
- 🖉 CJ: Choi-Jamiolkowski isomorphism



Postulates under exploration

FAITHE: There exists a faithful effectPURIFY: There exists a purification for each state



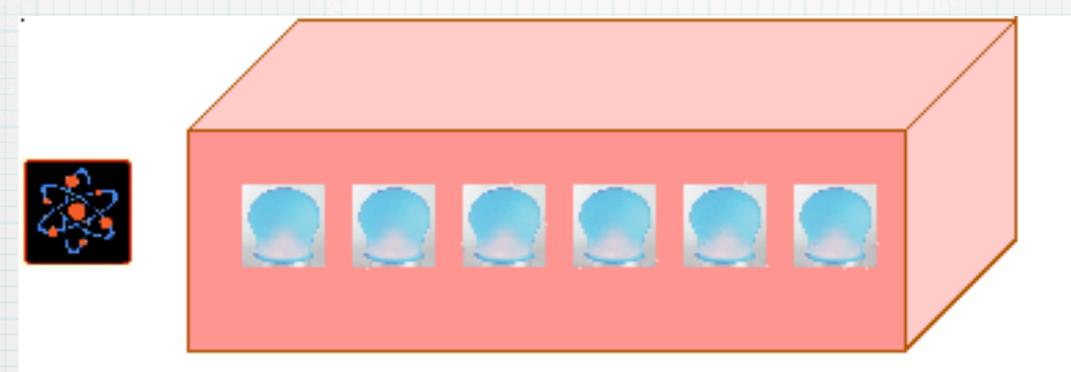
Probabilistic theories





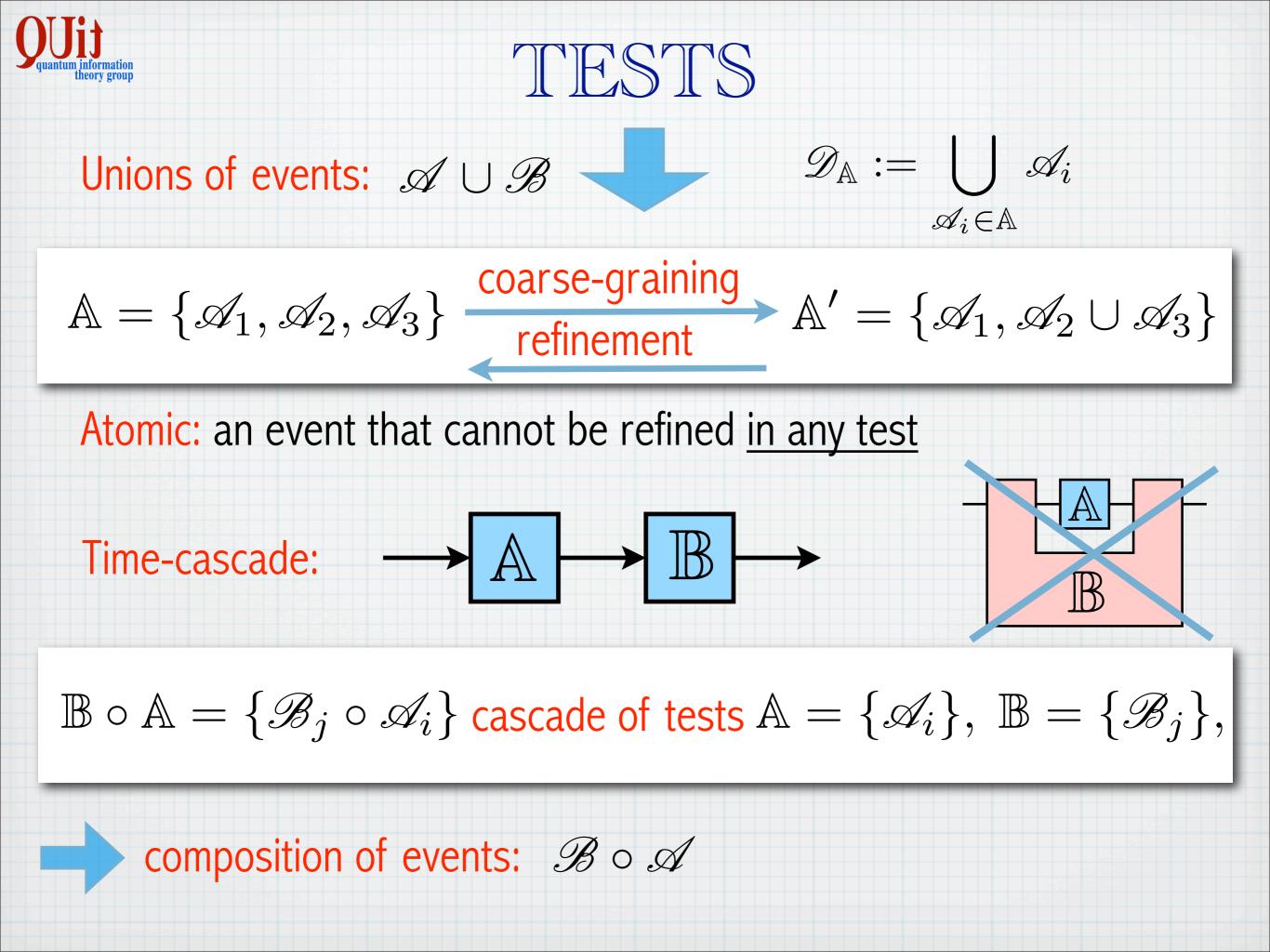


\Im Test/experiment: $\mathbb{A} \equiv \{\mathscr{A}_i\}$ set of possible events \mathscr{A}_i



(deterministic test/transformation: $\mathbb{D} = \{ \mathscr{D} \}$)

Notice: the same event can occur in different tests







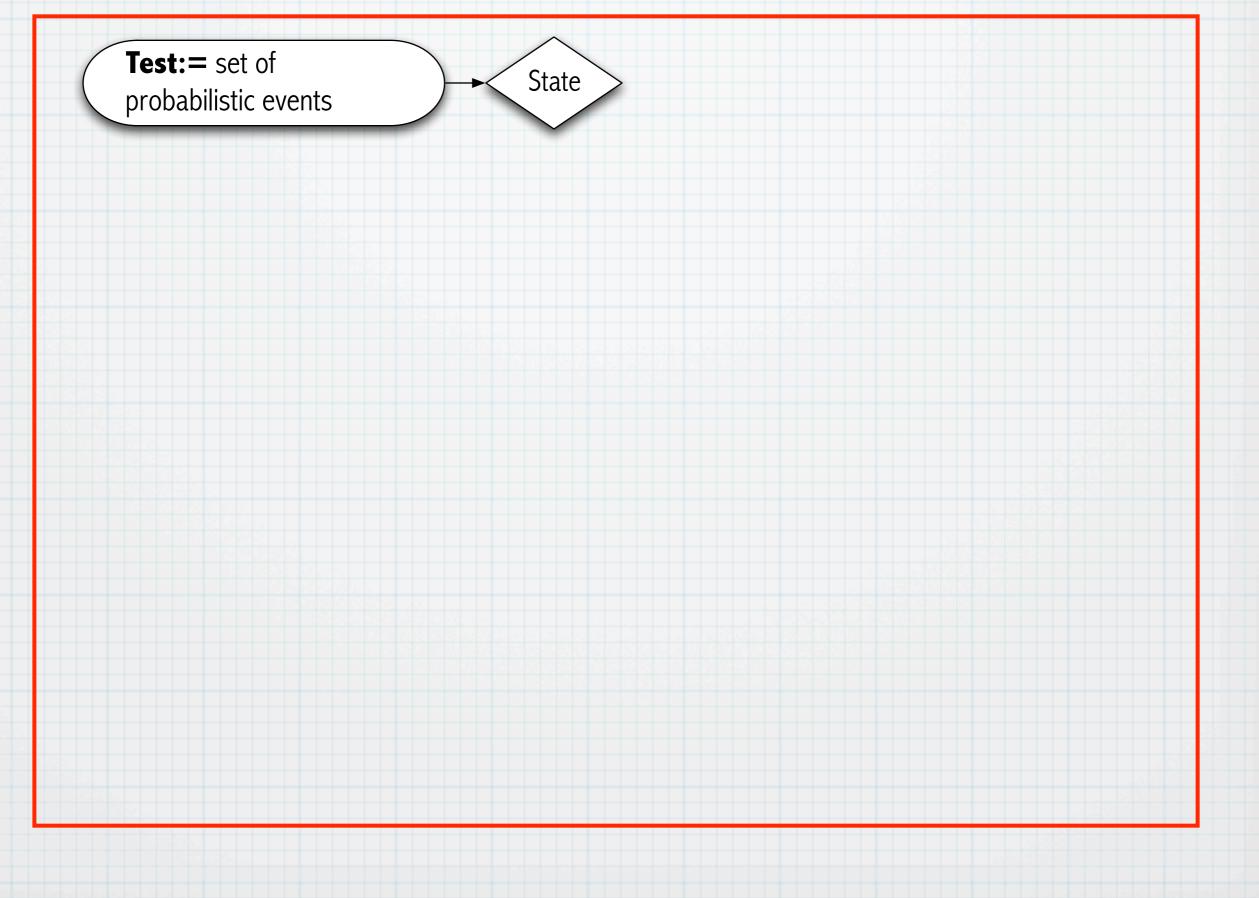
Copenhagen

 $S = \{\mathbb{A}, \mathbb{B}, \mathbb{C}, \ldots\}$ collection of tests closed under

- 🖉 coarse-graining
- conditioning
- cascading (mono-systemic)
- (convex combination)



Probabilistic theories





STATES

State ω : probability rule $\omega(\mathscr{A})$ for any possible event \mathscr{A} in any test

 $S = \{\omega_1, \omega_2, \ldots, A, \mathbb{B}, \mathbb{C}, \ldots\}$

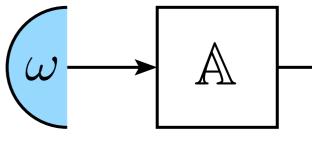
 $(\mathbf{5}$

Normalization: $\sum \omega(\mathscr{A}_j) = 1$

Convex set of states of a system:

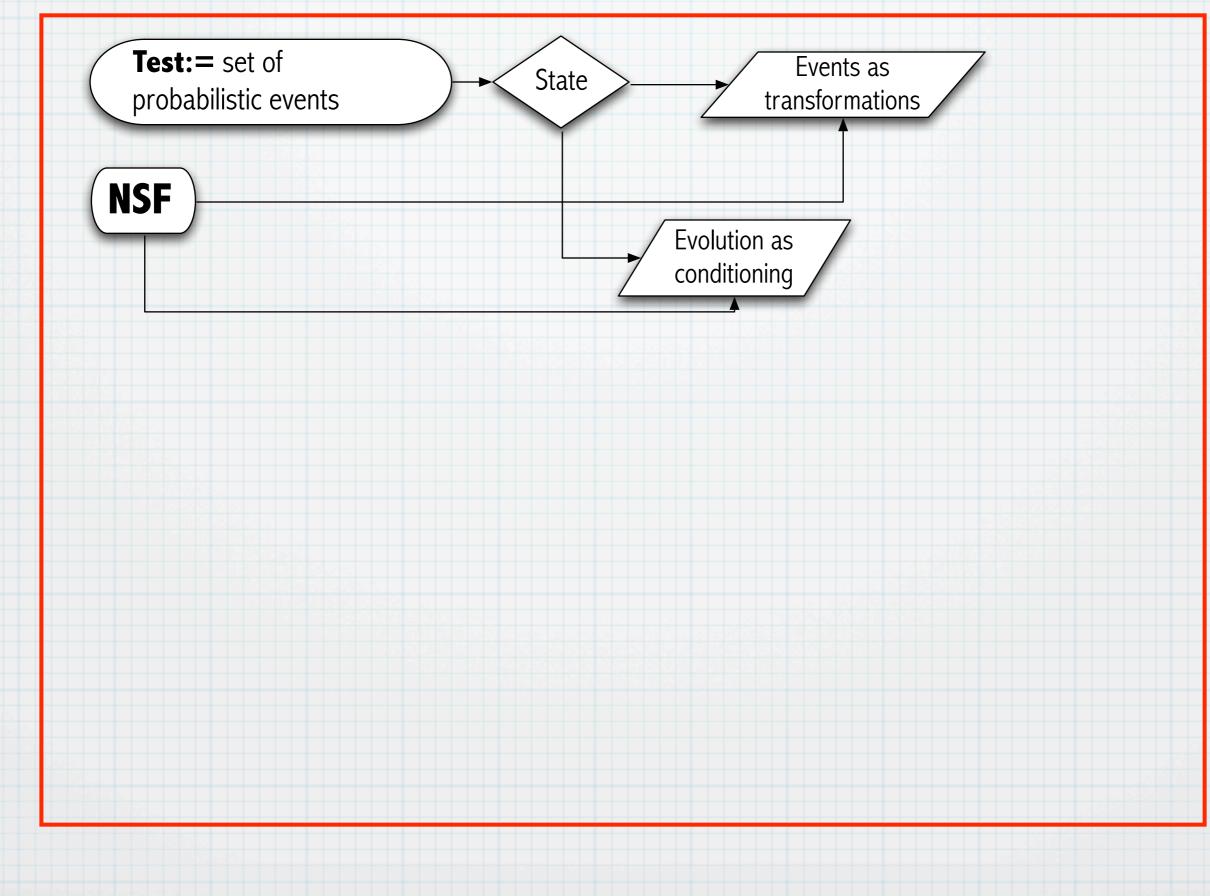
States will also be regarded as tests themselves ("preparation-tests".

 $\mathscr{A}_{i} \in \mathbb{A}$





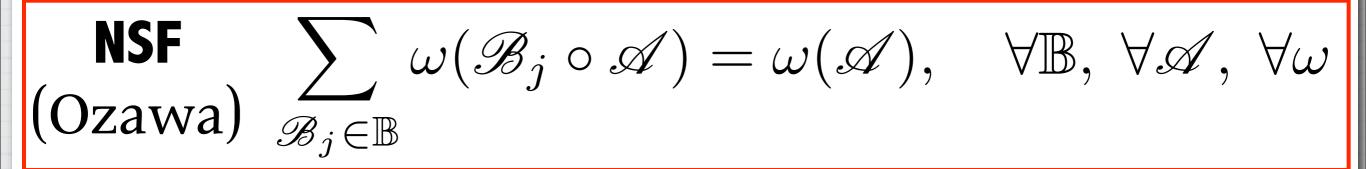
Probabilistic theories





Events \equiv transformations

Cascade: Event $\mathscr{B} \circ \mathscr{A}$: event $\mathscr{B} \in \mathbb{B}$ following $\mathscr{A} \in \mathbb{A}$

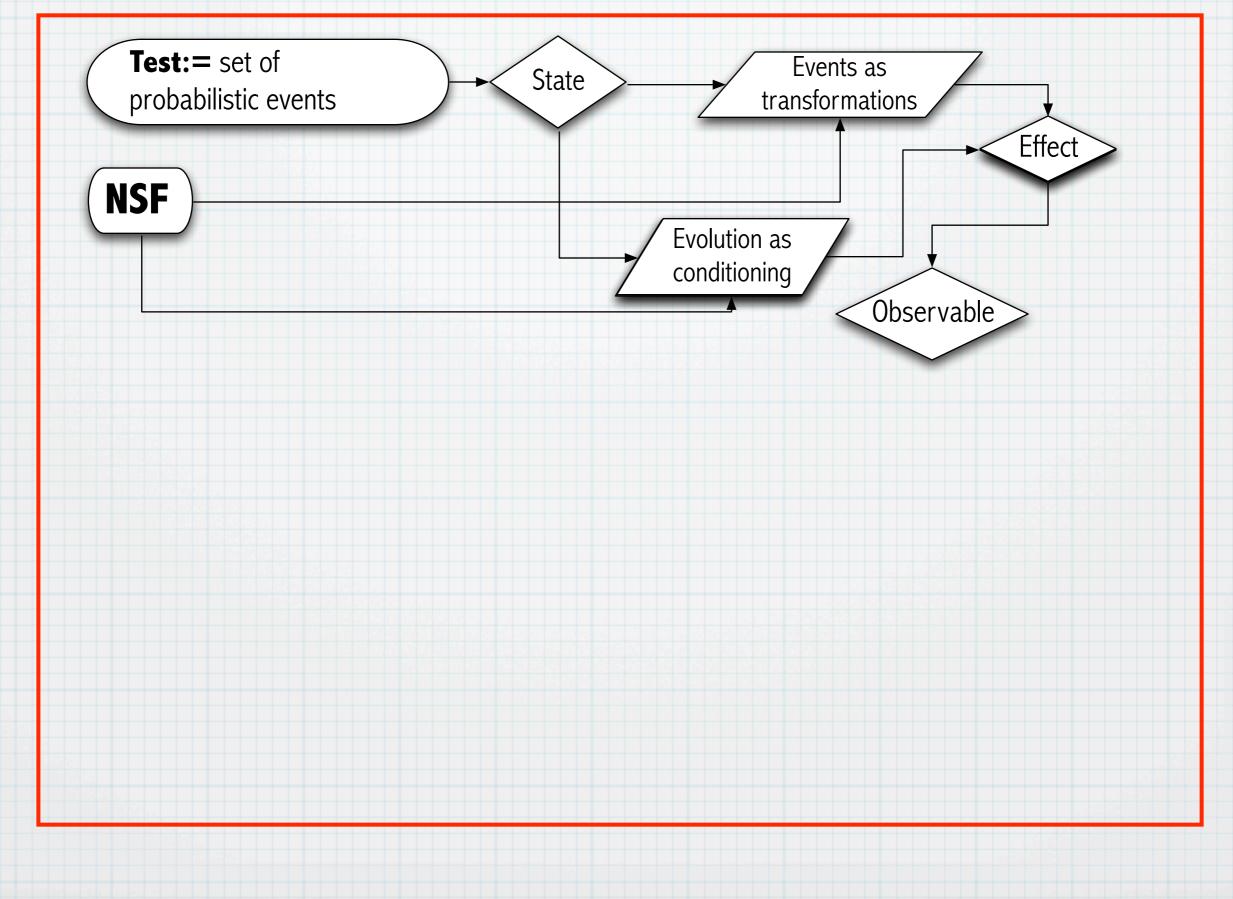


- \Rightarrow conditional probability: $p(\mathcal{B}|\mathcal{A}) = \omega(\mathcal{B} \circ \mathcal{A})/\omega(\mathcal{A})$
- \Rightarrow conditional state: $\omega_{\mathscr{A}} := \omega(\cdot \circ \mathscr{A}) / \omega(\mathscr{A})$
- \Rightarrow evolution \equiv state conditioning: $\mathscr{A}\omega := \omega(\cdot \circ \mathscr{A})$

events \equiv transformations Convex monoid of transformations: \mathfrak{L}



Probabilistic theories





Two transformations \mathscr{A} and \mathscr{B} are conditioning equivalent if

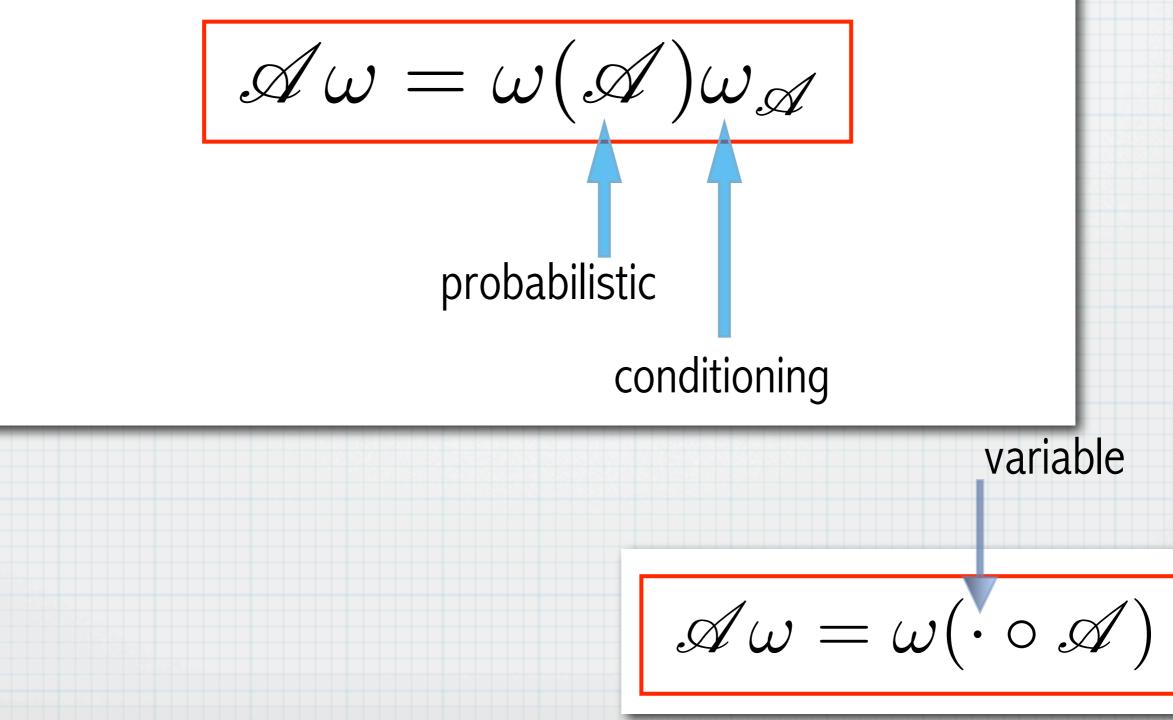
$\omega_{\mathscr{A}} = \omega_{\mathscr{B}} \quad \forall \omega \in \mathfrak{S}$

Conditioning-equivalence class

Two transformations \mathscr{A} and \mathscr{B} are probabilistically equivalent if $\omega(\mathscr{A}) = \omega(\mathscr{B}) \quad \forall \omega \in \mathfrak{S}$

Probabilistic equivalence class







Effects

Effect $\underline{\mathscr{A}}$: equivalence class of transformations occurring with the same probability as \mathcal{A} for all states.

$$\forall \omega \in \mathfrak{S} : \ \omega(\mathscr{A}) \equiv \omega(\mathscr{A})$$

$$a \text{ effect} \longrightarrow \mathscr{A} \in a \text{ means } \omega(\mathscr{A}) \equiv \omega(a)$$

:= convex set of effects

Duality: effects \mathfrak{E} positive linear functionals over states (bounded by 1) $a \in \mathfrak{E}, \ \omega \in \mathfrak{S}, \quad \omega(a) \equiv a(\omega)$

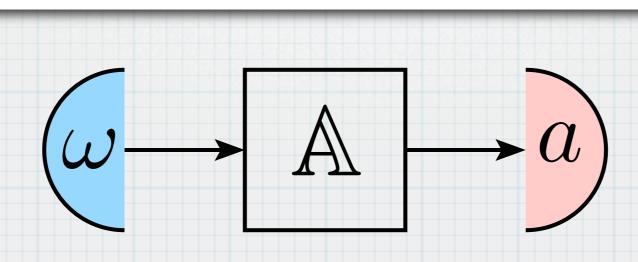
 $e \text{ deterministic effect} \quad \text{i.e.} \quad \omega(e) = 1 \quad \forall \omega \in \mathfrak{S}$





State-conditioning \Rightarrow Transformations act linearly over effects: $\underline{\mathscr{B}} \circ \mathscr{A} \in \underline{\mathscr{B}} \circ \mathscr{A}$ (Heisenberg picture)

Effects will also be regarded as tests themselves: "effect-tests"







Observable $\mathbb{L} = \{l_i\}$: complete set of effects of a test

Normalization: $\sum_{i \in \mathbb{L}} l_i = e$

Informationally complete observable:

$$\mathfrak{E}_{\mathbb{R}} = \mathsf{Span}_{\mathbb{R}}(\mathbb{L})$$

Convex sets, Cones and Linear spaces

Convex set of states:

Convex set of effects:

 \mathfrak{E} , cone: \mathfrak{E}_+

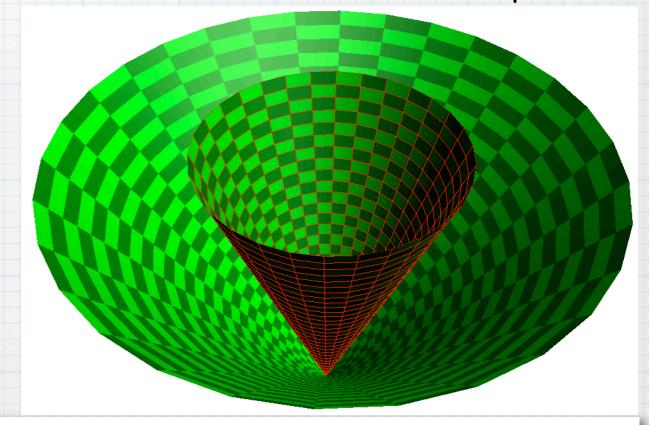
 \mathfrak{S} , cone: \mathfrak{S}_+

 \mathfrak{T} , cone: \mathfrak{T}_+

Convex monoid of transformations:

Linear spaces:

$$\begin{split} \mathfrak{S}_{\mathbb{R}} &= \mathsf{Span}_{\mathbb{R}}\mathfrak{S} \\ \mathfrak{S}_{\mathbb{C}} &= \mathsf{Span}_{\mathbb{C}}\mathfrak{S} \\ \mathfrak{E}_{\mathbb{R}}, \mathfrak{E}_{\mathbb{C}}, \mathfrak{T}_{\mathbb{R}}, \mathfrak{T}_{\mathbb{C}} \end{split}$$

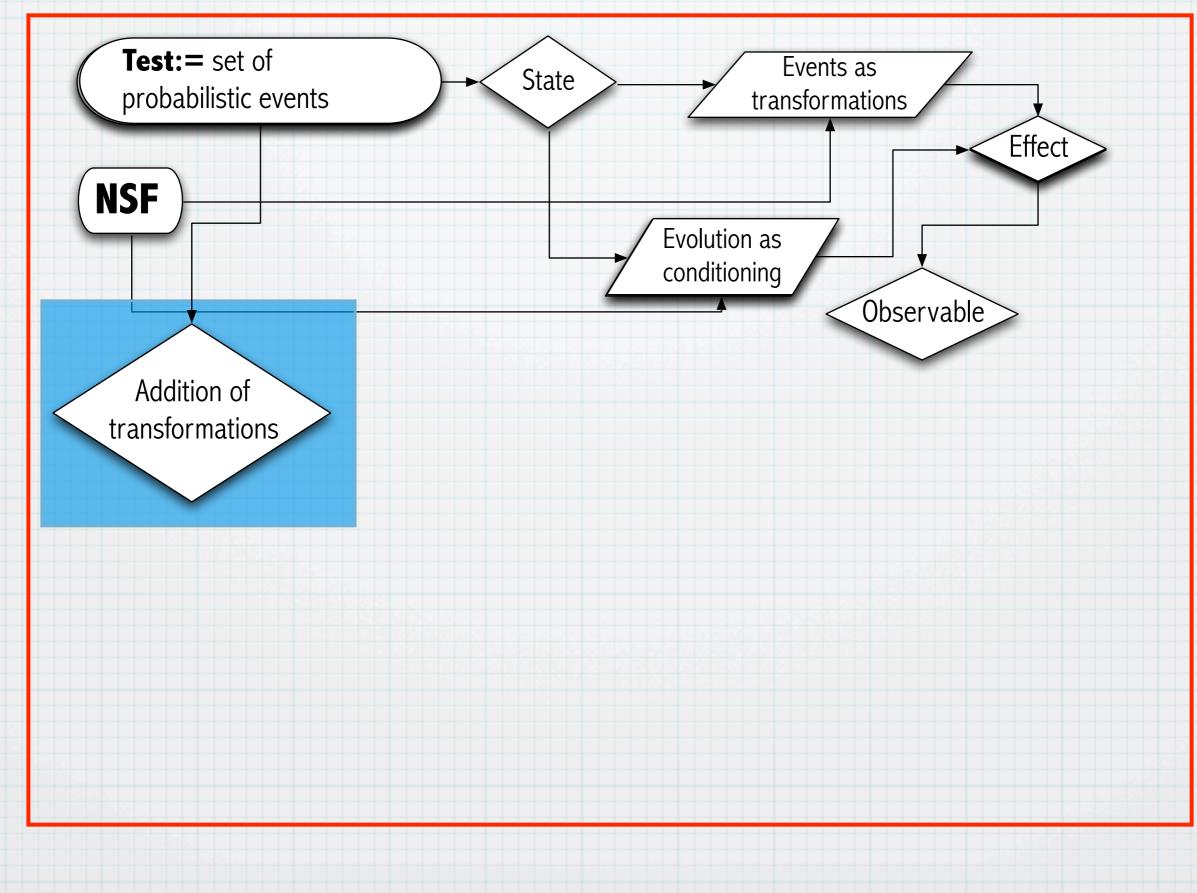


No-restriction hypotesis: (no limitations to preparability)

 $\mathfrak{S}_+ = (\mathfrak{E}_+)^*$



Probabilistic theories



OUII guantum information

Addition of transformations

Two transformations* \mathscr{A} and \mathscr{B} are test-compatible if for every state ω one has

$$\omega(\mathscr{A}) + \omega(\mathscr{B}) \le 1$$

For any two test-compatible transformations \mathscr{A}_1 and \mathscr{A}_2 we define the transformation $\mathscr{A}_1 + \mathscr{A}_2$ as the union event $\mathscr{A}_1 \cup \mathscr{A}_2$ (the apparatus signals that either \mathscr{A}_1 or \mathscr{A}_2 occurred)

$$\begin{split} &\omega(\mathscr{A}_1 + \mathscr{A}_2) = \omega(\mathscr{A}_1) + \omega(\mathscr{A}_2) & \text{(probabilistic class)} \\ &(\mathscr{A}_1 + \mathscr{A}_2)\omega = \mathscr{A}_1\omega + \mathscr{A}_2\omega & \text{(conditioning class)} \end{split}$$

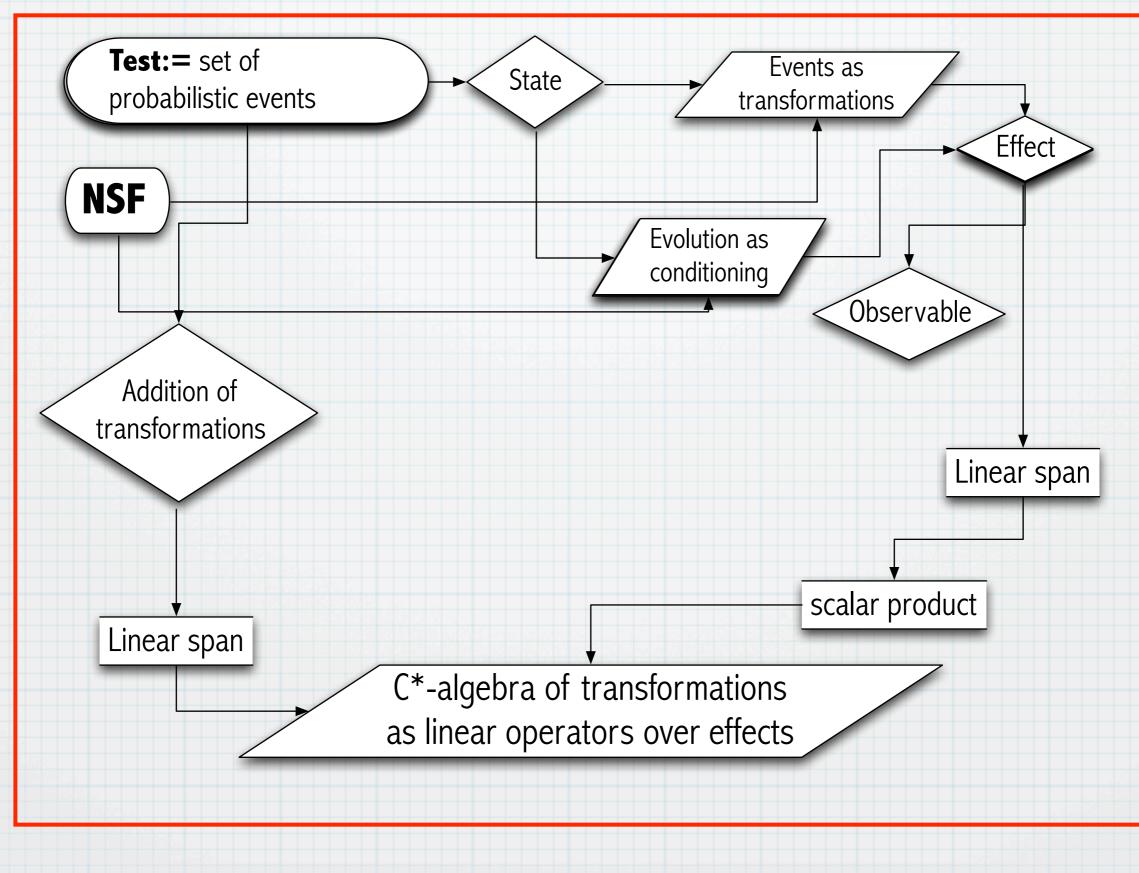
 $\omega_{\mathscr{A}_1+\mathscr{A}_2} = \frac{\omega(\mathscr{A}_1)}{\omega(\mathscr{A}_1+\mathscr{A}_2)}\omega_{\mathscr{A}_1} + \frac{\omega(\mathscr{A}_2)}{\omega(\mathscr{A}_1+\mathscr{A}_2)}\omega_{\mathscr{A}_2}$

 $\omega(b \circ (\mathscr{A}_1 + \mathscr{A}_2)) = \omega(b \circ \mathscr{A}_1) + \omega(b \circ \mathscr{A}_2), \quad \forall b \in \mathfrak{E}, \ \forall \omega \in \mathfrak{S}$

(*) occurring also in different tests



Probabilistic theories





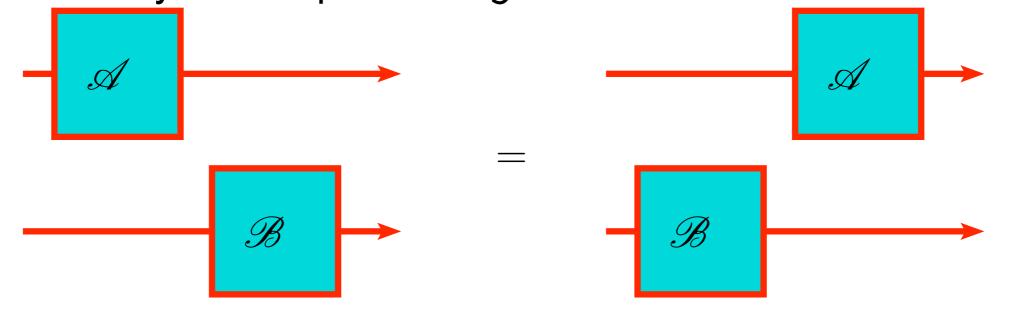
Transformations/events are linear maps over effects, i.e. they make a matrix algebra over effects

One can introduce a scalar product over effects ... \Rightarrow transformations become a C*-algebra ...

OUIT quantum information theory group

INDEPENDENT SYSTEMS

Two systems are independent if on each system it is possible to perform local tests for which on every joint state one has the commutativity of the pertaining transformations



 $A^{(1)} \circ \mathcal{B}^{(2)} = \mathcal{B}^{(2)} \circ \mathcal{A}^{(1)}$

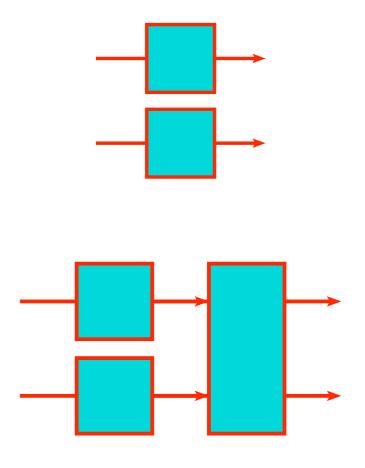
 $(\mathscr{A}, \mathscr{B}, \mathscr{C}, \ldots) \doteq \mathscr{A}^{(1)} \circ \mathscr{B}^{(2)} \circ \mathscr{C}^{(3)} \circ \ldots$

 $[(\mathscr{A}, \mathscr{B}, \mathscr{C}, \ldots)]_{\text{eff}} \equiv (\underline{\mathscr{A}}, \underline{\mathscr{B}}, \underline{\mathscr{C}}, \ldots)$

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COMPOSTING SYSTEMS

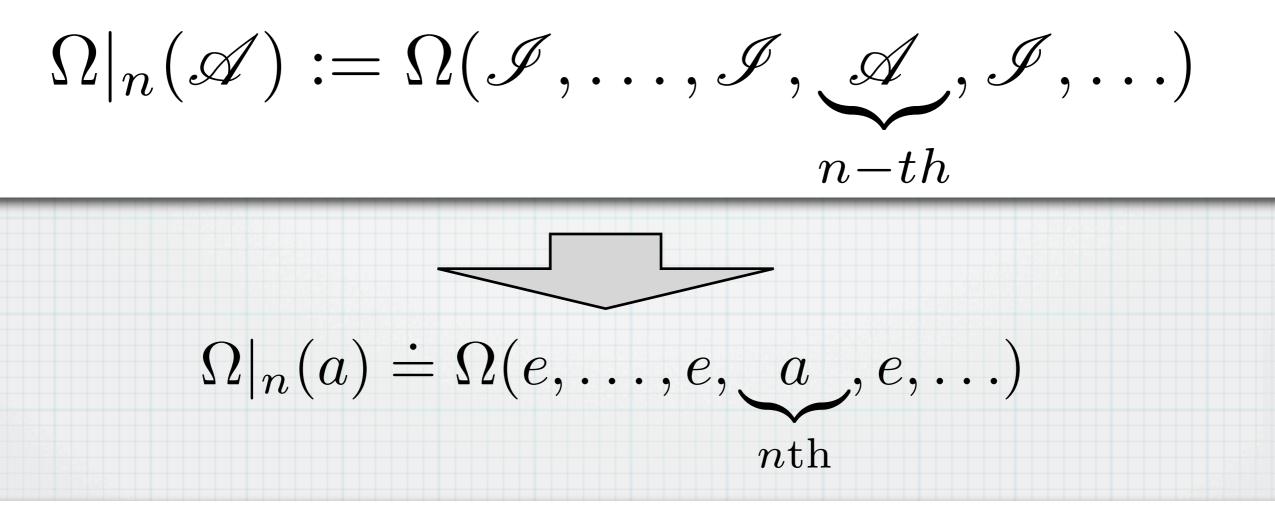
We compost the two systems S_1 and S_2 into the bipartite system $S_1 \odot S_2$ by embedding the local tests $S_1 \times S_2$ into the bipartite system $S_1 \odot S_2$ as $S_1 \odot S_2 \supseteq S_1 \times S_2$ and closing w.r.t. coarse graining, convex combination and cascading. Nonlocal tests: $S_1 \odot S_2 \setminus S_1 \times S_2$



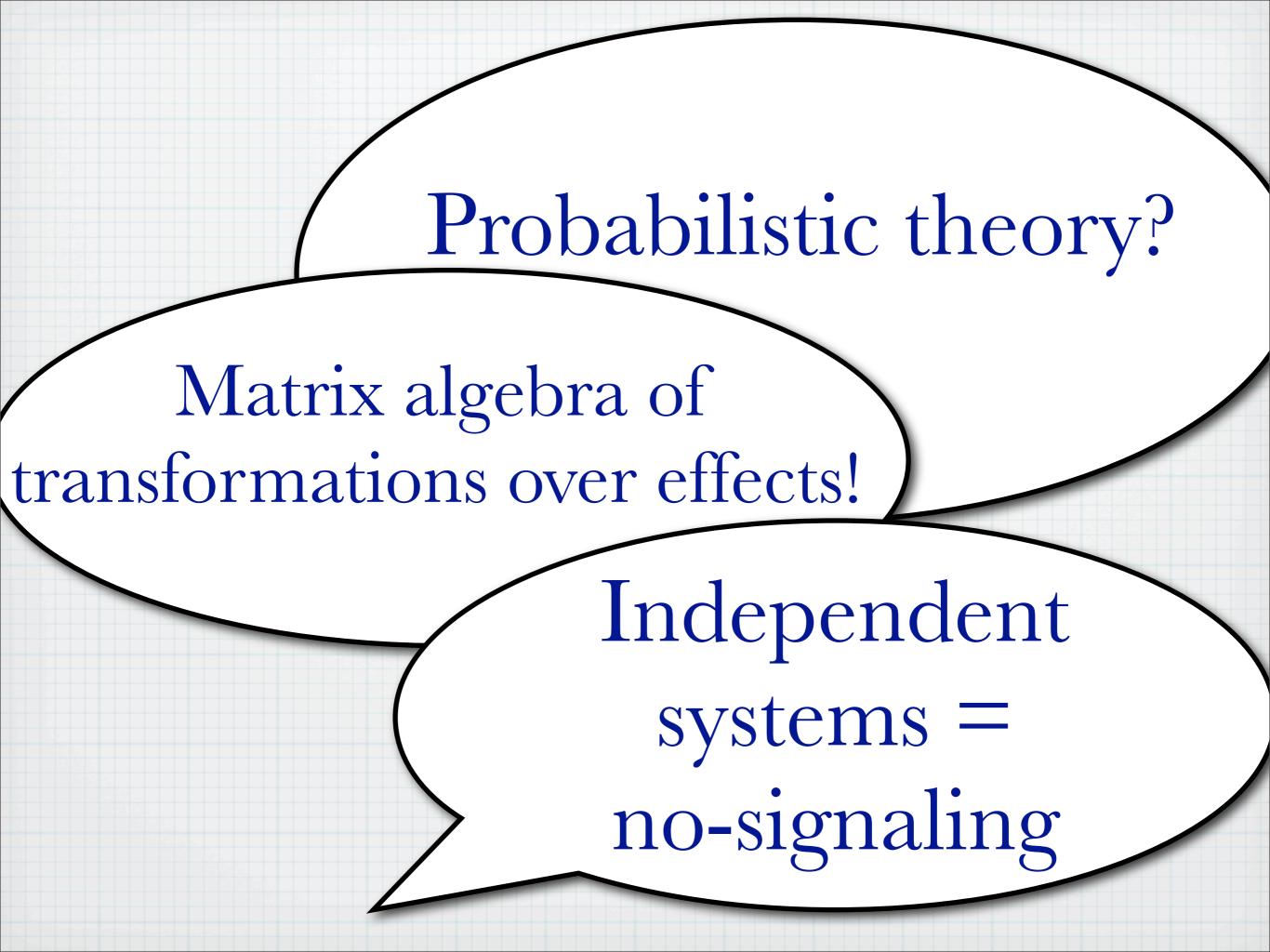
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MARGINAL STATE

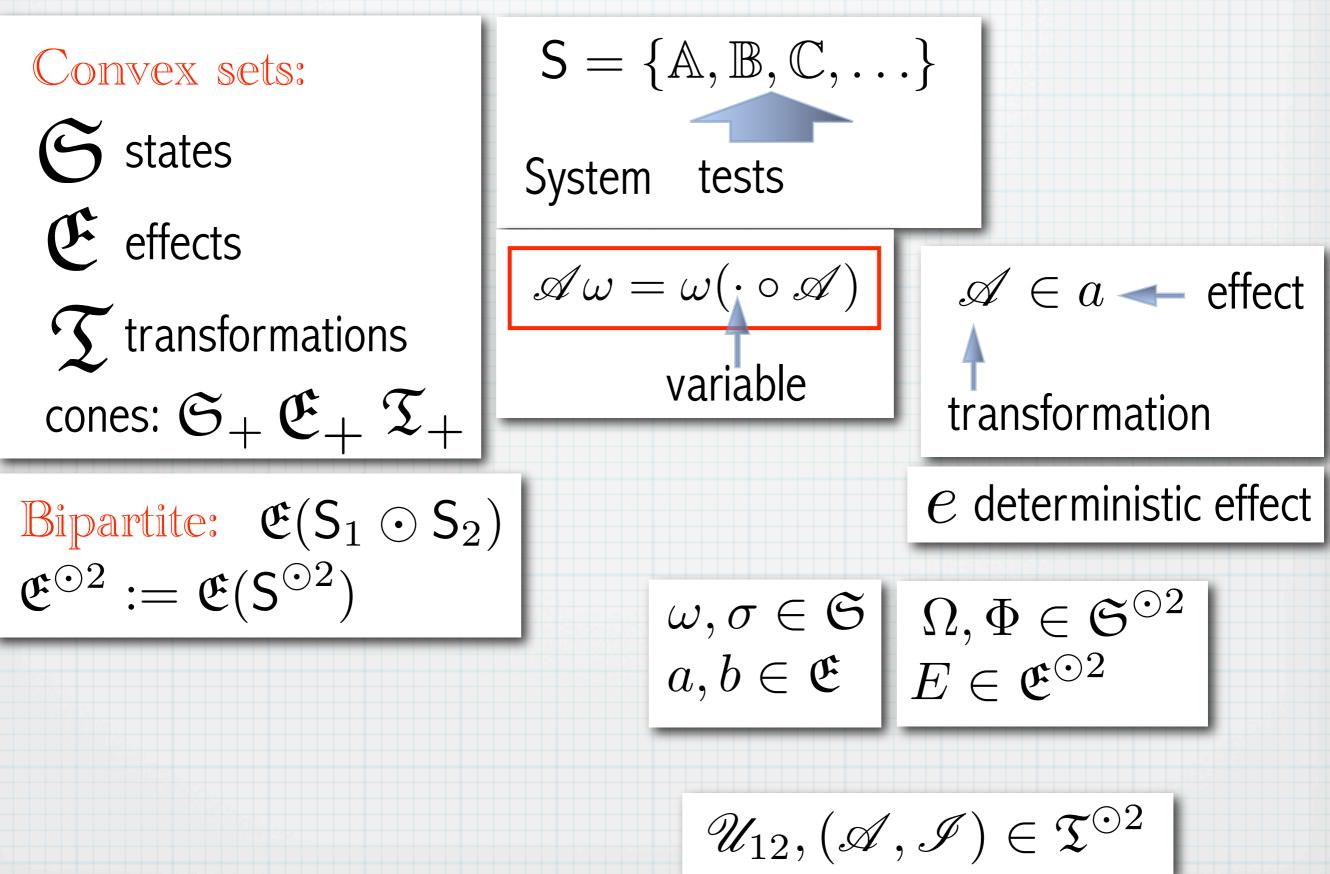
For a multipartite system we define the marginal state $\Omega|_n$ of the n-th system the state that gives the probability of any local transformation \mathscr{A} on the n-th system with all other systems untouched, namely

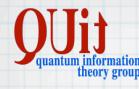


NS: (no-signaling) any local test on a system is equivalent to no-test on an another independent system.

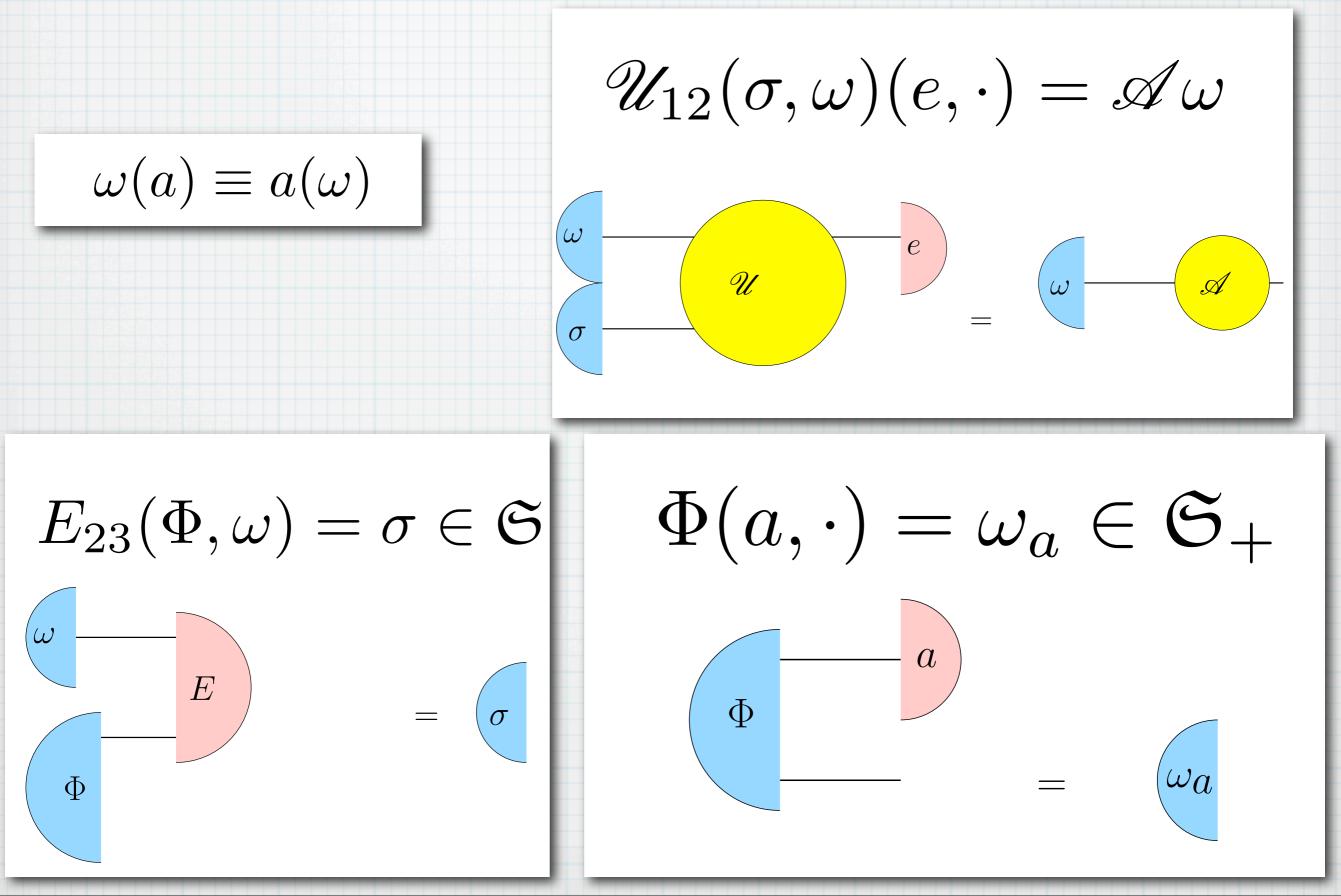


Review of notation





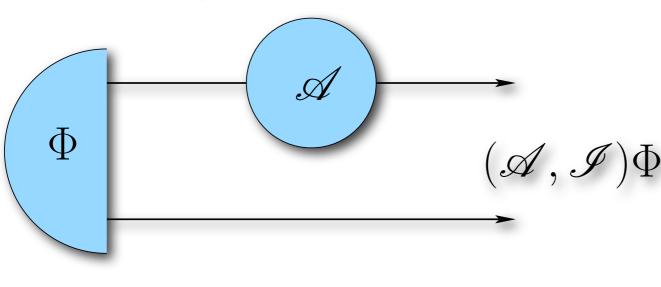
Review of notation



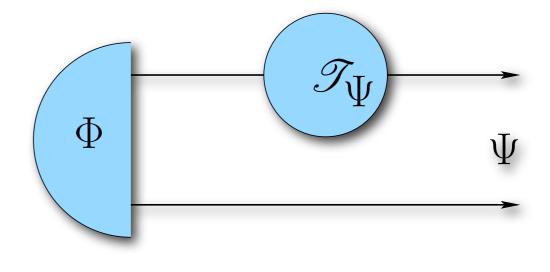


FAITHFUL STATES

A state Φ of a bipartite system is dynamically faithful when the output state $(\mathscr{A}, \mathscr{I})\Phi$ from a local transformation \mathscr{A} on one system is in 1-to-1 correspondence with the transformation \mathscr{A}



A state Φ of a bipartite system is preparationally faithful if every joint state Ψ can be achieved by a suitable local transformation \mathscr{T}_{Ψ} on one system occurring with nonzero probability





Postulate PFAITH

PFAITH: For any couple of identical systems, there exist a symmetric^{*} state Φ that is preparationally faithful.

(*) under permutation of the two systems

Theorem: $\Phi\,$ is also dynamically faithful.



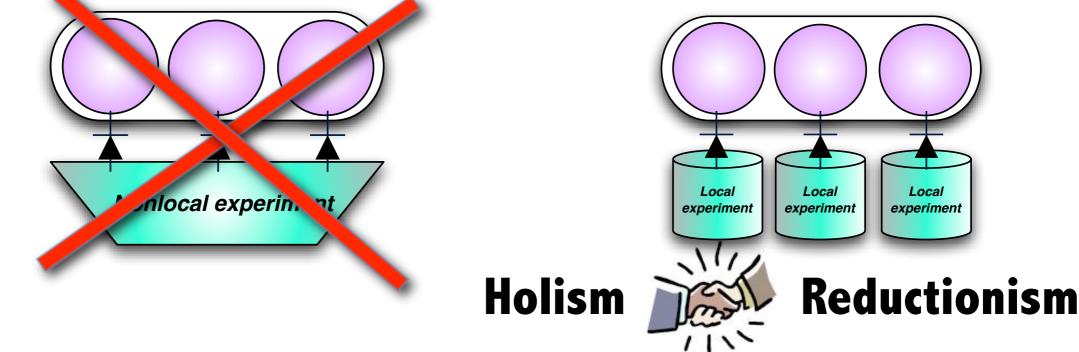
Consequences of PFAITH

Calibrability & Preparability

Impossibility of secure bit commitment

Marginal state $\chi=\Phi(e,\cdot)$ internal and invariant under a transposed channel

Local observability: There exist global info-complete observables made of local info-complete



$$\begin{split} & \overbrace{\mathbb{C}}^{e} \\ & \overbrace{\mathbb{C}}^{e} (S^{\odot^2}) \simeq \mathfrak{T}_{\mathbb{F}}(S) \\ & \overbrace{\mathbb{F}} = \mathbb{R}, \mathbb{C} \\ & \overbrace{\mathbb{C}}^{e} + \ni a \mapsto \omega_a = \Phi(a, \cdot) \in \mathfrak{S}_+ \\ & \overbrace{\mathbb{C}}^{e} = \omega_a \\ & \overbrace{\mathbb{C}}^{e} \\ & \overbrace{\mathbb{C}}^{e} = (S^{\odot^2}) = \mathfrak{E}_{\mathbb{F}}(S)^{\otimes 2} \\ & \overbrace{\mathbb{C}}^{e} (S^{\odot^2}) = \mathfrak{S}_{\mathbb{F}}(S)^{\otimes 2} \\ & \overbrace{\mathbb{C}}^{e} (S^{\odot^2}) = \mathfrak{S}_{\mathbb{F}}(S)^{\otimes 2} \\ & \overbrace{\mathbb{C}}^{e} = \mathbb{C}^{e} \\ & \overbrace{\mathbb{C}}^{e} = \mathbb{C}^{E$$

There exist states that are purifiable (e.g. $\mathscr{A}\chi$, with \mathscr{A} atomic)



The faithful state Φ provides a non-degenerate scalar product over effects via its Jordan form (ζ Jordan involution):

$$\forall a, b \in \mathfrak{E}_{\mathbb{R}}, \quad \Phi(b|a)_{\Phi} := |\Phi|(b,a) = \Phi(\varsigma(b),a)$$

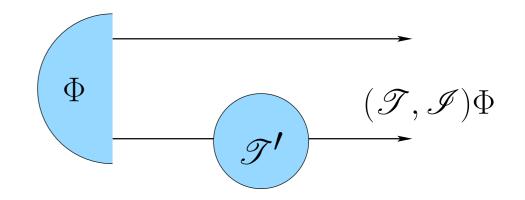
It allows to introduce an operational notion of transposition for transformations: $1 - (\mathcal{A} + \mathcal{B})' - \mathcal{A}' + \mathcal{B}'$

$$(\mathscr{T},\mathscr{I})\Phi=(\mathscr{I},\mathscr{T}')\Phi$$

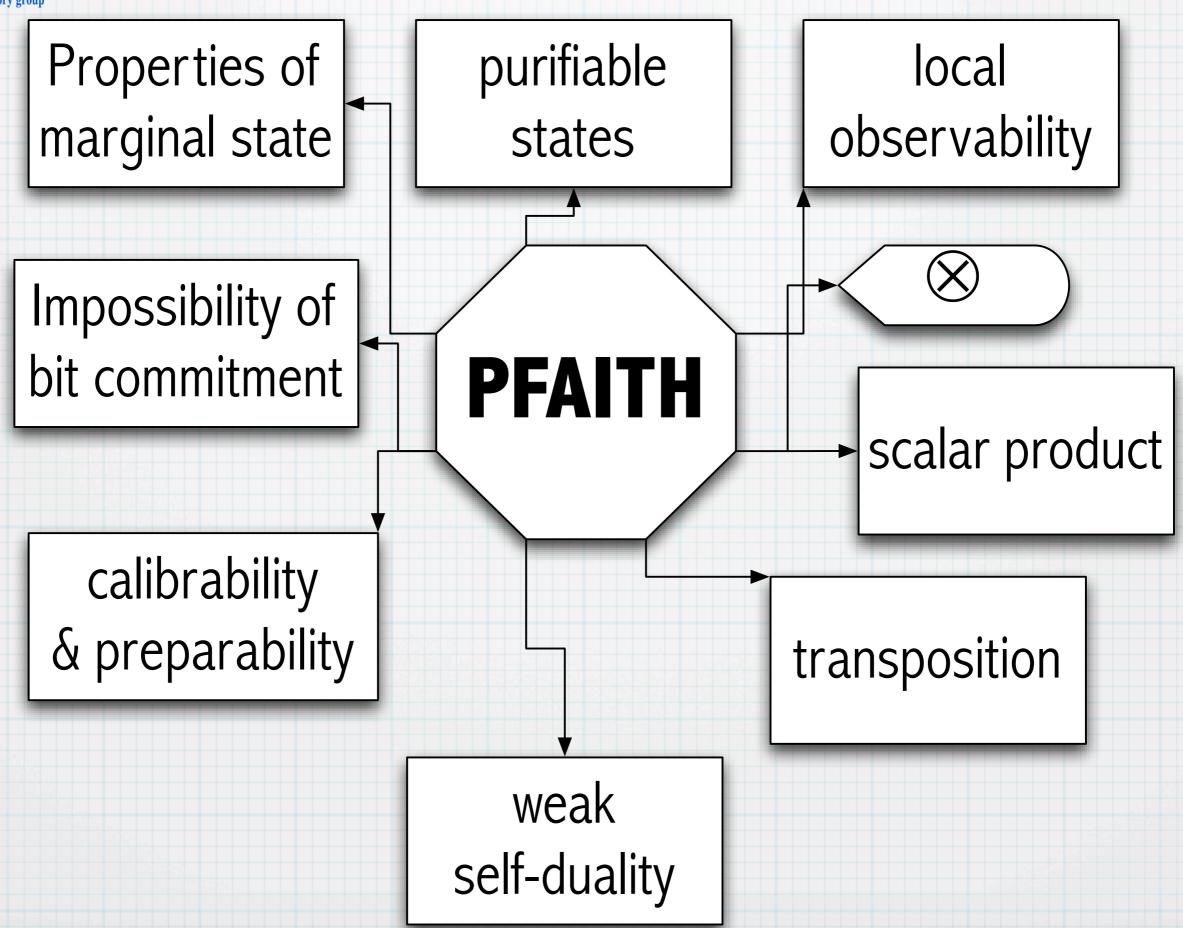
1.
$$(\mathscr{A} + \mathscr{B})' = \mathscr{A}' + \mathscr{B}'$$

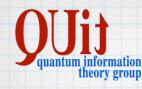
2. $(\mathscr{A}')' = \mathscr{A},$

3.
$$(\mathscr{A} \circ \mathscr{B})' = \mathscr{B}' \circ \mathscr{A}'$$









INTERLUDE

Exploring Postulates: FAITHE and PURIFY



Φ

a

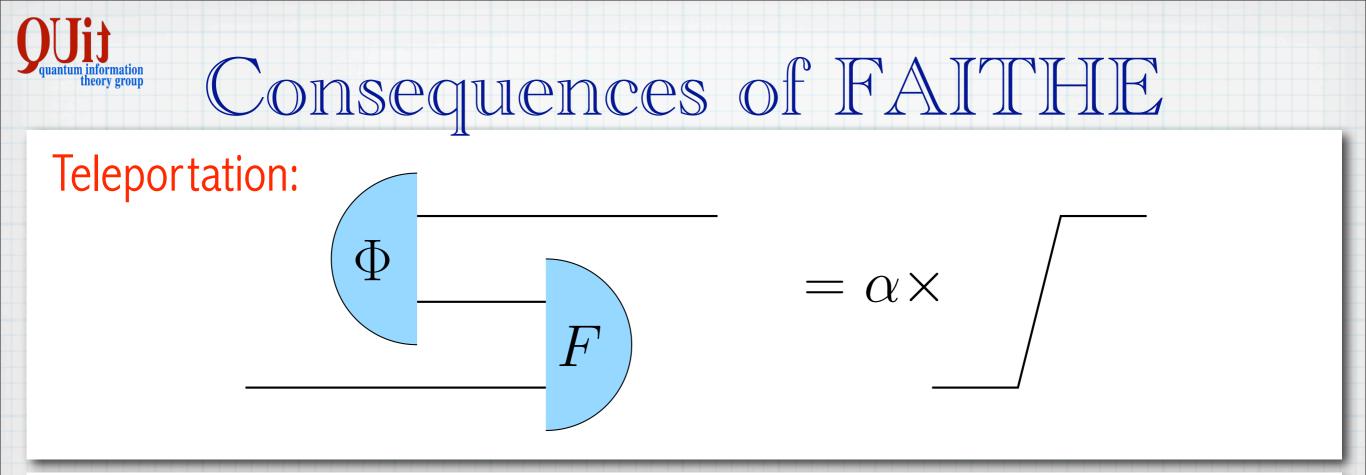
Faithful effect

Remind the cone-isomorphism from the faithful state Φ

$$\mathfrak{E}_+ \ni a \mapsto \omega_a = \Phi(a, \cdot) \in \mathfrak{S}_+$$

FAITHE: There exist a bipartite effect F achieving the inverse of the isomorphism $a \mapsto \omega_a = \Phi(a, \cdot)$ namely:

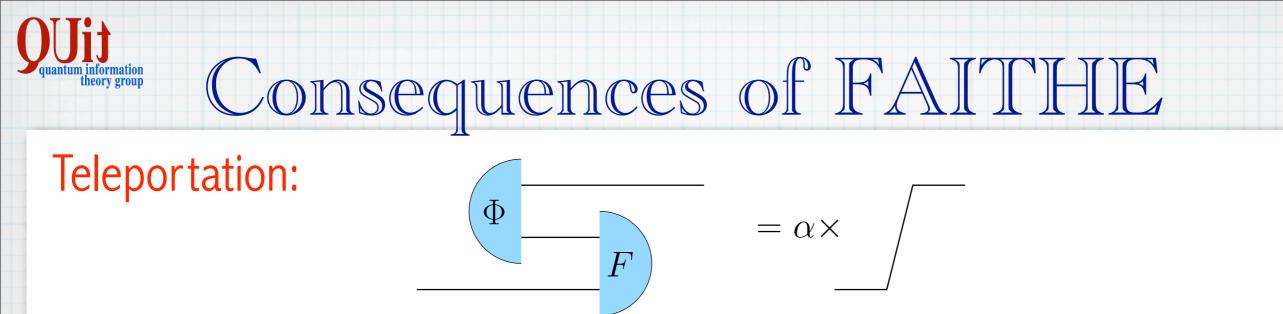
$$F_{23}(\omega_a)_2 = F_{23}\Phi_{12}(a, \cdot) = \alpha a_3, \qquad 0 < \alpha \leq 1$$



 $F \text{ is completely faithful, i.e. } F_{\mathscr{A}} := F \circ (\mathscr{I}, \mathscr{A}) \iff \mathscr{A}$ realizes the cone-isomorphism: $\mathfrak{E}_+(S^{\odot 2}) \simeq \mathfrak{T}_+(S)$

$$\mathfrak{E}_{+}(\mathsf{S}^{\odot 2}) \ni A \mapsto \Omega_{A} := A_{23}(\Phi, \Phi) \in \mathfrak{S}(\mathsf{S}^{\odot 2})$$

realizes the cone-isomorphism: $\mathfrak{S}_{+}(\mathsf{S}^{\odot 2}) \simeq \mathfrak{E}_{+}(\mathsf{S}^{\odot 2})$



$$\alpha(\mathsf{S}) = \max_{E \in \mathfrak{E}(\mathsf{S}^{\odot 2})} \{ (\Phi, \Phi)(e, E, e) \}$$

is a property of the system and depends on the particular probabilistic theory

In Quantum Mechanics: $\alpha = \dim(\mathsf{H})^{-2}$ $\omega_a = \sqrt{\alpha} \varsigma(a)$ $(\cdot, F)(\Phi, \cdot) = \sqrt{\alpha} |\Phi|$



Exploring PURIFY

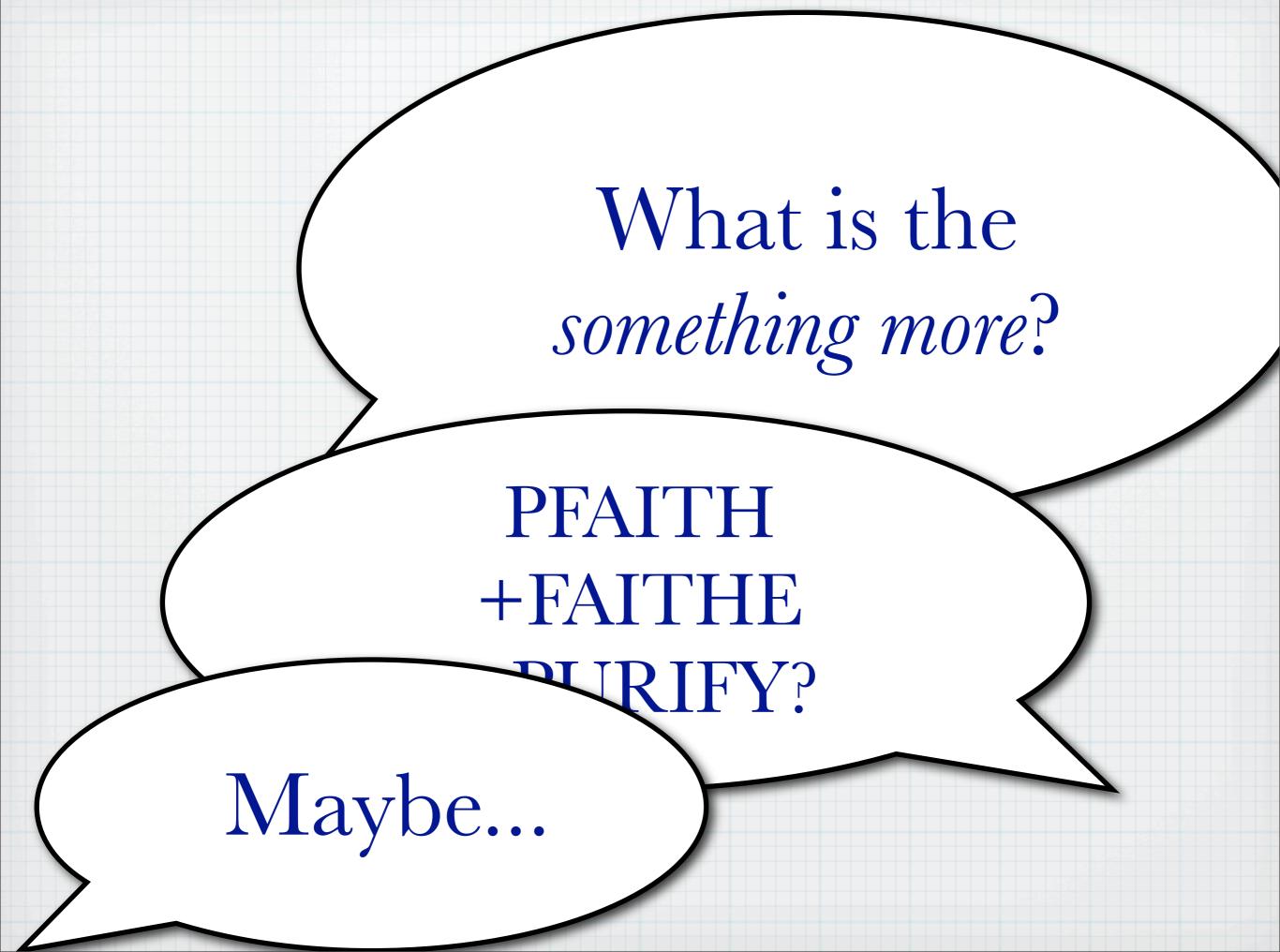
PURIFY: Every state has a purification on two identical systems.

Each state can be obtained by applying an atomic transformation to the marginal state $\,\chi=\Phi(e,\cdot)\,$

Each effect contains an atomic transformation.

 Φ is pure.

 \mathscr{I} is atomic.



What is the something more?

It must give that: effects make a C*-algebra

Reconstructing QM from probabilities

The axiomatic short-circuit of CJ+AE

Quantum Tomography for Measuring Experimentally the Matrix Elements of an Arbitrary Quantum Operation

G. M. D'Ariano and P. Lo Presti

at our disposal a general method for experimentally determining the quantum operation matrix, using any available quantum-tomographic scheme for the system in consideration, and a single fixed state at the input, which is an entangled (not even maximally) state. In the optical domain we show that one can achieve the tomographic reconstruction of the operation using exactly the same apparatus of the recently performed experiment of Ref. [9].

Let us consider for simplicity a "pure" quantum operation in the form (5). Given an orthonormal basis $\{|j\rangle\}$ corresponding to some physical observable, how can we determine the matrix $A_{ij} = \langle i|A|j \rangle$ experimentally? Instead of acting with the contraction A on an "isolated" system, we perform the map on a system which is entangled in the state $|\psi\rangle\rangle \in \mathcal{H} \otimes \mathcal{H}$ with an identical system; i.e.,

$$|\psi\rangle\rangle \rightarrow |\phi\rangle\rangle = \frac{A \otimes I|\psi\rangle\rangle}{\|A\psi\|_{HS}}.$$
 (6)

With the double ket we denote bipartite vectors $|\psi\rangle\rangle \in \mathcal{H} \otimes \mathcal{H}$, which, keeping the basis $\{|j\rangle\}$ as fixed, are in one-to-one correspondence with matrices as follows:

$$|\psi\rangle\rangle = \sum_{ij} \psi_{ij} |i\rangle \otimes |j\rangle.$$
⁽⁷⁾

$$A_{ij} = \kappa \langle E_{ij}(\psi) \rangle, \qquad (10)$$

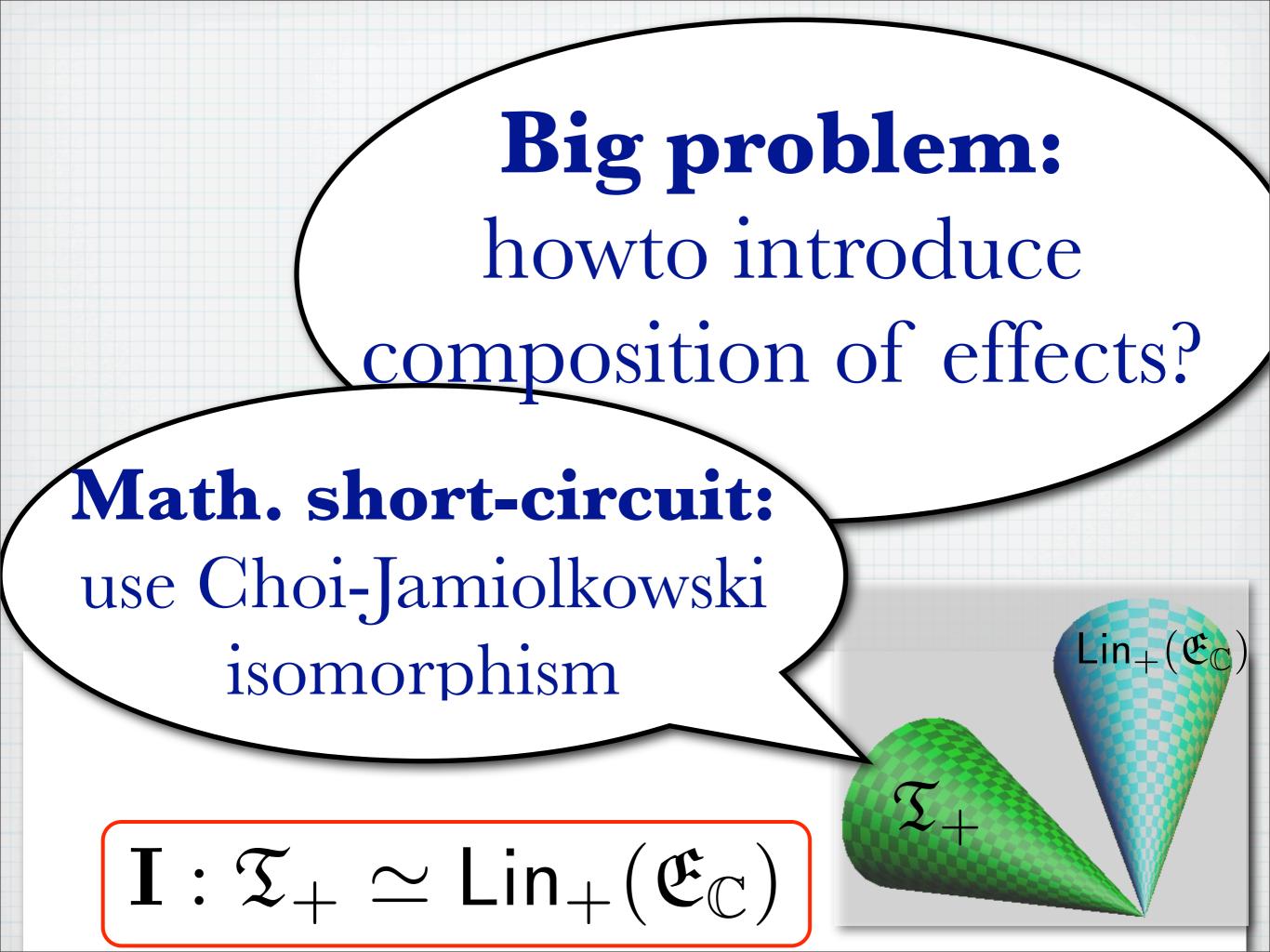
where the operator $E_{ij}(\psi)$ is given by

$$E_{ij}(\psi) = |i_0\rangle\langle i| \otimes |j_0\rangle\langle \psi^{-1*}(j)|, \qquad (11)$$

and the proportionality constant is given by

$$\kappa = e^{i\theta} \sqrt{\frac{p_A(\psi)}{\langle |i_0, j_0\rangle \rangle \langle \langle i_0, j_0| \rangle}}.$$
 (12)

Since A_{ii} is written only in terms of output ensemble averages, it can be estimated through quantum tomography. Quantum tomography [10,11] is a method to estimate the ensemble average $\langle H \rangle$ of any arbitrary operator H on \mathcal{H} by using only measurement outcomes of a quorum of observables $\{O(l)\}$. A quorum is just a set of operators $\{O(l)\}$ which are observable (i.e., have orthonormal resolution) and span the linear space of operators on \mathcal{H} . This means that any operator H can be expanded as $H = \sum_{l} \operatorname{Tr}[Q^{\dagger}(l)H]O(l)$, where $\{Q(l)\}$ and $\{O(l)\}$ form a biorthogonal set such that $\text{Tr}[Q^{\dagger}(i)O(j)] = \delta_{ij}$. Hence, the tomographic estimation of the ensemble average $\langle H \rangle$ is obtained as the double average—over both the ensemble and the quorum—of the unbiased $- \Gamma = + (1) + - T = - (1)$ • 1 1 1 1

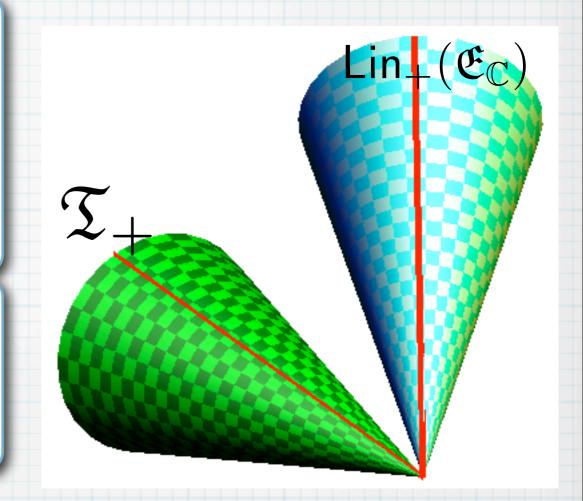


$\mathcal{G}_{\text{Heory group}}$ CJ Isomorphism \Rightarrow composition of effects

Effects are identified with "atomic" events

(apart from a phase) i.e. events that cannot be written as sum of other events

AE (Atomicity of evolution): the composition of "atomic" events is atomic



One can prove that the phase (two-cocycle) is trivial.

Introduce the generalized transformation via the polar identity:

$$\mathscr{T}_{a,b} := \frac{1}{4} \sum_{k=0} i^k \mathscr{T}_{a+i^k b}$$

3

compositon of effects as: $ab = e \circ \mathscr{T}_{e,a} \circ \mathscr{T}_{e,b}$

